

Continuum Calogero–Moser models

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The CCM models

- Continuum Calogero–Moser (CCM) equations:

$$i \frac{d}{dt} q = -q'' \pm 2iq C_+ (|q|^2)'$$

for $q(t)$ in the Hardy space

$$L^2_+ := \{f \in L^2 : \widehat{f}(\xi) = 0 \text{ for } \xi < 0\}$$

- Cauchy–Szegő projection C_+ :

$$\widehat{C_+ f}(\xi) := 1_{[0, \infty)}(\xi) \widehat{f}(\xi)$$

Physical origins

- Focusing CCM: continuum limit of Calogero–Moser particle system
[Abanov–Bettelheim–Wiegmann '09]

$$x_1(t), \dots, x_n(t) \quad \text{solve CM:} \quad \frac{d^2 x_j}{dt^2} = \sum_{k \neq j} \frac{1}{(x_j - x_k)^3}$$

$$\rightsquigarrow \quad \rho(t, x) = \sum \delta(x - x_j(t)), \quad v(t, x) \quad \text{solve} \quad \partial_t \rho + (\rho v)' = 0$$

$$\rightsquigarrow \quad q(t, x) = \sqrt{\rho} e^{i \int^x (v + \pi \rho) dx} \quad \text{solves focusing CCM}$$

- Defocusing CCM: modulation theory for the setting of the Benjamin–Ono equation [Pelinovsky '95]

$$u(t, x) = \epsilon^{-1} q(t, x) e^{i\theta(t, x; \epsilon)} + O(1) \quad \text{solves ILW}$$

$$\xRightarrow{\epsilon \rightarrow 0} \quad q(t, x) \quad \text{solves Intermediate NLS}$$

$$\xRightarrow{\text{depth} \rightarrow \infty} \quad q(t, x) \quad \text{solves defocusing CCM}$$

Conserved quantities

- Mass:

$$M(q) = \int |q|^2 dx$$

- Momentum:

$$P(q) = \int -i\bar{q}q' \mp \frac{1}{2}|q|^4 dx$$

- Hamiltonian:

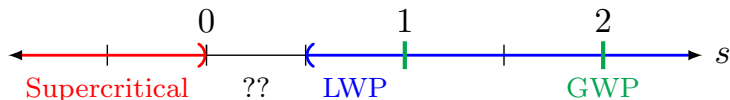
$$H(q) = \frac{1}{2} \int |q' \mp iqC_+(|q|^2)|^2 dx$$

⋮

- Completely integrable: infinite sequence of conserved quantities

Defocusing case

- Well-posedness in H^s spaces on \mathbb{R} and \mathbb{T} : [Gérard–Lenzmann '22]

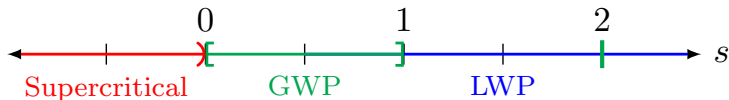


- Scaling symmetry:

$$q(t, x) \mapsto q_\lambda(t, x) = \sqrt{\lambda} q(\lambda^2 t, \lambda x) \quad \text{for } \lambda > 0$$

- Critical H^s regularity is $s = 0$

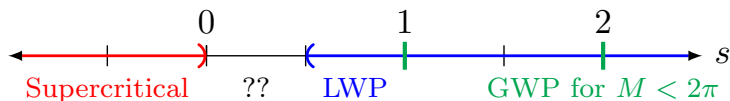
Theorem 1 (Killip–L.–Vişan ‘23). Fix $0 \leq s < 1$. The defocusing CCM equation is globally well-posed in $H_+^s(\mathbb{R})$.



