

# The anisotropic Navier-Stokes system with small unidirectional derivative on the torus

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In progress, joint work with Yanlin Liu (Beijing Normal University)

# The model and functional setting

In  $\mathbb{R}^3$ , we consider the incompressible Navier-Stokes equation

$$\partial_t u + u \cdot \nabla u - \Delta u + \nabla p = 0, \quad \operatorname{div} u = 0, \quad u(0, \cdot) = u_0, \quad \text{where } t \geq 0, x \in \mathbb{R}^3 \quad (\text{NS})$$

and the anisotropic Navier-Stokes equation with only horizontal dissipation

$$\partial_t u + u \cdot \nabla u - \Delta_h u + \nabla p = 0, \quad \operatorname{div} u = 0, \quad u(0, \cdot) = u_0, \quad \text{where } t \geq 0, x \in \mathbb{R}^3 \quad (\text{ANS})$$

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For homogenous tempered distributions  $a, u$ , we define the anisotropic Sobolev space  $H^{s, s'}(\mathbb{R}^3)$  as

$$\|a\|_{H^{s, s'}}^2 := \int_{\mathbb{R}^3} |\xi_h|^{2s} |\xi_3|^{2s'} |\hat{a}(\xi)|^2 d\xi, \quad \text{with } \xi_h = (\xi_1, \xi_2), \quad \xi = (\xi_h, \xi_3)$$

and the anisotropic Besov spaces  $\mathcal{B}, B_{p, q}^{s_1, s_2}$  as

$$\|a\|_{\mathcal{B}} := \sum_{l \in \mathbb{Z}} 2^{\frac{l}{2}} \|\Delta_l^v a\|_{L^2(\mathbb{R}^3)}, \quad \|u\|_{B_{p, q}^{s_1, s_2}} := \left( \sum_{k \in \mathbb{Z}} \sum_{l \in \mathbb{Z}} 2^{qks_1} 2^{qls_2} \|\Delta_k^h \Delta_l^v u\|_{L^p}^q \right)^{1/q}$$

with  $\Delta_k^h, \Delta_l^v$  being the anisotropic dyadic operators

$$\Delta_k^h a = \mathcal{F}^{-1}(\varphi(2^{-k}|\xi_h|)\hat{a}(\xi)), \quad \Delta_l^v a = \mathcal{F}^{-1}(\varphi(2^{-l}|\xi_3|)\hat{a}(\xi))$$

$\varphi(\tau)$  smooth, satisfies  $\operatorname{Supp} \varphi \subset \{3/4 \leq |\tau| \leq 8/3\}$ ,  $\sum_{j \in \mathbb{Z}} \varphi(2^{-j}\tau) = 1$ .

# Global well-posedness of (NS) with small unidirectional derivative

## Theorem

(Y.Liu, P.Zhang, Arch. Ration. Mech. Anal. 2020) Let  $\delta \in (0, 1)$ , and  $u_0 = (u_0^h, u_0^3) \in H^{-\delta, 0} \cap H^2(\mathbb{R}^3)$  with  $\partial_3 u_0^h \in H^{-\delta, 0} \cap H^{-\frac{1}{2}, 0}$ . There exists a universal small positive constant  $\varepsilon_0$  such that if

$$\|\partial_3 u_0\|_{H^{-\frac{1}{2}, 0}}^2 \exp(CA_\delta(u_0^h) + CB_\delta(u_0)) \leq \varepsilon_0$$

where

$$A_\delta(u_0^h) := \left( \frac{\|\nabla_h u_0^h\|_{L_v^\infty(L_h^2)}^2 \|u_0^h\|_{L_v^\infty(B_{2, \infty}^{-\delta})_h}^{2/\delta}}{\|u_0^h\|_{L_v^\infty(L_h^2)}^{2/\delta}} + \|u_0^h\|_{L_v^\infty(L_h^2)}^2 \right) \cdot \exp(C_\delta(1 + \|u_0^h\|_{L_v^\infty(L_h^2)}^4)),$$

$$B_\delta(u_0) := \|u_0\|_{H^{-\delta, 0}}^{\frac{1}{2}} \|u_0\|_{H^{\delta, 0}}^{\frac{1}{2}} \|\partial_3 u_0\|_{H^{-\delta, 0}}^{\frac{1}{2}} \|\partial_3 u_0\|_{H^{\delta, 0}}^{\frac{1}{2}} \exp(CA_\delta(u_0^h))$$

then (NS) has a unique global solution  $u \in C(\mathbb{R}^+, H^{\frac{1}{2}}) \cap L^2(\mathbb{R}^+, H^{\frac{3}{2}})$ .

