Some advances in the selection problem in ergodic optimization Renaud Leplaideur

April 2, 2024



Joint work with Jairo Mengue, UFRGS Porto Alegre Brazil



Some advances in the selection problem in er

Thermodynamic formalism for dynamical systems

Some advances in the selection problem in er

Orbits

Goal of dynamics : describe orbits for (X, T), X compact metric space, $T : X \rightarrow X$ continuous.

Usually too difficult. Then describe only almost all trajectories. For that purpose, need to select some special invariant measure = aim for thermodynamic formalism.

Definition

(X, T)= Dynamical System, $A : X \to \mathbb{R}$ continuous called *potential*.

• Pressure =
$$P(A) := \sup_{\mu T - inv} \{h_{\mu}(T) + \int A d\mu\},$$

Any measure realizing maximum = equilibrium state for A.

 $h_{\mu}(T)$ = Kolmogorov entropy.

A (10) A (10) A (10) A

If (X, T) uniformly hyperbolic and A Hölder continuous, then there exists a unique equilibrium state for A. Furthermore, $\mathcal{P}(\beta) = \mathcal{P}(\beta.A)$, $\beta \in \mathbb{R}$ is analytic.

my settings : subshift of finite type + A Lipschitz. Easy to check $\mathcal{P}(\beta)$ is convex with asymptote at $+\infty$: Set μ_{β} for unique equilibrium state for β .A.

If (X, T) uniformly hyperbolic and A Hölder continuous, then there exists a unique equilibrium state for A. Furthermore, $\mathcal{P}(\beta) = \mathcal{P}(\beta.A)$, $\beta \in \mathbb{R}$ is analytic.

my settings : subshift of finite type + *A* Lipschitz. Easy to check $\mathcal{P}(\beta)$ is convex with asymptote at $+\infty$: Set μ_{β} for unique equilibrium state for β .*A*.

< 回 > < 三 > < 三 >

If (X, T) uniformly hyperbolic and A Hölder continuous, then there exists a unique equilibrium state for A. Furthermore, $\mathcal{P}(\beta) = \mathcal{P}(\beta.A)$, $\beta \in \mathbb{R}$ is analytic.

my settings : subshift of finite type + *A* Lipschitz. Easy to check $\mathcal{P}(\beta)$ is convex with asymptote at $+\infty$:



Set μ_{β} for unique equilibrium state for β .*A*.

If (X, T) uniformly hyperbolic and A Hölder continuous, then there exists a unique equilibrium state for A. Furthermore, $\mathcal{P}(\beta) = \mathcal{P}(\beta.A)$, $\beta \in \mathbb{R}$ is analytic.

my settings : subshift of finite type + *A* Lipschitz. Easy to check $\mathcal{P}(\beta)$ is convex with asymptote at $+\infty$:



Set μ_{β} for unique equilibrium state for β .*A*.

On the road to ergodic optimization

Here
$$m(A) = \max_{\mu \ T-inv} \left\{ \int A d\mu \right\}.$$

Definition

An invariant measure is said to be A-maximizing if it realizes maximum.

h= maximal entropy among *A*-maximizing measures.

Theorem (Folklore)

Any accumulation for μ_{β} is A-maximizing with entropy h.

Question

Is there convergence for μ_{β} ? If yes why ? If not why ? How selects different accumulation points.

A (1) > A (2) > A

On the road to ergodic optimization

Here
$$m(A) = \max_{\mu \ T-inv} \left\{ \int A d\mu \right\}.$$

Definition

An invariant measure is said to be A-maximizing if it realizes maximum.

h= maximal entropy among *A*-maximizing measures.

Theorem (Folklore)

Any accumulation for μ_{β} is A-maximizing with entropy h.

Question

Is there convergence for μ_{β} ? If yes why? If not why? How selects different accumulation points.

