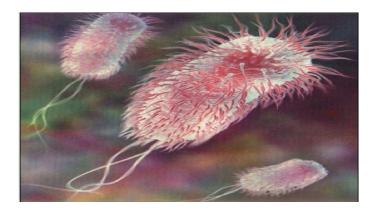
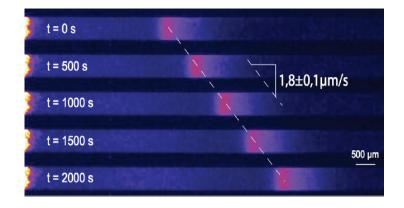


Bacterial movement by run and tumble : models, patterns, pathways, scales

Benoît Perthame





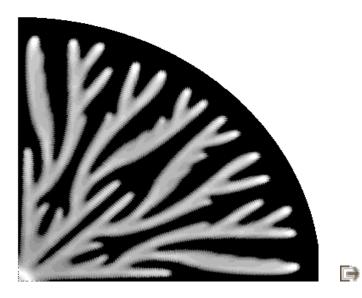
Motivation



Paradigm for collective organisation



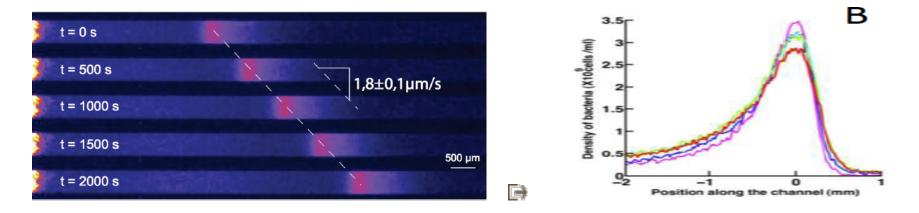
Courtesy S. Seror, B. Holland (Paris-Sud),



Numerical simulation of a mathematical model

Motivation

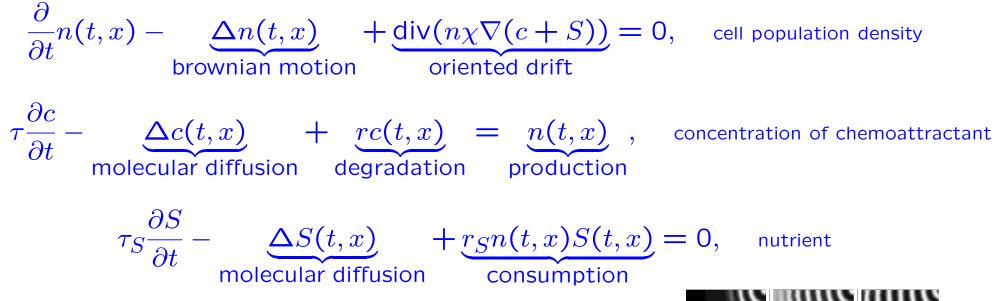




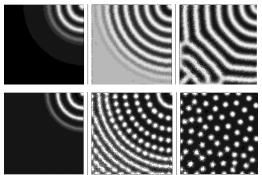
- Adler's famous experiment for *E. Coli* (1966)
- Self-attraction + attraction towards fresh nutrient
- Explain this pattern; its asymmetry (experiments Curie institute)
- *E. coli* is a chemotactic bacterium
- Several strains are used; the phenomena is robust
- Fluid dynamics is not dominant

Chemotaxis : Keller-Segel





- Internal mathematical interest
- Finite time blow-up
- Cannot sustain robust traveling bands
- In opposition with kinetic/hyperbolic models



Chemotaxis : Flux Limited K.-S.