# Global well-posedness of (ANS) with small unidirectional derivative

## Theorem

(M.Paicu, P.Zhang, *Comm. Math. Phys.* 2011) There exist universal constants  $L > 0$  and a sufficiently small constant  $\varepsilon_0 > 0$  such that

$$\|u_0^h\|_{\mathcal{B}} \exp(L\|u_0^3\|_{\mathcal{B}}^4) \leq \varepsilon$$

then (ANS) has a unique global solution  $u \in C([0, \infty), \mathcal{B})$  with  $\nabla_h u \in L^2([0, \infty), \mathcal{B})$ .

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## Theorem

(Y.Liu, M.Paicu and P.Zhang, *Arch. Ration. Mech. Anal.* 2020) Let  $u_0, \partial_3 u_0$  satisfy  $u_0, \partial_3 u_0 \in L^2 \cap H^{-1,1}$ , and there exist universal constants  $N_1, N_2 > 0$  and a sufficiently small constant  $\varepsilon_0 > 0$  such that

$$\begin{aligned} & \|\partial_3 u_0\|_{H^{-1,0}} \|\partial_3^2 u_0\|_{H^{-1,0}} \\ & \times \exp(N_1(1 + \|u_0^3\|_{L^2}^2 \|\partial_3 u_0^3\|_{L^2}^2)) \exp(\exp(N_2\|u_0^h\|_{L^2}^2 \|\partial_3 u_0^h\|_{L^2}^2)) \leq \varepsilon_0 \end{aligned}$$

then (ANS) has a unique global solution  $u \in C([0, \infty), \mathcal{B})$  with  $\nabla_h u \in L^2([0, \infty), \mathcal{B})$ .

# The anisotropic Navier-Stokes equation on the torus

We want to give a more delicate description for the smallness of the derivative on vertical variable, and we start from the torus  $\mathbb{T}_h^2 \times \mathbb{R}_v$ . For the solution  $u$ , we do the Fourier expansion

$$u = \sum_{k \in \mathbb{Z}^2} \hat{u}_k e^{ikx_h}, \quad u_k = \bar{u}_{-k}$$

then we have

$$\begin{cases} \partial_t u_0 = - \begin{pmatrix} 0 \\ \partial_3 p_0 \end{pmatrix} - \partial_3 \left( \sum_{m \neq 0} u_m^3 u_{-m} \right) \\ \partial_t u_k + k^2 u_k = - \begin{pmatrix} ik \\ \partial_3 \end{pmatrix} p_k - ik_j \sum_{m+n=k} u_m^j u_n - \partial_3 \left( \sum_{m+n=k} u_m^3 u_n \right), \quad k \neq 0 \end{cases} \quad (\text{ANST})$$

# Main result

## Theorem

There exist universal constants  $N_1, N_2 > 0$  and a constant  $\varepsilon_0 > 0$  small enough, such that if the initial data  $u_{k,int}$  satisfies

$$\begin{aligned} & \left( \|\partial_3 \hat{u}_{0,int}\|_{L^2} + \sum_{k \neq 0} \frac{1}{|k|} \|\partial_3 \hat{u}_{k,int}\|_{L^2} \right) \left( \|\partial_3^2 \hat{u}_{0,int}\|_{L^2} + \sum_{k \neq 0} \frac{1}{|k|} \|\partial_3^2 \hat{u}_{k,int}\|_{L^2} \right) \\ & \times \exp((N_1 + N_1 \|u_{0,int}^3\|_{L^2}^2 \|\partial_3 u_{0,int}^3\|_{L^2}^2 + N_1 \sum_{k \neq 0} \|\partial_3 u_{k,int}^3\|_{L^2}^4) \\ & \times \exp(\exp(N_2 \|u_{0,int}^h\|_{L^2}^2 \|\partial_3 u_{0,int}^h\|_{L^2}^2 + N_2 \sum_{k \neq 0} \|\partial_3 u_{k,int}^h\|_{L^2}^4))) \leq \varepsilon_0 \end{aligned}$$

then (ANST) has a unique global solution.