- Analogous result on \mathbb{T} already known [Badreddine ‘23]

Focusing case

- Well-posedness in H^s spaces on \mathbb{R} and \mathbb{T} : [Gérard–Lenzmann '22]



- Solitons:

$$q(t, x) = \sqrt{\lambda} R(\lambda x + x_0), \quad R(x) = \frac{\sqrt{2}}{x + i}$$

- Mass is $\|q\|_{L^2}^2 = \|R\|_{L^2}^2 = 2\pi$ for all $\lambda > 0$

- N -soliton resembles N interacting solitons

$$q(t, x) = \sum_{j=1}^N \frac{a_j(t)}{x - z_j(t)}$$

- The poles $z_j(t) \in \mathbb{C}_-$ solve a complexified CM particle system:
[Gérard–Lenzmann '22]

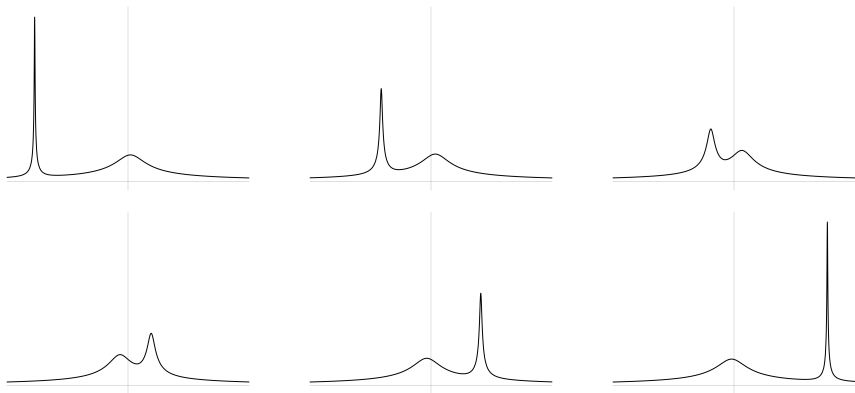
$$\frac{d^2 z_j}{dt^2} = \sum_{k \neq j} \frac{8}{(z_j - z_k)^3}$$

- Mass of an N -soliton is $2\pi N$

- N -solitons with $N \geq 2$ exhibit **turbulent behavior**:

$$\|q(t)\|_{H^s} \sim |t|^{2s} \quad \text{as } t \rightarrow \pm\infty$$

for any $s > 0$ [Gérard–Lenzmann '22]



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$$\|q(t)\|_{H^s} \sim |t|^{2s} \quad \text{as } t \rightarrow \pm\infty$$

for any $s > 0$ [Gérard–Lenzmann ‘22]

- Mass of an N -soliton is $2\pi N$
- There exist solutions with mass $2\pi + \epsilon$ that either exhibit similar behavior, or blow up in finite time [Hogan–Kowalski ‘24]
- There exist smooth solutions with mass $2\pi + \epsilon$ so that $\|q(t)\|_{H^1}$ blows up in finite time [Kim–Kim–Kwon ‘24]

Theorem 2 (Killip–L.–Vişan ‘23). Fix $0 \leq s < 1$. The focusing CCM equation is globally well-posed in the space

$$\{q \in H_+^s(\mathbb{R}) : \|q\|_{L^2}^2 < 2\pi\}.$$

- Analogous result on \mathbb{T} already known [Badreddine ‘23]
- There exist smooth solutions with mass $2\pi + \epsilon$ so that $\|q(t)\|_{H^1}$ blows up in finite time [Kim–Kim–Kwon ‘24]

Equicontinuity

- $\mathcal{F} \subseteq L^2_+$ is **equicontinuous**

$$\iff \sup_{q \in \mathcal{F}} \sup_{|y| < \delta} \|q(\cdot + y) - q(\cdot)\|_{L^2} \rightarrow 0 \quad \text{as } \delta \rightarrow 0$$

$$\iff \sup_{q \in \mathcal{F}} \int_{\kappa}^{\infty} |\widehat{q}(\xi)|^2 d\xi \rightarrow 0 \quad \text{as } \kappa \rightarrow \infty$$

- At critical regularity,

Equicontinuity of orbits	\iff	No concentration in physical space
	\iff	No blowup

Theorem 3 (Defocusing case). If $\mathcal{F} \subseteq L^2_+(\mathbb{R})$ is bounded and equicontinuous, then the set of orbits

$$\mathcal{F}^* := \{q(t) : q(0) \in \mathcal{F}, t \in \mathbb{R}\}$$

for defocusing CCM is also bounded and equicontinuous.