A (10) > A (10) > A (10)

One example (Baraviera-L-Lopes)

Considering in $\Sigma := \{0, 1, 2\}^{\mathbb{N}}$ and

$$A(x) = \begin{cases} -d(x,0^{\infty}) & \text{if } x \in [0] \\ -3d(x,1^{\infty}) & \text{if } x \in [1] \\ -\alpha < 0 & \text{if } x \in [2]. \end{cases}$$

Theorem

Let ρ be the golden mean $\rho := \frac{1 + \sqrt{5}}{2}$. Then

- for lpha > 1, μ_{eta} converges to $rac{1}{2}(\delta_{0^{\infty}} + \delta_{1^{\infty}})$ as eta goes to $+\infty$,
- (a) for $\alpha = 1$, μ_{β} converges to $\frac{1}{1+\rho^2}(\rho^2 \delta_{0\infty} + \delta_{1\infty})$ as β goes to $+\infty$,
- 3) for $\alpha < 1$, μ_{β} converges to $\delta_{0^{\infty}}$ as β goes to $+\infty$.

A (10) < A (10) < A (10) </p>

One example (Baraviera-L-Lopes)

Considering in $\Sigma := \{0, 1, 2\}^{\mathbb{N}}$ and

$$\mathcal{A}(x) = \begin{cases} -d(x,0^{\infty}) & \text{if } x \in [0] \\ -3d(x,1^{\infty}) & \text{if } x \in [1] \\ -\alpha < 0 & \text{if } x \in [2]. \end{cases}$$

Theorem

Let ρ be the golden mean $\rho := \frac{1 + \sqrt{5}}{2}$. Then • for $\alpha > 1$, μ_{β} converges to $\frac{1}{2}(\delta_{0\infty} + \delta_{1\infty})$ as β goes to $+\infty$, • for $\alpha = 1$, μ_{β} converges to $\frac{1}{1+\rho^2}(\rho^2\delta_{0\infty} + \delta_{1\infty})$ as β goes to $+\infty$, • for $\alpha < 1$, μ_{β} converges to $\delta_{0\infty}$ as β goes to $+\infty$.

A good way to see ergodic optimization

Some advances in the selection problem in er

< 回 > < 三 > < 三 >



◆□ → ◆□ → ▲目 → ▲目 → ▲□ →

On one island = can walk without being wet. To go from 1 island to another, need to go to the water = to pay.

Definition

Seeing A as a cost, Aubry set = some islands in X, invariant for T where we can walk without paying.

Good tool for that = calibrated subaction.

Remind transfert operator gives all informations in thermodynamic formalism:

$$L(g)(x) = \sum_{y, T(y)=x} e^{\beta A(y)} g(y).$$

 $e^{\mathcal{P}(\beta)}$ = unique single dominating eigenvalue, H_{β} = eigenfunction, ν_{β} = eigenmeasure (for L^*), $\mu_{\beta} = H_{\beta} \otimes \nu_{\beta}$.

4 **A** N A **B** N A

On one island = can walk without being wet. To go from 1 island to another, need to go to the water = to pay.

Definition

Seeing A as a cost, Aubry set = some islands in X, invariant for T where we can walk without paying.

Good tool for that = calibrated subaction.

Remind transfert operator gives all informations in thermodynamic formalism:

$$L(g)(x) = \sum_{y, T(y)=x} e^{\beta A(y)} g(y).$$

 $e^{\mathcal{P}(\beta)}$ = unique single dominating eigenvalue, H_{β} = eigenfunction, ν_{β} = eigenmeasure (for L^*), $\mu_{\beta} = H_{\beta} \otimes \nu_{\beta}$.

A (1) > A (1) > A

On one island = can walk without being wet. To go from 1 island to another, need to go to the water = to pay.

Definition

Seeing A as a cost, Aubry set = some islands in X, invariant for T where we can walk without paying.

Good tool for that = calibrated subaction.

Remind transfert operator gives all informations in thermodynamic formalism:

$$L(g)(x) = \sum_{y, T(y)=x} e^{\beta A(y)} g(y).$$

 $e^{\mathcal{P}(\beta)}$ = unique single dominating eigenvalue, H_{β} = eigenfunction, ν_{β} = eigenmeasure (for L^*), $\mu_{\beta} = H_{\beta} \otimes \nu_{\beta}$.

