

Stability, cutoff phenomenon and Eyring-Kramers estimates for the Langevin dynamics

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Analysis and simulation of metastable systems, (C.I.R.M.), April 5, 2023

Motivation

- **Molecular Dynamics** methods are used in **Biology, Material Science, Nuclear Physics** (protein folding, nuclear fuels propagation inside the nuclear reactor).

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with $F \in C^2(\mathbb{R}^d, \mathbb{R}^d)$, $\gamma, \epsilon > 0$.

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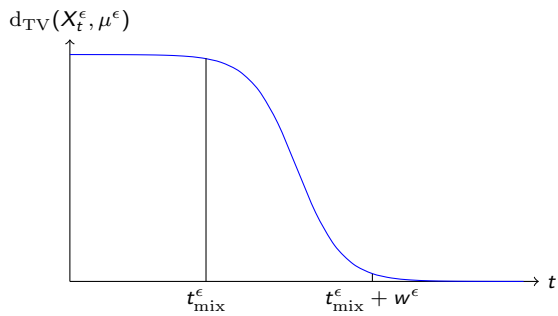
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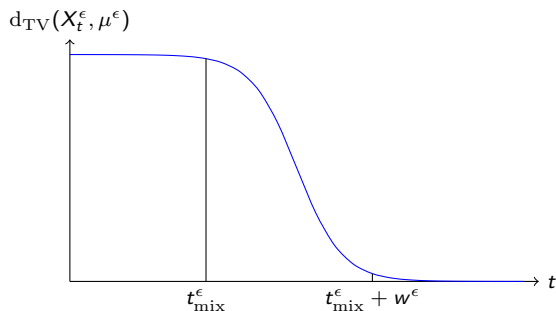
with $F \in C^2(\mathbb{R}^d, \mathbb{R}^d)$, $\gamma, \epsilon > 0$.

Question: How fast does the convergence to the equilibrium distribution μ^ϵ happen?

Cutoff phenomenon



Cutoff phenomenon



Cutoff if $w^\epsilon / t_{mix}^\epsilon \xrightarrow{\epsilon \rightarrow 0} 0$ and

$$\lim_{c \rightarrow +\infty} \limsup_{\epsilon \rightarrow 0} d_{TV}(X_{t_{mix}^\epsilon + cw^\epsilon}^\epsilon, \mu^\epsilon) = 0,$$

$$\lim_{c \rightarrow -\infty} \liminf_{\epsilon \rightarrow 0} d_{TV}(X_{t_{mix}^\epsilon + cw^\epsilon}^\epsilon, \mu^\epsilon) = 1.$$

Overdamped Langevin dynamics case

Let

$$dX_t^\epsilon = -F(X_t^\epsilon)dt + \sqrt{2\epsilon}dB_t, \quad X_0 = x.$$

Theorem (G. Barrera, M. Jara (2020))

Assume that there exist $\alpha > 0$ such that for all $y, q \in \mathbb{R}^d$,

$$DF(y)q \cdot q \geq \alpha|q|^2.$$

Then, there exist $\eta(x), \nu(x), T(x)$ such that cutoff is satisfied with

$$t_{\text{mix}}^\epsilon = \frac{1}{2\eta(x)} \log\left(\frac{1}{2\epsilon}\right) + \frac{\nu(x) - 1}{\eta(x)} \log \log\left(\frac{1}{2\epsilon}\right) + T(x), \quad w^\epsilon = \frac{1}{\eta(x)} + o(1).$$

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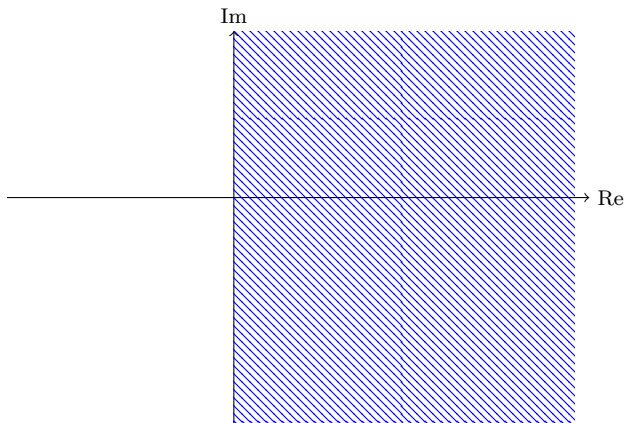
Extension to *Langevin* dynamics with global attractor at 0 ?