Theorem 4 (Focusing case). If $\mathcal{F} \subseteq L^2_+(\mathbb{R})$ is equicontinuous and

$$\sup_{q \in \mathcal{F}} \|q\|_{L^2}^2 < 2\pi,$$

then the set of orbits \mathcal{F}^* for focusing CCM is bounded and equicontinuous.

- The constant 2π here is sharp

Lax pair

- Lax pair:

$$q(t) \text{ solves CCM} \quad \Longleftrightarrow \quad \frac{d}{dt}L = [P, L],$$

$$L := -i\partial \mp qC_+\bar{q}, \quad P := \dots \quad \text{on } L_+^2$$

- L is symmetric, P is antisymmetric
- Formally,

$$U(t) \text{ solves } \frac{d}{dt}U = PU$$

$$\implies L(t) = U(t)L(0)U(t)^* \quad \text{with } U(t) \text{ unitary}$$

- So spectrum of $L(t)$ is conserved

Proof of equicontinuity

Goal: If $\mathcal{F} \subseteq L_+^2$ is equicontinuous and

$$\sup_{q \in \mathcal{F}} \|q\|_{L^2}^2 < \begin{cases} 2\pi & \text{(focusing)} \\ \infty & \text{(defocusing)} \end{cases}$$

then the set of orbits \mathcal{F}^* is also equicontinuous.

- The spectrum of

$$L_q := -i\partial \mp qC_+ \bar{q} \quad \text{on } L_+^2$$

is conserved in time (formally)

Proposition. The quantity

$$\mathrm{tr}(R_q - R_0), \quad \text{where} \quad R_q = (L_q + \kappa)^{-1},$$

is finite for $q \in L_+^2$. Moreover, it is conserved for H_+^∞ solutions.

- Expand as a series in q :

$$\mathrm{tr}(R_q - R_0) = \sum_{\ell \geq 1} (\pm 1)^\ell \mathrm{tr} \{ (R_0 q C_+ \bar{q})^\ell R_0 \}$$

- The leading order term is

$$\mathrm{tr} (R_0 q C_+ \bar{q} R_0) = \frac{1}{2\pi} \int_0^\infty \frac{|\widehat{q}(\xi)|^2}{\xi + \kappa} d\xi$$

- Build a conserved quantity to estimate the high frequencies:

$$\beta(\kappa, q) := \|q\|_{L^2}^2 \mp 2\pi\kappa \operatorname{tr}(R_q - R_0)$$

- Quadratic term is

$$\beta^{[2]} := \|q\|_{L^2}^2 - 2\pi\kappa \operatorname{tr}(R_0 q C_+ \bar{q} R_0) = \int_0^\infty \frac{\xi}{\xi + \kappa} |\widehat{q}(\xi)|^2 d\xi$$

- For $\mathcal{F} \subseteq L_+^2$ bounded,

$$\begin{aligned} \mathcal{F} \text{ is equicontinuous} & \iff \sup_{q \in \mathcal{F}} \beta^{[2]}(\kappa, q) \rightarrow 0 \quad \text{as } \kappa \rightarrow \infty \\ & \implies \sup_{q \in \mathcal{F}} \beta(\kappa, q) \rightarrow 0 \quad \text{as } \kappa \rightarrow \infty \end{aligned}$$

Defocusing case

$$L_q = -i\partial + qC_+\bar{q}, \quad qC_+\bar{q} \geq 0$$

$$\begin{aligned} \Rightarrow \quad -\operatorname{tr} \left\{ \frac{1}{L_q + \kappa} - \frac{1}{L_0 + \kappa} \right\} &= \operatorname{tr} \left\{ \frac{1}{\sqrt{L_0 + \kappa}} \frac{qC_+\bar{q}}{L_0 + qC_+\bar{q} + \kappa} \frac{1}{\sqrt{L_0 + \kappa}} \right\} \\ &\leq \operatorname{tr} \left\{ \frac{1}{\sqrt{L_0 + \kappa}} \frac{qC_+\bar{q}}{L_0 + \kappa} \frac{1}{\sqrt{L_0 + \kappa}} \right\} \end{aligned}$$

$$\Rightarrow \quad \beta(\kappa, q) \geq \beta^{[2]}(\kappa, q)$$

$$\Rightarrow \quad \beta^{[2]}(\kappa, q(t)) \leq \beta(\kappa, q(t)) \equiv \beta(\kappa, q(0)) \xrightarrow{\kappa \rightarrow \infty} 0$$