The Flux Limited Keller-Segel (Dolak-Schmeiser, Erban-Othmer)

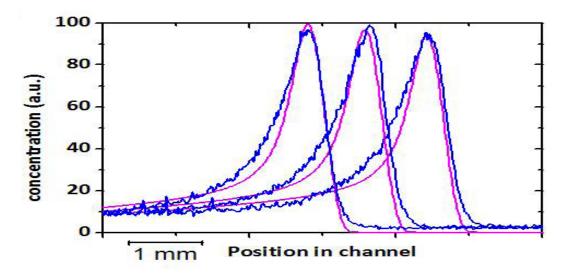
$$\begin{aligned} \int \frac{\partial}{\partial t} n(t,x) - \Delta n(t,x) + \operatorname{div}(nU) &= 0, \\ U &= \chi_c(c_t,c_x) \frac{\nabla c}{|\nabla c|} - \chi_S(S_t,S_x) \frac{\nabla S}{|\nabla S|} \\ \frac{\partial}{\partial t} c - D_c \Delta c &= n(t,x) \\ \int \frac{\partial}{\partial t} S - D_S \Delta S &= -n(x,t) S(t,x) \end{aligned}$$

admits traveling band solutions

Can be fit to the experimental data

Chemotaxis : Flux Limited K.-S.





Superimposition of the calculated (pink) and the experimental (blue) concentration profiles at three different times.

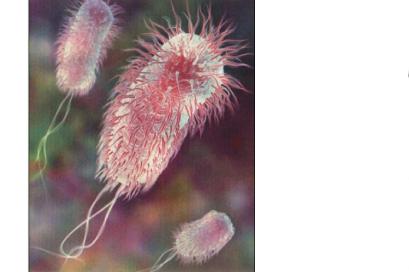
• Where does the FLKS comes from?

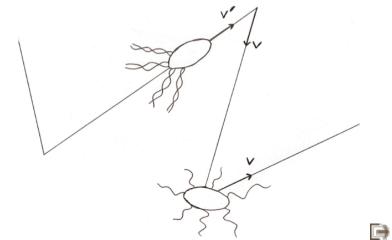
The kinetic formalism uses rules for individual behaviour

Kinetic models (80-90's)



E. Coli is known to move by run and tumble Alt, Dunbar, Othmer, Stevens, Hillen...





A beautiful example of multiscale motion

Kinetic models (80-90's)

■ $f(t, x, \xi)$ the population density of cells moving with the velocity ξ ■ c(t, x) the chemoattractant concentration

 $\frac{\partial}{\partial t}f(t,x,\xi) + \underbrace{\xi \cdot \nabla_x f}_{\text{run}} = \underbrace{\mathcal{K}[c,S]f}_{\text{tumble}},$

$$\mathcal{K}[c,S]f = \int_B K(c,S;\xi,\xi')f(\xi')d\xi' - \int_B K(c,S;\xi',\xi)d\xi' f,$$

Various forms of the tumbling kernel K[c, S] have been proposed
Typical K(c; \(\xi, \xi') = k\) (c(x - \(\epsilon \xi')\)) (with chemoattractant only)



Kinetic models : asymptotic



Based on the run time ε : $K(c; \xi, \xi') = k(c(x - \varepsilon \xi'))$

$$\frac{\partial}{\partial t}f_{\varepsilon}(t,x,\xi) + \frac{\xi \cdot \nabla_x f_{\varepsilon}}{\varepsilon} = \frac{\mathcal{K}[c]}{\varepsilon^2},$$

Before blow-up time

$$f_{\varepsilon}(t, x, \xi) \to n(t, x), \qquad c_{\varepsilon}(t, x) \to c(t, x),$$

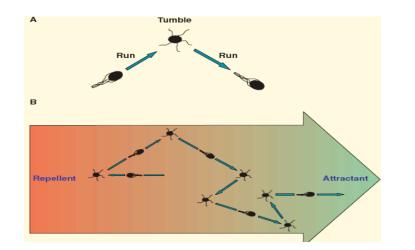
 $rac{\partial}{\partial t}n(t, x) - \operatorname{div}[D \nabla n(t, x)] + \operatorname{div}(n\chi \nabla c) = 0,$