$$\begin{cases} dq_t = p_t dt, \\ dp_t = -F(q_t)dt - \gamma p_t dt + \sqrt{2\epsilon}dB_t, \end{cases}$$

Stability of the dynamical system

$$\frac{dX_t}{dt} = -F(X_t), \quad X_0 = x.$$

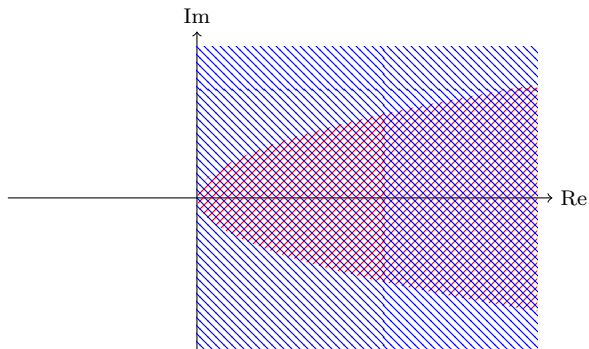
$(X_t)_{t \geq 0}$ locally stable $\iff \exists \alpha > 0$, such that $|x| \leq \alpha \Rightarrow |X_t| \xrightarrow[t \rightarrow \infty]{} 0$
 $\iff \text{Sp}_{\mathbb{C}}(DF(0)) \subset \{a + ib, a > 0, b \in \mathbb{R}\}$.



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 $\iff \text{Sp}_{\mathbb{C}}(DF(0)) \subset \{a + ib, a > 0, b^2 < \gamma^2 a\}$.



Remark: $\forall q \in \mathbb{R}^d, F(q) \cdot q \geq \alpha |q|^2$ not enough to guarantee stability!

Stability condition

$$\begin{cases} dq_t = p_t dt, \\ dp_t = -F(q_t) dt - \gamma p_t dt, \end{cases}$$

Theorem (Global stability)

Assume that

$$F(q) = \nabla U(q) + \ell(q),$$

with U non-negative. Besides, there exists $\alpha > 0$, $\beta \in (0, \gamma)$ such that

$$F(q) \cdot q \geq \alpha(|q|^2 + U(q)) + \frac{|\ell(q)|^2}{\beta^2}.$$

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Then, there exist $C, \lambda > 0$ such that

$$|(q_t, p_t)| \leq C|(q_0, p_0)|e^{-\lambda t}.$$

Linear case

$$\begin{cases} dq_t = p_t dt, \\ dp_t = -Fq_t dt - \gamma p_t dt, \end{cases}$$

where F is a constant normal matrix ($FF^T = F^T F$). Then,

$$\text{Stability} \iff S + \frac{A^2}{\gamma^2} \text{ is definite positive,}$$

where $S = \frac{1}{2}(F + F^T)$, $A = \frac{1}{2}(F - F^T)$.

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Conclusion: Almost sharp condition!

Cutoff in the underdamped setting

Theorem

Assume F satisfies the previous condition then there exist $\eta(x), \nu(x), T(x)$ such that cutoff is satisfied for the underdamped Langevin dynamics starting from $x \in \mathbb{R}^{2d}$ with cutoff time t_{mix}^ϵ and window w^ϵ :

$$t_{\text{mix}}^\epsilon = \frac{1}{2\eta(x)} \log\left(\frac{1}{2\epsilon}\right) + \frac{\nu(x) - 1}{\eta(x)} \log \log\left(\frac{1}{2\epsilon}\right) + T(x), \quad w^\epsilon = \frac{1}{\eta(x)} + o(1).$$

