# Energy balance for 2D incompressible fluid flow

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Smooth inviscid flows ( $\nu = 0$ ) conserve kinetic energy

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## Wild solutions, anomalous dissipation

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Involves studying optimal conditions for energy flux to vanish.

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- wild solutions: no control on integrability of vorticity

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#### Definition

Fix T > 0 and  $u_0 \in L^2(\mathbb{T}^2)$  with initial vorticity  $\omega_0 = \operatorname{curl} u_0 \in L^p(\mathbb{T}^2)$ , for some  $p \ge 1$ .
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Fix T > 0 and let  $u \in C(0, T; L^2_{weak}(\mathbb{T}^2))$  be a weak solution with  $\omega \equiv \text{ curl } u \in L^{\infty}(0, T; L^{3/2}(\mathbb{T}^2)).$ 

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Helena J. Nussenzveig Lopes (IM-UFRJ) Energy balance 2D incompressible flow

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There exists a divergence free vector field  $u \in B_{3,\infty}^{1/3} \cap W^{1,p}(\mathbb{T}^2)$ , for any  $1 \le p < 3/2$ , such that  $\limsup_{q \to \infty} \prod_q [u] \ne 0$ ,

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Suggests exists dynamical mechanism preventing anomalous dissipation in 2D even for supercritical (less than 1/3 regular) flows

#### Definition

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•  $u^{\nu}(0,\cdot) \equiv u_0^{\nu} \rightarrow u_0$  strongly in  $L^2(\mathbb{T}^2)$ .

The family  $\{u^{\nu}\}$  is called a *physical realization* of *u*.





#### Theorem (Cheskidov,Lopes Filho, N-L, Shvydkoy; 2016)

Let  $u \in C(0, T; L^2(\mathbb{T}^2))$  be a physically realizable weak solution of the incompressible 2D Euler equations.

### Energy

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Obs. 1 'Onsager supercritical'.

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Proof: Assume  $\omega_0 \in L^p(\mathbb{T}^2)$  for some p < 2, and  $\omega_0 \notin L^2(\mathbb{T}^2)$  otherwise, the result is trivial. u is physically realizable  $\Longrightarrow \exists$  physical realization  $\{u^{\nu}\}$  solutions of Navier-Stokes with  $\{\omega_0^{\nu}\}$  bounded in  $L^p$ .

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$$-2\nu \|\nabla \omega^{\nu}\|_{L^{2}}^{2} \leq -2\nu \|\omega^{\nu}\|_{L^{2}}^{\frac{4}{2-\rho}} \|\omega^{\nu}\|_{L^{p}}^{-\frac{2\rho}{2-\rho}}.$$

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Multiply the vorticity equation by  $|\omega^{\nu}|^{p-2}\omega^{\nu}$  and integrate on torus  $\Longrightarrow$ 

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## Energy identity for 2D Navier-Stokes:
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(4)

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The proof is concluded.

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- Lanthaler *et alli* analysis relies on *L*<sup>2</sup>-based *structure function* for *u*; play the role of vorticity estimates.

# Forced fluid flow and energy balance

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$$u_t + u \cdot \nabla u = -\nabla p + \nu \Delta u + F,$$
  
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Physically realizable weak solutions:

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Then u is energy balanced.

Proof: Suppose  $\omega_0 \notin L^2$ .

$$\partial_t \omega^{\nu} + u^{\nu} \cdot \nabla \omega^{\nu} = \nu \Delta \omega^{\nu} + g^{\nu}.$$

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$$\frac{d}{dt} \|\omega^{\nu}\|_{L^{2}}^{2} \leq -C(\|\omega_{0}\|_{L^{p}}, \|\boldsymbol{g}^{\nu}\|_{L^{1}(L^{p})}) \nu \|\omega^{\nu}\|_{L^{2}}^{\frac{4}{2-p}} + \|\boldsymbol{g}^{\nu}\|_{L^{\infty}(L^{2})} \|\omega^{\nu}\|_{L^{2}}.$$

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#### Proposition

Under the hypotheses of the Theorem,

$$\lim_{\nu\to 0^+}\nu\int_0^t\|\omega^\nu(\boldsymbol{s},\cdot)\|_{L^2}^2\,d\boldsymbol{s}\to 0.$$

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$$\frac{dz^{\nu}}{dt} \leq -A\nu(z^{\nu})^{\frac{2}{2-p}} + B(z^{\nu})^{\frac{1}{2}},$$

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Then  $0 \le z^{\nu}(t) \le m^{\nu}(t)$  all  $t \in (\delta, T)$ .

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$$\limsup_{\nu \to 0^+} \frac{z^{\nu}(0)}{R_{\nu}^*} = \begin{cases} <1 \\ =1 \\ >1. \end{cases}$$

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$$\|
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$$egin{aligned} & 
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$$\begin{split} \nu \int_0^t \|\omega^{\nu}(\boldsymbol{s},\cdot)\|_{L^2}^2 \, d\boldsymbol{s} &\leq \nu \int_0^t \boldsymbol{R}_{\nu}^* \, d\boldsymbol{s} \\ &\leq \nu \int_0^t \left(\frac{\boldsymbol{B}}{\boldsymbol{A}\nu}\right)^{\frac{2}{2\alpha-1}} \, d\boldsymbol{s} \equiv \boldsymbol{C}\nu^{1-\frac{2}{2\alpha-1}} \\ &= \boldsymbol{C}\nu^{\frac{2\alpha-3}{2\alpha-1}} \to 0 \text{ as } \nu \to 0. \end{split}$$

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This is enough to conclude the proof of the Proposition.

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