Phase Space Sampling for

Trajectory Simulations

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- in reactants / products,
- "energy leaking" below ground quantum level.

...not taken into account by <u>classical</u> trajectories.

2

may change reaction rate by orders of magnitude!

Quantum effects in reactive collisions

- Tunnel effects:
 - through barrier,
 - corner cutting.
- Zero-point energy:

 - all along trajectories,



Quantum effects in trajectory simulations

- Can we account for quantum effects—at least in part—by means of *trajectories*?
- Here we will look specifically at the *sampling* of trajectory initial conditions.

First, need to get back to 1D quantum trajectories



Quantum trajectory approach the "no-wave" formulation

- Originates in (but quite different from) deBroglie-Bohm.
- The *trajectory ensemble itself* is the fundamental quantum state entity, rather than the wave function Ψ .

- Bill Poirier, Bohmian mechanics without pilot waves, Chem. Phys. 370, 4 (2010).
- J. Schiff and Bill Poirier, Quantum mechanics without wavefunctions, JCP 136, 031102 (2012).
- G. Parlant, Y.-C. Ou, K. Park, and B. Poirier, Classical-like trajectory simulations for accurate computation of quantum reactive scattering probabilities. *Comput. Theor. Chem.*, **990**, 3 (2012).

Stationary scattering state - 1D (#1)

Hamilton eqs of motion –

(x, p) "real space"

$$\dot{x} = \frac{s}{m}$$
$$\dot{p} = -\frac{\partial V}{\partial x}$$

(x, p) "real space"

(r, s) "quantum space"

$$\dot{x} = \frac{s}{m}$$
$$\dot{p} = -\frac{\partial V}{\partial x}$$

$$\dot{r} = \frac{p-s}{m} - \frac{8r^2s^3}{m\hbar^2}$$
$$\dot{s} = \frac{4rs^4}{m\hbar^2}$$

(x, p) "real space"

(r, s) "quantum space"

. 8	∂H
$x = -\frac{1}{m}$	$= \overline{\partial p}$
∂V	∂H
$p = -\frac{1}{\partial x}$	$=-\frac{1}{\partial x}$

$$\dot{r} = \frac{p - s}{m} - \frac{8r^2 s^3}{m\hbar^2} = \frac{\partial H}{\partial s}$$
$$\dot{s} = \frac{4rs^4}{m\hbar^2} = -\frac{\partial H}{\partial r}$$

(x,p) "real space"

(r, s) "quantum space"

. 8	∂H	$p-s 8r^2s^3$	∂H
x = -m	$=\overline{\partial p}$	$r = -\frac{m}{m} - \frac{m\hbar^2}{m\hbar^2}$	$=$ $\overline{\partial s}$
∂V	∂H	$. 4rs^4$	∂H
$p = -\frac{1}{\partial x}$	$=-\overline{\partial x}$	$s = \overline{m\hbar^2}$	$=-\overline{\partial r}$

Hamiltonian:

$$H(\boldsymbol{x},\boldsymbol{p},\boldsymbol{r},\boldsymbol{s}) = \frac{\boldsymbol{s}(2\boldsymbol{p}-\boldsymbol{s})}{2m} + V(\boldsymbol{x}) - \frac{2\boldsymbol{r}^2\boldsymbol{s}^4}{\underline{m}\hbar^2}$$

H(x, p, r, s) reduces to classical Hamiltonian when r = 0 and s = p

(x,p) "real space"

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. 8	∂H	$p-s = 8r^2s^3$	∂H
x = -m	$=\overline{\partial p}$	r =	$= \overline{\partial s}$
∂V	∂H	$4rs^4$	∂H
$p = -\frac{\partial x}{\partial x}$	$=-\overline{\partial x}$	$s = \overline{m\hbar^2}$	$=-\overline{\partial r}$

Hamiltonian:

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H(x, p, r, s) reduces to classical Hamiltonian when r = 0 and s = p

four initial conditions:

1 for time translation / 1 for total energy / 2 to specify particular quantum state $\Psi(x)$



(x,p) "real space"

(r,s) "quantum space"

· _ S	$- \partial H$	$\dot{r} - p - s = 8r^2s^3$	∂H
$x = \frac{1}{m}$	$-\overline{\partial p}$	$m = \frac{1}{m} = \frac{1}{m\hbar^2}$	$=\overline{\partial s}$
. ∂V	∂H	$. 4rs^4$	∂H
$p = -\frac{1}{\partial x}$	$=-\overline{\partial x}$	$s = \overline{m\hbar^2}$	$=-\overline{\partial r}$

Hamiltonian:

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Stationary scattering state - 1D (#2)

The classical limit!...