A I > A I > A

On one island = can walk without being wet. To go from 1 island to another, need to go to the water = to pay.

Definition

Seeing A as a cost, Aubry set = some islands in X, invariant for T where we can walk without paying.

Good tool for that = calibrated subaction.

Remind transfert operator gives all informations in thermodynamic formalism:

$$L(g)(x) = \sum_{y, T(y)=x} e^{\beta A(y)} g(y).$$

 $e^{\mathcal{P}(\beta)}$ = unique single dominating eigenvalue, H_{β} = eigenfunction, ν_{β} = eigenmeasure (for L^*), $\mu_{\beta} = H_{\beta} \otimes \nu_{\beta}$.

On one island = can walk without being wet. To go from 1 island to another, need to go to the water = to pay.

Definition

Seeing A as a cost, Aubry set = some islands in X, invariant for T where we can walk without paying.

Good tool for that = calibrated subaction.

Remind transfert operator gives all informations in thermodynamic formalism:

$$L(g)(x) = \sum_{y, T(y)=x} e^{\beta A(y)} g(y).$$

 $e^{\mathcal{P}(\beta)}$ = unique single dominating eigenvalue, H_{β} = eigenfunction, ν_{β} = eigenmeasure (for L^*), $\mu_{\beta} = H_{\beta} \otimes \nu_{\beta}$.

On one island = can walk without being wet. To go from 1 island to another, need to go to the water = to pay.

Definition

Seeing A as a cost, Aubry set = some islands in X, invariant for T where we can walk without paying.

Good tool for that = calibrated subaction.

Remind transfert operator gives all informations in thermodynamic formalism:

$$L(g)(x) = \sum_{y, T(y)=x} e^{\beta A(y)} g(y).$$

 $e^{\mathcal{P}(\beta)}$ = unique single dominating eigenvalue, H_{β} = eigenfunction, ν_{β} = eigenmeasure (for L^*), $\mu_{\beta} = H_{\beta} \otimes \nu_{\beta}$.

(人間) トイヨト イヨト

On one island = can walk without being wet. To go from 1 island to another, need to go to the water = to pay.

Definition

Seeing A as a cost, Aubry set = some islands in X, invariant for T where we can walk without paying.

Good tool for that = calibrated subaction.

Remind transfert operator gives all informations in thermodynamic formalism:

$$L(g)(x) = \sum_{y, T(y)=x} e^{\beta A(y)} g(y).$$

 $e^{\mathcal{P}(\beta)}$ = unique single dominating eigenvalue, H_{β} = eigenfunction, ν_{β} = eigenmeasure (for L^*), $\mu_{\beta} = H_{\beta} \otimes \nu_{\beta}$.

$$e^{\mathcal{P}(\beta)}H_{\beta}(x) = \sum_{y} e^{\beta \cdot A(y)}H_{\beta}(y)$$

$$e^{\beta\left(\frac{1}{\beta}\mathcal{P}(\beta) + \frac{1}{\beta}\log H_{\beta}(x)\right)} = \sum_{y} e^{\beta(A(y) + \frac{1}{\beta}\log H_{\beta}(y))}$$

$$\frac{1}{\beta}\log + \beta \to +\infty$$

$$m(A) + V(x) = \max_{y} \{V(y) + A(y)\}$$

Rewritten as
$$A(y) + V(y) - V \circ \underbrace{T(y)}_{=x} - m(A) \leq 0$$
.

Up to coboundary + constant =no change for equilibrium states, A always non positive. Aubry set \approx bigger invariant set where $A \equiv 0$.

A (10) A (10)

1

$$e^{\mathcal{P}(\beta)}H_{\beta}(x) = \sum_{y} e^{\beta \cdot A(y)}H_{\beta}(y)$$

$$e^{\beta\left(\frac{1}{\beta}\mathcal{P}(\beta) + \frac{1}{\beta}\log H_{\beta}(x)\right)} = \sum_{y} e^{\beta(A(y) + \frac{1}{\beta}\log H_{\beta}(y))}$$

$$\frac{1}{\beta}\log + \beta \to +\infty$$

$$m(A) + V(x) = \max_{y} \{V(y) + A(y)\}$$

Rewritten as
$$A(y) + V(y) - V \circ \underbrace{T(y)}_{=x} - m(A) \leq 0$$
.