and the transport coefficients are given by

$$D(c) = \frac{D_0}{k(c)}, \qquad \chi(c) = \chi_0 \frac{k'(c)}{k(c)}.$$

Pulse waves





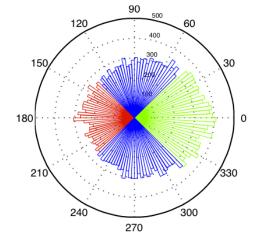
When *c* increases, jumps are longer

$$\frac{\partial}{\partial t}f(t,x,\xi) + \xi \cdot \nabla_x f = \mathcal{K}[c]f$$
$$K(c;\xi,\xi') = \mathbf{K}_{\varepsilon} \left(\frac{\partial c}{\partial t} + \xi' \cdot \nabla c\right)$$

■ Macroscopic limit is the Flux Limited K.-S. system

Pulse waves





angular distribution and mean run time

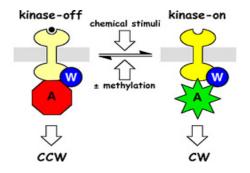
$$\frac{\partial}{\partial t}f(t, x, \xi) + \xi \cdot \nabla_x f = \mathcal{K}[c]f$$
$$K(c; \xi, \xi') = \mathbf{K}_{\varepsilon} \left(\frac{\partial c}{\partial t} + \xi' \cdot \nabla c\right)$$

Can one explain the tumbling rate

Biochemical pathways



Can one explain the tumbling rate $\mathbf{K}_{\varepsilon} \left(\frac{\partial c}{\partial t} + \xi' \cdot \nabla c \right)$?



Use the internal biochemical pathway controling tumbling (Erban-Othmer, Dolak-Schmeiser),

 $g(t, x, \xi, m)$ receptor methylation level (internal state).

Biochemical pathways



Principle : internal state m adapts to the external state

 $m \approx M(c)$

$$\frac{\partial}{\partial t}g_{\varepsilon}(t,x,\xi,m) + \xi \cdot \nabla_{x}g_{\varepsilon} + \frac{1}{\varepsilon}\frac{\partial}{\partial m}[R(m-M(c))g_{\varepsilon}] = \mathcal{K}_{\varepsilon}[m,c][g_{\varepsilon}]$$

 $\mathcal{K}_{\varepsilon}[m,c][g_{\varepsilon}] = \int \left[K(\frac{m-M(c)}{\varepsilon},\xi,\xi')g_{\varepsilon}(x,\xi',m) - K(...,\xi',\xi)g_{\varepsilon}(t,x,\xi,m) \right] d\xi'$

Fast adaptation, stiff response

Theorem The limit $\varepsilon \to 0$ gives $g_{\varepsilon} \to \delta(m - M(c))f(s, \xi, t)$ and $\mathbf{K}\left(\frac{\partial c}{\partial t} + \xi' \cdot \nabla c\right)$

Abnormal diffusions



ARTICLE Received 28 Mar 2015 | Accepted 19 Aug 2015 | Published 25 Sep 2015 Dol: 10.1038/ncomme9396 OPEN Swarming bacteria migrate by Lévy Walk

Gil Ariel¹, Amit Rabani², Sivan Benisty², Jonathan D. Partridge³, Rasika M. Harshey³ & Avraham Be'er²

$$\varepsilon^{1+\mu}\frac{\partial}{\partial t}g(t,x,\xi,m) + \varepsilon\xi \cdot \nabla_x g + \varepsilon^s \Delta_m g = \mathcal{K}[m,c][g]$$

When $\mathcal{K}[m,c][g]$ degenerates,

 $\mathcal{K}[m,c][g]\approx 0 \quad \text{as} \ m\rightarrow\infty,$

the limiting behaviour is fractional Laplacian

$$\frac{\partial n}{\partial t} - \Delta^{\alpha} n = 0$$



Conclusion



Flux-Limited Keller-Segel system relies on a multiscale approach (molecule to cell to population)

It is possible to fit quantitatively the experimental data

Numerous mathematical questions (singularities, asymptotic, fractional derivatives, waves...)

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