Focusing case

$$L_q = -i\partial - qC_+\bar{q}, \quad qC_+\bar{q} \leq \frac{M}{2\pi}L_0$$

$$\implies L_q \geq (1 - \frac{M}{2\pi})L_0$$

$$\implies (L_q + \kappa)^{-1} \leq \frac{2\pi}{2\pi - M}(L_0 + \kappa)^{-1}$$

$$\implies \dots$$

Explicit formula

Theorem 5 (Killip–L.–Vişan ‘23). If $q^0 \in L_+^2$ (and $\|q^0\|_{L^2}^2 < 2\pi$ in the focusing case), then the global solution $q(t)$ to CCM satisfies

$$q(t, z) = \frac{1}{2\pi i} I_+ \{ (X + 2tL_{q^0} - z)^{-1} q^0 \} \quad \forall \operatorname{Im} z > 0.$$

- Here, X is an extension of multiplication-by- x on $L_+^2(\mathbb{R})$
- I_+ is an extension of $q \mapsto \int q \, dx$ on $L_+^2(\mathbb{R})$
- Cubic Szegő equation: [Gérard–Grellier ‘15, Gérard–Pushnitski ‘24]
- Benjamin–Ono equation: [Gérard ‘23]

Proof of well-posedness

Goal: If $\{q_n^0\}_{n \geq 1} \subseteq \mathcal{S}$ converges in L_+^2 and

$$\sup_n \|q_n^0\|_{L^2}^2 < \begin{cases} 2\pi & (\text{focusing}) \\ \infty & (\text{defocusing}) \end{cases}$$

then the solutions $q(t)$ converge in $C_t([-T, T]; L_+^2)$.

Steps:

1. The explicit formulas for q_n^0 converge for fixed t, z
2. The solutions $q_n(t)$ converge weakly to $q(t)$ in L_+^2 for fixed t
3. The solutions $q_n(t)$ are precompact in $C_t([-T, T]; L_+^2)$

1. The explicit formulas for q_n^0 converge for fixed t, z

- The operator $X + 2tL_0 = X - 2it\partial$ on L_+^2 is maximally accretive
- $A_0 := (X + 2tL_0 - z)^{-1}$ is bounded $L_+^2 \rightarrow L_+^2$ for $\text{Im } z > 0$
- A_0 is also bounded $L_+^2 \rightarrow L_+^\infty$ and $L_+^1 \rightarrow L_+^2$, provided $t \neq 0$
- The series

$$(X + 2tL_q - z)^{-1} = \underset{L^2 \leftarrow L^2}{A_0} \pm \underset{L^2 \leftarrow L^1}{A_0} \underset{L^\infty \leftarrow L^2}{2tqC_+ \bar{q}} \underset{L^\infty \leftarrow L^2}{A_0} + \dots$$

converges for fixed t, z

2. The solutions $q_n(t)$ converge weakly to $q(t)$ in L_+^2 for fixed t

- Fix t , and pass to any subsequence of $\{q_n(t)\}_{n \geq 1}$
- Conservation of mass $\Rightarrow q_n(t) \rightharpoonup \tilde{q}(t)$ in L_+^2 along a subsequence
- Poisson integral formula: For $\operatorname{Im} z > 0$,

$$q_n(t, z) = \int \frac{\operatorname{Im} z}{\pi |x - z|^2} q_n(t, x) dx$$

- So $\tilde{q}(t, z)$ is given by the explicit formula for all $\operatorname{Im} z > 0$
- This does not depend on the subsequence

3. The solutions $q_n(t)$ are precompact in $C_t([-T, T]; L^2_+)$

- Equicontinuity in space: already have
- Equicontinuity in time:

$$\|q_n(t+h) - q_n(t)\|_{C_t L^2_x} \leq |h| \left\| \frac{d}{dt} P_{\leq N} q_n \right\|_{L_t^\infty L^2_x} + 2 \|P_{> N} q_n\|_{C_t L^2_x}$$