Up to coboundary + constant =no change for equilibrium states, A always non positive. Aubry set \approx bigger invariant set where $A \equiv 0$.

・ 同 ト ・ ヨ ト ・ ヨ ト

$$e^{\mathcal{P}(\beta)}H_{\beta}(x) = \sum_{y} e^{\beta \cdot A(y)}H_{\beta}(y)$$

$$e^{\beta\left(\frac{1}{\beta}\mathcal{P}(\beta) + \frac{1}{\beta}\log H_{\beta}(x)\right)} = \sum_{y} e^{\beta(A(y) + \frac{1}{\beta}\log H_{\beta}(y))}$$

$$\frac{1}{\beta}\log + \beta \to +\infty$$

$$m(A) + V(x) = \max_{y} \{V(y) + A(y)\}$$

Rewritten as
$$A(y) + V(y) - V \circ \underbrace{T(y)}_{=x} - m(A) \leq 0$$
.

Up to coboundary + constant =no change for equilibrium states, A always non positive. Aubry set \approx bigger invariant set where $A \equiv 0$.

4 **A** N A **B** N A **B** N

$$e^{\mathcal{P}(\beta)}H_{\beta}(x) = \sum_{y} e^{\beta.A(y)}H_{\beta}(y)$$

$$e^{\beta\left(\frac{1}{\beta}\mathcal{P}(\beta) + \frac{1}{\beta}\log H_{\beta}(x)\right)} = \sum_{y} e^{\beta(A(y) + \frac{1}{\beta}\log H_{\beta}(y))}$$

$$\frac{1}{\beta}\log + \beta \to +\infty$$

$$m(A) + V(x) = \max_{y} \{V(y) + A(y)\}$$

Rewritten as
$$A(y) + V(y) - V \circ \underbrace{T(y)}_{=x} - m(A) \le 0$$
.

Up to coboundary + constant =no change for equilibrium states, A always non positive. Aubry set \approx bigger invariant set where $A \equiv 0$.

< ロ > < 同 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ >

$$e^{\mathcal{P}(\beta)}H_{\beta}(x) = \sum_{y} e^{\beta \cdot A(y)}H_{\beta}(y)$$

$$e^{\beta\left(\frac{1}{\beta}\mathcal{P}(\beta) + \frac{1}{\beta}\log H_{\beta}(x)\right)} = \sum_{y} e^{\beta(A(y) + \frac{1}{\beta}\log H_{\beta}(y))}$$

$$\frac{1}{\beta}\log + \beta \to +\infty$$

$$m(A) + V(x) = \max_{y} \{V(y) + A(y)\}$$

Rewritten as
$$A(y) + V(y) - V \circ \underbrace{T(y)}_{=x} - m(A) \leq 0$$

Up to coboundary + constant =no change for equilibrium states, A always non positive. Aubry set \approx bigger invariant set where $A \equiv 0$.

$$e^{\mathcal{P}(\beta)}H_{\beta}(x) = \sum_{y} e^{\beta.A(y)}H_{\beta}(y)$$

$$e^{\beta\left(\frac{1}{\beta}\mathcal{P}(\beta) + \frac{1}{\beta}\log H_{\beta}(x)\right)} = \sum_{y} e^{\beta(A(y) + \frac{1}{\beta}\log H_{\beta}(y))}$$

$$\frac{1}{\beta}\log + \beta \to +\infty$$

$$m(A) + V(x) = \max_{y} \{V(y) + A(y)\}$$

Rewritten as
$$A(y) + V(y) - V \circ \underbrace{T(y)}_{=x} - m(A) \leq 0$$

Up to coboundary + constant =no change for equilibrium states, A always non positive. Aubry set \approx bigger invariant set where $A \equiv 0$.