# Strategies of the proof

We approximate the solution by

$$\partial_t \bar{u}^h + \bar{u}^h \cdot \nabla_h \bar{u}^h - \Delta_h \bar{u}^h = -\nabla_h \bar{p}, \quad \operatorname{div}_h \bar{u}^h = 0$$

We introduce Helmholtz decomposition for the initial data. For any vector field  $u^h$  in dimension 2, we decompose  $u^h$  as

$$u^h = \nabla_h^\perp \Delta_h^{-1}(\operatorname{curl}_h u^h) + \nabla_h \Delta_h^{-1}(\operatorname{div}_h u^h) := u_{\operatorname{curl}}^h + u_{\operatorname{div}}^h$$

where  $\operatorname{div}_h u_{\operatorname{curl}}^h = 0$ . Thus we can choose  $\bar{u}_{t=0}^h := \nabla_h^\perp \Delta_h^{-1}(\operatorname{curl}_h u_0^h)$  and the reminder term  $v = u - \bar{u}$  satisfies

$$\begin{cases} \partial_t v^h + v \cdot \nabla v^h + \bar{u}^h \cdot \nabla_h v^h + v \cdot \nabla \bar{u}^h - \Delta_h v^h = -\nabla_h p + \nabla_h \bar{p}, \\ \partial_t v^3 + v \cdot \nabla v^3 + \bar{u}^h \cdot \nabla_h v^3 - \Delta_h v^3 = -\partial_3 p, \\ \operatorname{div} v = 0, v|_{t=0} = (-\nabla_h \Delta_h^{-1}(\partial_3 u_0^3), u_0^3). \end{cases}$$

# Strategies of the proof

Let the Fourier multiplier  $\Lambda_h := (-\Delta_h)^{\frac{1}{2}}$  in (ANS). The key step is to give the following energy estimates

$$\|\bar{u}^h(t)\|_{L^2}^2 + 2\|\nabla_h \bar{u}^h(t)\|_{L_t^2(L^2)}^2 = \|\bar{u}_0^h\|_{L^2}^2$$

and

$$\|\Lambda_h^{-1} \partial_3^\alpha \bar{u}^h(t)\|_{L^2}^2 + \|\partial_3^\alpha \bar{u}^h\|_{L_t^2(L^2)}^2 \leq \|\Lambda_h^{-1} \partial_3^\alpha \bar{u}_0^h\|_{L^2}^2 \underbrace{\exp(\dots \exp(C\|\bar{u}_0^h\|_{L^2}^2 \|\partial_3 \bar{u}_0^h\|_{L^2}^2))}_{\alpha \text{ times}}$$

for  $\alpha = 1, 2$ .

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for  $\alpha = 1, 2$ . For the equation of remainder  $v$ , we study the time weighted Chemin-Lerner space as follows. For some  $\lambda > 0$ , define

$$v_\lambda(t) := v(t) \exp\left(-\lambda \int_0^t f(t') dt'\right)$$

where  $f$  needs to be chosen differently for  $v^h$  and  $v^3$ . And we multiply the equation of  $v^h, v^3$  by  $\exp(-\lambda \int_0^t f(t') dt')$  to give the estimates in some Hilbert space  $\mathcal{H}$ .

# Strategies of the proof

Finally, we prove the main theorem by a continuity argument. For the maximal solution of (ANST) with lifespan  $T^*$ , consider the following continuity argument:

$$\|v^h\|_{\tilde{L}_t^\infty(\mathcal{H})} + \|v^h\|_{\tilde{L}_{t,f}^2(\mathcal{H})} + \|\nabla_h v^h\|_{\tilde{L}_t^2(\mathcal{H})} \leq A(u_0), \quad (\text{A1})$$

$$\|v^3\|_{\tilde{L}_t^\infty(\mathcal{H})} + \|\nabla_h v^3\|_{\tilde{L}_t^2(\mathcal{H})} \leq B(u_0), \quad (\text{A2})$$

here  $B(u_0) \ll A(u_0)$ . Denote

$$T^{**} := \sup\{T \in (0, T^*) : (\text{A1}), (\text{A2}) \text{ hold for any } t \in [0, T]\}$$

we can prove that if  $A(u_0)$  is small enough, then  $T^{**} = T^*$ . In view of the blow-up criterion, we have  $T^* = \infty$ . ( see M. Paicu, Equation anisotrope de Navier-Stokes dans des espaces critiques, Rev. Mat. Iberoamericana 2005) And the global well-posedness is proved.

## Future work

- A natural question is that for general  $\partial_X u_0$  small, where  $X$  is a variable direction (which can depend on time and space), could we obtain global well-posedness result?
- What could we say about the anisotropic Navier-Stokes equation with only vertical dissipation? As far as we have known, there are many open problems about it.

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Thanks for your attention!