- Tightness in space:
 - For fixed $b > 0$, the functions $x \mapsto q_n(t, x + ib)$ are tight
 - Poisson integral formula: $q_n(t, x + ib) = [e^{-b|\partial|} q_n](t, x)$
 - $\|q_n(t)\|_{L^2(|x| \geq R)} \leq \|P_{\leq N} q_n(t)\|_{L^2(|x| \geq R)} + 2 \|P_{> N} q_n(t)\|_{L^2}$

Dispersive decay

- Dispersive estimate for the free Schrödinger flow:

$$\|e^{it\Delta}q\|_{L^\infty} \lesssim |t|^{-\frac{1}{2}} \|q\|_{L^1}$$

Theorem 6 (Killip–L.–Viřan ‘25+). There exists $\delta > 0$ so that for any initial data $q \in L^2_+(\mathbb{R})$ satisfying

$$\|q\|_{L^2} \leq \delta \quad \text{and} \quad q \in L^1,$$

the corresponding global solution $q(t)$ of CCM satisfies

$$\|q(t)\|_{L^\infty} \lesssim |t|^{-\frac{1}{2}}$$

for all $t \neq 0$.

- Cubic NLS: Dispersive decay for...
 - small q in $H^{1,1}$ [Deift–Zhou ‘03, Kato–Pusateri ‘11]
 - small q in $H^{\frac{1}{2}+, \frac{1}{2}+}$ [Hayashi–Naumkin ‘98]
 - small q in $H^{0, \frac{1}{2}+}$ [Ifrim–Tataru ‘15]
- Benjamin–Ono: Dispersive decay fails for small q in $L^1 \cap L^2$

Proof sketch

- Explicit formula:

$$q(t, z) = \frac{1}{2\pi i} I_+ \{ (X + 2tL_q - z)^{-1} q \} \quad \forall \operatorname{Im} z > 0.$$

- Free resolvent:

$$A_0 = (X + 2tL_0 - z)^{-1}$$

$$[(X + 2tL_0)f]^\wedge(\xi) = (i\partial_\xi + 2t\xi)\hat{f}(\xi) = e^{it\xi^2} i\partial_\xi e^{-it\xi^2} \hat{f}(\xi)$$

$$\implies A_0 = e^{-it\Delta} (X - z)^{-1} e^{it\Delta}$$

$$\begin{aligned}
I_+(A_0 f) &= I_+ [e^{-it\Delta} (X - z)^{-1} e^{it\Delta} f] \\
&= I_+ [(X - z)^{-1} e^{it\Delta} f] \\
&= 2\pi i [e^{it\Delta} f](z)
\end{aligned}$$

- Expand as a series in q :

$$\begin{aligned}
q(t, z) &= \frac{1}{2\pi i} I_+ [(X + 2tL_q - z)^{-1} q] \\
&= \frac{1}{2\pi i} I_+ A_0 q \pm \frac{1}{2\pi i} I_+ A_0 2tq C_+ \bar{q} A_0 q + \dots \\
&= [e^{it\Delta} q \pm e^{it\Delta} 2tq C_+ \bar{q} A_0 q + \dots](z)
\end{aligned}$$

- $A_0 := (X + 2tL_0 - z)^{-1}$ is bounded $L_+^2 \rightarrow L_+^2$ for $\operatorname{Im} z > 0$
- A_0 is also bounded $L_+^1 \rightarrow L_+^\infty$ for $t \neq 0$:

$$\|A_0\|_{L_+^1 \rightarrow L_+^\infty} \lesssim |t|^{-1}$$

$$q(t, z) = \left[\underset{L^\infty \leftarrow L^1}{e^{it\Delta} q} \pm \underset{L^\infty \leftarrow L^1}{e^{it\Delta} 2tqC_+ \bar{q}} \underset{L^\infty \leftarrow L^1}{A_0 q} + \dots \right](z)$$

- Poisson integral formula:

$$f(z) = \int \frac{\operatorname{Im} z}{\pi |x - z|^2} f(x) dx \quad \implies \quad \|f(x + ib)\|_{L_x^\infty} \lesssim \|f\|_{L^\infty}$$

- Conclude:

$$\|q(t, x + ib)\|_{L_x^\infty} \lesssim |t|^{-\frac{1}{2}} \quad \forall b > 0$$

Thank you!