A (1) > A (1) > A

$$e^{\mathcal{P}(\beta)}H_{\beta}(x) = \sum_{y} e^{\beta.A(y)}H_{\beta}(y)$$

$$e^{\beta\left(\frac{1}{\beta}\mathcal{P}(\beta) + \frac{1}{\beta}\log H_{\beta}(x)\right)} = \sum_{y} e^{\beta(A(y) + \frac{1}{\beta}\log H_{\beta}(y))}$$

$$\frac{1}{\beta}\log + \beta \to +\infty$$

$$m(A) + V(x) = \max_{y} \{V(y) + A(y)\}$$

Rewritten as
$$A(y) + V(y) - V \circ \underbrace{T(y)}_{=x} - m(A) \leq 0$$

Up to coboundary + constant =no change for equilibrium states, A always non positive. Aubry set \approx bigger invariant set where $A \equiv 0$.

< 回 > < 三 > < 三 >

Theorem (L-Mengue, in progress)

If the Aubry set is a subshift of finite type, then the pressure function converges exponentially fast to the asymptote:

$$\lim_{eta
ightarrow+\infty}rac{1}{eta}\log(\mathcal{P}(eta)-h)=\gamma<0.$$

Furthermore, γ is identified as the unique eigenvalue for the matrix of cost between the irreducible components for the Aubry set, within the **Max-Plus formalism**

$$\begin{array}{l} \mathsf{Max-plus}: \oplus = \max, \otimes = +.\\ \mathsf{lim} \ \frac{1}{\beta} \log(e^{a\beta} + e^{c\beta}) = \max\{a, c\}. \ + \to \max.\\ \mathsf{lim} \ \frac{1}{\beta} \log(e^{a\beta} \times e^{c\beta}) = a + c. \ \times \to +. \end{array}$$

Theorem (L-Mengue, in progress)

If the Aubry set is a subshift of finite type, then the pressure function converges exponentially fast to the asymptote:

$$\lim_{eta
ightarrow+\infty}rac{1}{eta}\log(\mathcal{P}(eta)-h)=\gamma<0.$$

Furthermore, γ is identified as the unique eigenvalue for the matrix of cost between the irreducible components for the Aubry set, within the Max-Plus formalism

$$\begin{array}{l} \mathsf{Max-plus}: \oplus = \mathsf{max}, \otimes = +.\\ \mathsf{lim} \ \frac{1}{\beta} \log(e^{a\beta} + e^{c\beta}) = \mathsf{max}\{a, c\}. \ + \to \mathsf{max}.\\ \mathsf{lim} \ \frac{1}{\beta} \log(e^{a\beta} \times e^{c\beta}) = a + c. \ \times \to +. \end{array}$$

- **→ → →**

Theorem (L-Mengue, in progress)

If the Aubry set is a subshift of finite type, then the pressure function converges exponentially fast to the asymptote:

$$\lim_{eta o +\infty} rac{1}{eta} \log(\mathcal{P}(eta) - h) = \gamma < 0.$$

Furthermore, γ is identified as the unique eigenvalue for the matrix of cost between the irreducible components for the Aubry set, within the Max-Plus formalism

 $\begin{aligned} & \mathsf{Max-plus}: \oplus = \mathsf{max}, \otimes = +. \\ & \lim \frac{1}{\beta} \log(e^{a\beta} + e^{c\beta}) = \mathsf{max}\{a, c\}. + \to \mathsf{max}. \\ & \lim \frac{1}{\beta} \log(e^{a\beta} \times e^{c\beta}) = a + c. \times \to +. \end{aligned}$

Theorem (L-Mengue, in progress)

If the Aubry set is a subshift of finite type, then the pressure function converges exponentially fast to the asymptote:

$$\lim_{eta
ightarrow+\infty}rac{1}{eta}\log(\mathcal{P}(eta)-h)=\gamma<0.$$

Furthermore, γ is identified as the unique eigenvalue for the matrix of cost between the irreducible components for the Aubry set, within the Max-Plus formalism

$$\begin{array}{l} \mathsf{Max-plus}: \oplus = \max, \otimes = +.\\ \mathsf{lim} \ \frac{1}{\beta} \log(\boldsymbol{e}^{\boldsymbol{a}\beta} + \boldsymbol{e}^{\boldsymbol{c}\beta}) = \max\{\boldsymbol{a}, \boldsymbol{c}\}. \ + \to \max.\\ \mathsf{lim} \ \frac{1}{\beta} \log(\boldsymbol{e}^{\boldsymbol{a}\beta} \times \boldsymbol{e}^{\boldsymbol{c}\beta}) = \boldsymbol{a} + \boldsymbol{c}. \ \times \to +. \end{array}$$

