# Energy balance for 2D incompressible fluid flow 

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## Collaborators:

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Smooth inviscid flows $(\nu=0)$ conserve kinetic energy

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Brief history...

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Involves studying optimal conditions for energy flux to vanish.


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(B) $\operatorname{div}\left[u^{\varepsilon}\left(\frac{\left|u^{\varepsilon}\right|^{2}}{2}+p^{\varepsilon}\right)\right] \rightarrow \operatorname{div}\left[u\left(\frac{|u|^{2}}{2}+p\right)\right]$ in the sense of distributions;
(C) $u^{\varepsilon} \cdot \mathcal{R}^{\varepsilon} \rightarrow 0$ strongly in $L^{\infty}\left(0, T ; L^{1}\left(\mathbb{T}^{2}\right)\right)$.
(A) and (B) are subcritical for $\omega \in L^{3 / 2}$. In fact, they require $\omega \in L^{6 / 5}$. It is the convergence of the energy flux term, which is (C), that requires $\omega \in L^{3 / 2}$. (Good behavior of the energy flux term is the key point in all results along these lines.)

The proof of $(\mathrm{C})$ uses convergence of mollifications together with the Sobolev imbedding: $\omega \in L^{3 / 2} \Longrightarrow u^{\varepsilon}$ bounded in $L^{\infty}\left(0, T ; L^{6}\left(\mathbb{T}^{2}\right)\right)$.

Key fact: $u \cdot(u \cdot \nabla) u \in L^{1} ;\|u \cdot[(u \cdot \nabla) u]\| \lesssim\|u\|_{W^{1, \frac{3}{2}}}^{3}$.
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Note $S_{q}$ is a convolution with a mollifier.

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The family $\left\{u^{\nu}\right\}$ is called a physical realization of $u$.

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Obs. $1<p<3 / 2$ 'Onsager supercritical'.

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Write $y=y(t)=\left\|\omega^{\nu}\right\|_{L^{2}}^{2}$ and $C_{0}=\left\|\omega_{0}^{\nu}\right\|_{L^{p}}^{-\frac{2 p}{2^{2-p}}}$.

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\lim _{\nu \rightarrow 0}\left\|u^{\nu}(t, \cdot)\right\|_{L^{2}}^{2}-\left\|u_{0}^{\nu}\right\|_{L^{2}}^{2}=0
$$

DiPerna-Majda 1987, $\omega \in L^{p}, p>1$, non-concentration result:

$$
\lim _{\nu \rightarrow 0}\left\|u^{\nu}(t, \cdot)\right\|_{L^{2}}^{2}=\|u(t, \cdot)\|_{L^{2}}^{2}
$$

Strong convergence of initial data, hypothesis, not compactness:

$$
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The proof is concluded.

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## Proposition

Under the hypotheses of the Theorem,

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\lim _{\nu \rightarrow 0^{+}} \nu \int_{0}^{t}\left\|\omega^{\nu}(s, \cdot)\right\|_{L^{2}}^{2} d s \rightarrow 0
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Step 1 Fix $\delta>0$. Let $m^{\nu}$ solution of

$$
\left\{\begin{array}{l}
m^{\prime}=-A \nu m^{\alpha}+B \sqrt{m} \\
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Step 1 Fix $\delta>0$. Let $m^{\nu}$ solution of

$$
\left\{\begin{array}{l}
m^{\prime}=-A \nu m^{\alpha}+B \sqrt{m} \\
m(\delta)=z^{\nu}(\delta)
\end{array}\right.
$$

Then $0 \leq z^{\nu}(t) \leq m^{\nu}(t)$ all $t \in(\delta, T)$.

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Step 3:

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\nu \int_{0}^{t} z^{\nu}(s) \mathrm{d} s \leq \nu \int_{R_{\nu}^{* *}}^{z^{\nu}(0)} \Phi_{\nu}(y) \mathrm{d} y+\nu t R_{\nu}^{* *}+\nu R_{\nu}^{* *} \Phi_{\nu}\left(z^{\nu}(0)\right)
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This is enough to conclude the proof of the Proposition.

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- Lanthaler et al equivalence criterion with forcing? Less regular forcing?
- Extension to approximations by vortex blob method, $L^{\infty}\left(W_{l o c}^{1, p}\right)$, $p>1$ and local energy balance $p \geq 6 / 5$. Ciampa, Crippa, Spirito 2020.
- Extension to axisymmetric Euler : Nobili Seis 2022. Initial vorticity $\omega_{0}$ nonnegative, $|x| \omega_{0}(\cdot)$ integrable, $\omega_{0} / r \in L^{p}(r \mathrm{~d} r \mathrm{~d} z), p>3 / 2$.
- Energy conservation in the case $p=1$ ? No tools. There is a discrepancy wrt conservation of $L^{p}$-norms! Less ambitious: $p=1$, $u$ physically realizable, can $u$ be attainable by convex integration? Work in progress.
- Lanthaler et al equivalence criterion with forcing? Less regular forcing? Also work in progress.


## Thank you!

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Merci!