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- Lyapunov matrix equation

$$\begin{pmatrix} 0 & -DF(0)^T \\ I_d & -\gamma I_d \end{pmatrix} \Sigma + \Sigma \begin{pmatrix} 0 & I_d \\ -DF(0) & -\gamma I_d \end{pmatrix} = - \begin{pmatrix} 0 & 0 \\ 0 & I_d \end{pmatrix}.$$

Eyring-Kramers estimates for the Langevin dynamics

$$\begin{cases} dq_t = p_t dt, \\ dp_t = -\nabla U(q_t) dt - \gamma p_t dt + \sqrt{2\gamma\epsilon} dB_t, \end{cases}$$

where $\gamma, \epsilon > 0$.

Assumption: U is a Morse function with suitable growth condition. U admits a local minimum and a global minimum.

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Remark: Guarantees positive recurrence, tightness on the stationary distribution.

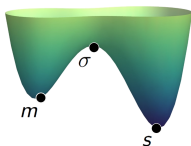


Figure: Double well potential.

Issues:

- No reversibility.
- No ellipticity (hypoelliptic).

Main result

Let μ^ϵ be the stationary distribution of the Langevin process

$$d\mu^\epsilon = \frac{1}{Z_\epsilon} e^{-V(q,p)/\epsilon} dq dp,$$

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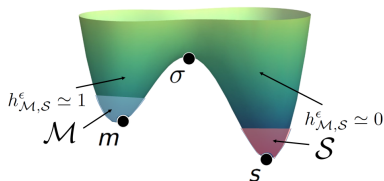
Theorem (Eyring-Kramers estimates)

Let \mathbb{H}^m (resp. \mathbb{H}^σ) be the Hessian matrices of V at m (resp. σ). Then,

$$E_m(\tau_S) = (1 + o_\epsilon(1)) \frac{2\pi}{\lambda^\sigma} \sqrt{\frac{-\det \mathbb{H}^\sigma}{\det \mathbb{H}^m}} \exp\left(\frac{U(\sigma) - U(m)}{\epsilon}\right),$$

where $-\lambda^\sigma$ is the unique negative eigenvalue of $\mathbb{H}^m \begin{pmatrix} 0 & I_d \\ -I_d & \gamma I_d \end{pmatrix}$ and $S = \mathcal{B}(s, \epsilon)$.

Scheme of proof



Potential theory approach.

$$E_\nu(\tau_S) = \frac{1}{C_\epsilon(\mathcal{M}, S)} \int_{\mathbb{R}^{2d}} h_{\mathcal{M}, S}^{\epsilon, *} d\mu^\epsilon.$$

Estimates:

- $E_\nu(\tau_S) = E_m(\tau_S) + o_\epsilon(1).$
- $\int_{\mathbb{R}^{2d}} h_{\mathcal{M}, S}^{\epsilon, *} d\mu^\epsilon(x) = (1 + o_\epsilon(1)) Z_\epsilon^{-1} (2\pi\epsilon)^d e^{-U(m)/\epsilon} \frac{1}{\sqrt{\det \mathbb{H}^m}}.$
- $C_\epsilon(\mathcal{M}_\epsilon, S_\epsilon) = (1 + o_\epsilon(1)) Z_\epsilon^{-1} (2\pi\epsilon)^d e^{-U(\sigma)/\epsilon} \frac{\lambda^\sigma}{2\pi} \frac{1}{\sqrt{-\det \mathbb{H}^\sigma}}.$

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- $\nu_{\text{eq}} := \nabla h_{\mathcal{M}, \mathcal{S}}^{\epsilon, *}(x) \cdot n_{\partial \mathcal{M}}(x) \mu^{\epsilon}(x) \sigma(dx)$ well defined ?
- $\text{Cap}(\mathcal{M}, \mathcal{S}) := \epsilon \int_{(\mathcal{M} \cup \mathcal{S})^c} |\nabla_p h_{\mathcal{M}, \mathcal{S}}^{\epsilon, *}(x)|^2 d\mu^{\epsilon}(x)$ well defined ?

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Remark: Analogous work regarding the spectral gap by Bony, Le Peutrec, Michel (2022)