- **→ → →**

Another good way to see ergodic optimization

Some advances in the selection problem in er



◆□ → ◆□ → ▲目 → ▲目 → ▲□ →

Aubry set has zero measure for every μ_{β} .

Question

Can we however see emerge the properties of the convergence/accumulation point from the transfer operator at $\beta < +\infty$, as island emerging when tide is going out ?

Corollary

If we know that for some A μ_{β} converge, can we get the same result if we change a little bit A ?

$$L(g)(x) = \sum_{y, T(y)=x} e^{\beta A(y)} g(y).$$

Need control on Lipschitz norm for the potential: Blows up as $\beta \to +\infty$.

A I > A = A A

Aubry set has zero measure for every μ_{β} .

Question

Can we however see emerge the properties of the convergence/accumulation point from the transfer operator at $\beta < +\infty$, as island emerging when tide is going out ?

Corollary

If we know that for some A μ_{β} converge, can we get the same result if we change a little bit A ?

$$L(g)(x) = \sum_{y, T(y)=x} e^{\beta A(y)} g(y).$$

Need control on Lipschitz norm for the potential: Blows up as $\beta \to +\infty$.

A (10) A (10)

Aubry set has zero measure for every μ_{β} .

Question

Can we however see emerge the properties of the convergence/accumulation point from the transfer operator at $\beta < +\infty$, as island emerging when tide is going out ?

Corollary

If we know that for some A μ_β converge, can we get the same result if we change a little bit A ?

$$L(g)(x) = \sum_{y, T(y)=x} e^{\beta A(y)} g(y).$$

Need control on Lipschitz norm for the potential: Blows up as $\beta \to +\infty$.

Aubry set has zero measure for every μ_{β} .

Question

Can we however see emerge the properties of the convergence/accumulation point from the transfer operator at $\beta < +\infty$, as island emerging when tide is going out ?

Corollary

If we know that for some A μ_β converge, can we get the same result if we change a little bit A ?

$$L(g)(x) = \sum_{y, T(y)=x} e^{\beta A(y)} g(y).$$

Need control on Lipschitz norm for the potential: Blows up as $\beta \to +\infty$.

< ロ > < 同 > < 回 > < 回 >

Aubry set has zero measure for every μ_{β} .

Question

Can we however see emerge the properties of the convergence/accumulation point from the transfer operator at $\beta < +\infty$, as island emerging when tide is going out ?

Corollary

If we know that for some A μ_{β} converge, can we get the same result if we change a little bit A ?

$$L(g)(x) = \sum_{y, T(y)=x} e^{\beta A(y)} g(y).$$

Need control on Lipschitz norm for the potential: Blows up as $\beta \to +\infty$.

Result 2

Only for $X = \{0, 1\}^{\mathbb{N}}$ = full 2 shift & A = special Walters potential. First Theorem holds.

Theorem (L-Mengue, in progress)

If B_{β} is a family of Lipschitz potentials going exponentially to 0 faster than $e^{\gamma,\beta}$ (= speed for pressure), then,

$$\lim_{\beta \to +\infty} \mu_{\beta.A+B_{\beta}} = \lim_{\beta \to +\infty} \mu_{\beta.A},$$

+ more properties.

Result 2

Only for $X = \{0, 1\}^{\mathbb{N}}$ = full 2 shift & A = special Walters potential. First Theorem holds.

Theorem (L-Mengue, in progress)

If B_{β} is a family of Lipschitz potentials going exponentially to 0 faster than $e^{\gamma.\beta}$ (= speed for pressure), then,

$$\lim_{\beta \to +\infty} \mu_{\beta.A+B_{\beta}} = \lim_{\beta \to +\infty} \mu_{\beta.A},$$

+ more properties.