Stationary scattering state - 1D (#2)

The classical limit!...



Tunneling through a 1D Eckart barrier (#1)



$$\Psi(x \to +\infty) = k_R^{-1/2} \mathbf{T} \exp[ik_R x]$$
$$\Psi(x \to -\infty) = k_L^{-1/2} \left(\exp[ik_L x] + \mathbf{R} \exp[-ik_L x]\right)$$

Initial conditions: $x_0 = x_R$ $\dot{x}_0 = \sqrt{2E/m}$ $\ddot{x}_0 = \dddot{x}_0 = 0$

Transmission proba: $P_T = |T|^2$ Reflection proba: $P_R = |R|^2$

Tunneling through a 1D Eckart barrier (#2)

-solution quantum trajectory-



 $\alpha = 3.0 \, a.u.$ $m = 2000 \, a.u.$ $V_0 = 400 \, cm^{-1}$

Eckart potential:

 $V(x) = V_0 \operatorname{sech}^2(\alpha x)$

 ODE integration by means of a simple (adaptive step-size) Runge-Kutta propagator.

• typical CPU time = 0.03 s

extremely deep tunneling

E/V ₀		5.48 (-43)	
Exact F	• Transmission	1.5641888074199	039 (- 43)
Computed F	Transmission	1.5641888074199	345 (-43)
Relative Erro	or	~ 2.0 (-14)	

One single trajectory gives the exact solution

Extension to multi-D systems assume tunneling is most important along reaction coordinate 'x'

- *x* → follows quantum motion along reaction path,
 - exact tunneling / interferences.
- $y \rightarrow$ classical vibration,
 - will need to choose the (undetermined vibrational phase *at random*.

X

Write down quantum and classical motion eqs





A + BC

2D Hamilton eqs of motion

notice the tunnel / interference extra space!

$$\dot{x}=rac{\partial H}{\partial p_{x}} \quad \dot{p_{x}}=-rac{\partial H}{\partial x};$$

$$\dot{r}=rac{\partial H}{\partial p_r} ~~ \dot{p_r}=-rac{\partial H}{\partial r};$$

quantum

tunnel / interference

$$\dot{y} = \frac{\partial H}{\partial p_y} \quad \dot{p_y} = -\frac{\partial H}{\partial y}$$

$$\textbf{x} \Rightarrow \textbf{y} \text{ coupling}$$

$$\textbf{classical}$$

$$H = \frac{p_r(2p_x - p_r)}{2m} - \frac{2r^2p_r^4}{m\hbar^2} + \frac{p_y^2}{2m} + V(x, y)$$

Hamiltonian

prototype 2D chemical system

2D bottleneck potential -







Trajectory ensemble sampling (#1) "Traditional" classical sampling

- Takes care of tunnel effects but does not account for *quantization* of reactant/product states.
- *State-resolved 'v' reaction probabilities* obtained by "binning" trajectories into energy regions centered on individual quantum states 'v' of the classical y-oscillator.
- Solution? Restrict asymp. channel phase space to EBK quantization:
 - Problem: lack of micro-reversibility,
 - Problem: computationally inefficient.

To fix this we propose the Phase Space Approximation



v=1

v=0

Trajectory ensemble sampling (#2) Phase Space Approximation (PSA)

Bill Poirier, Algebraically Self-Consistent Quasiclassical Approximation on Phase Space, Found. Phys. 30, 1191 (2000).

- Derives from exact Wigner-Weyl formalism (Bill Poirier).
- Vibrational quantum state \rightarrow classical phase space region of area $2\pi\hbar$.
- Main difference with traditional sampling: for a given quantum state v all trajectories have *the same kinetic energy*
 - \implies energy conservation is not enforced.

Provides exact reactive scattering for separable systems



trajectory ensemble sampling (#3)

...more on PSA sampling -

- Individual 1D quantum states for the y-oscillator are represented by phase-space regions of area $2\pi\hbar$
- All trajectories have different energies: $E_{\text{traj}} \neq E_{\text{total}}$



PSA vs traditional sampling

Standard sampling *E*_{tot} *is conserved* PSA sampling *E*_{tot} *is not conserved*



cumulative reaction probabilities - CRPs

— x and y motions decoupled (b=0) —

cumulative reaction probabilities - CRPs

x and y motions decoupled (b=0)









cumulative reaction probabilities - CRPs





cumulative reaction probabilities - CRPs narrower bottleneck (b=1)



cumulative reaction probabilities - CRPs wider bottleneck (b= -0.25)





- New trajectory sampling procedure.
- Correct quantum state thresholds.
- Easy to use.
- Would be interesting to investigate for fully classical systems.
- Try to use it all along trajectories to avoid "energy leak"?