< 回 > < 三 > < 三 >

Set Ω =Aubry set = subshift of finite type.

Set S(x, y)= Mañé potential

$$S(x,y) = \lim_{\varepsilon \to 0} \sup\{S_n(A)(z), \sigma^n(z) = y, \ d(x,y) \le \varepsilon\}.$$

Then,

Irreducible components for Ω as subshift of finite type coincide with equivalence classes for

$$S(x,y)+S(y,x)=0.$$



不同 トイモトイモ

Set Ω =Aubry set = subshift of finite type. Set S(x, y)= Mañé potential

$$S(x,y) = \lim_{\varepsilon \to 0} \sup \{ S_n(A)(z), \sigma^n(z) = y, \ d(x,y) \le \varepsilon \}.$$

Then,

 Irreducible components for Ω as subshift of finite type coincide with equivalence classes for

$$S(x,y)+S(y,x)=0.$$

Set
$$\Omega = \underbrace{\Omega_1 \sqcup \ldots \sqcup \Omega_k}_{\text{entropy } h} \sqcup \underbrace{\Omega_{k+1} \sqcup \ldots \sqcup \Omega_l}_{\text{entropy } < h}$$
Set $a_{ij} := \sup_{x \in \Omega_j} \sup_{y \notin \Omega_i, \ \sigma(y) \in \Omega_i} S(x, y).$

< 回 > < 三 > < 三 >

Set Ω =Aubry set = subshift of finite type. Set S(x, y)= Mañé potential

$$S(x,y) = \lim_{\varepsilon \to 0} \sup \{ S_n(A)(z), \sigma^n(z) = y, \ d(x,y) \le \varepsilon \}.$$

Then,

 Irreducible components for Ω as subshift of finite type coincide with equivalence classes for

$$S(x,y)+S(y,x)=0.$$



(I) > (A) > (A) > (A) > (A)

Set Ω =Aubry set = subshift of finite type. Set S(x, y)= Mañé potential

$$S(x,y) = \lim_{\varepsilon \to 0} \sup \{ S_n(A)(z), \sigma^n(z) = y, \ d(x,y) \le \varepsilon \}.$$

Then,

 Irreducible components for Ω as subshift of finite type coincide with equivalence classes for

$$S(x,y)+S(y,x)=0.$$



Set Ω =Aubry set = subshift of finite type. Set S(x, y)= Mañé potential

$$S(x,y) = \lim_{\varepsilon \to 0} \sup \{ S_n(A)(z), \sigma^n(z) = y, \ d(x,y) \le \varepsilon \}.$$

Then,

 Irreducible components for Ω as subshift of finite type coincide with equivalence classes for

$$S(x,y)+S(y,x)=0.$$

Set
$$\Omega = \underbrace{\Omega_1 \sqcup \ldots \sqcup \Omega_k}_{\text{entropy } h} \sqcup \underbrace{\Omega_{k+1} \sqcup \ldots \sqcup \Omega_l}_{\text{entropy } < h}$$
Set $a_{ij} := \sup_{x \in \Omega_j} \sup_{y \notin \Omega_i, \ \sigma(y) \in \Omega_i} S(x, y).$

- 4 回 ト 4 回 ト

If V =accumulation point for $\frac{1}{\beta} \log(H_{\beta}) =$ one selected subaction.

$$\gamma \otimes \begin{bmatrix} V(\Omega_1) \\ \vdots \\ V(\Omega_k) \end{bmatrix} = \begin{pmatrix} a_{11} & a_{12} & \dots & a_{1k} \\ \vdots & \vdots & \vdots & \vdots \\ a_{k1} & a_{k2} & \dots & a_{kk} \end{pmatrix} \otimes \begin{bmatrix} V(\Omega_1) \\ \vdots \\ V(\Omega_k) \end{bmatrix}$$

γ is uniquely determined by the "cycles" in the matrix.
 But V is not necessarily unique.

< 回 > < 三 > < 三 >

If V =accumulation point for $\frac{1}{\beta} \log(H_{\beta}) =$ one selected subaction.

$$\gamma \otimes \begin{bmatrix} V(\Omega_1) \\ \vdots \\ V(\Omega_k) \end{bmatrix} = \begin{pmatrix} a_{11} & a_{12} & \dots & a_{1k} \\ \vdots & \vdots & \vdots & \vdots \\ a_{k1} & a_{k2} & \dots & a_{kk} \end{pmatrix} \otimes \begin{bmatrix} V(\Omega_1) \\ \vdots \\ V(\Omega_k) \end{bmatrix}$$

•

< 回 > < 回 > < 回 >

γ is uniquely determined by the "cycles" in the matrix.
 But V is not necessarily unique.

If V =accumulation point for $\frac{1}{\beta} \log(H_{\beta}) =$ one selected subaction.

$$\gamma \otimes \begin{bmatrix} V(\Omega_1) \\ \vdots \\ V(\Omega_k) \end{bmatrix} = \begin{pmatrix} a_{11} & a_{12} & \dots & a_{1k} \\ \vdots & \vdots & \vdots & \vdots \\ a_{k1} & a_{k2} & \dots & a_{kk} \end{pmatrix} \otimes \begin{bmatrix} V(\Omega_1) \\ \vdots \\ V(\Omega_k) \end{bmatrix}$$

.

< 回 > < 回 > < 回 >

γ is uniquely determined by the "cycles" in the matrix.
 But V is not necessarily unique.

If V =accumulation point for $\frac{1}{\beta} \log(H_{\beta}) =$ one selected subaction.

$$\gamma \otimes \begin{bmatrix} V(\Omega_1) \\ \vdots \\ V(\Omega_k) \end{bmatrix} = \begin{pmatrix} a_{11} & a_{12} & \dots & a_{1k} \\ \vdots & \vdots & \vdots & \vdots \\ a_{k1} & a_{k2} & \dots & a_{kk} \end{pmatrix} \otimes \begin{bmatrix} V(\Omega_1) \\ \vdots \\ V(\Omega_k) \end{bmatrix}$$

.

< 回 ト < 三 ト < 三

- γ is uniquely determined by the "cycles" in the matrix.
- 2 But \vec{V} is not necessarily unique.

- O Can we extend this when Ω not of finite type ?
- Can we get convergence for $\frac{1}{\beta} \log H_{\beta}$ to some special eigenvector ?
- What is the role of components with smaller entropy ?
- I Can we detect/decide convergence in function of the phase diagram that gives γ ?

- On the extend this when Ω not of finite type ?
- Can we get convergence for $\frac{1}{\beta} \log H_{\beta}$ to some special eigenvector ?
- What is the role of components with smaller entropy ?
- 3 Can we detect/decide convergence in function of the phase diagram that gives γ ?

- On the extend this when Ω not of finite type ?
- Can we get convergence for $\frac{1}{\beta} \log H_{\beta}$ to some special eigenvector ?
- What is the role of components with smaller entropy ?
- 3 Can we detect/decide convergence in function of the phase diagram that gives γ ?

- On the extend this when Ω not of finite type ?
- 2 Can we get convergence for $\frac{1}{\beta} \log H_{\beta}$ to some special eigenvector ?
- What is the role of components with smaller entropy ?
- 3 Can we detect/decide convergence in function of the phase diagram that gives γ ?

- On the extend this when Ω not of finite type ?
- 2 Can we get convergence for $\frac{1}{\beta} \log H_{\beta}$ to some special eigenvector ?
- What is the role of components with smaller entropy ?
 - Scan we detect/decide convergence in function of the phase diagram that gives γ ?

- On the extend this when Ω not of finite type ?
- 2 Can we get convergence for $\frac{1}{\beta} \log H_{\beta}$ to some special eigenvector ?
- What is the role of components with smaller entropy ?
- 3 Can we detect/decide convergence in function of the phase diagram that gives γ ?

Thank You

Some advances in the selection problem in er

イロン イ理 とく ヨン イヨン

æ