

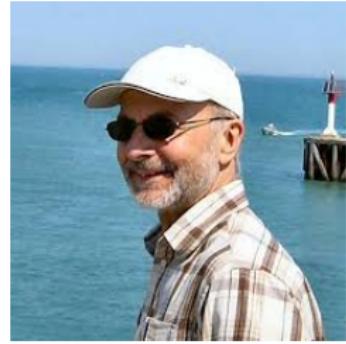
Propagation of Wave Packets and Herman-Kluk Propagators

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Université Paris Est - Créteil & CNRS

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Introduction

Schedule of the talk.

- ① **Quantum Dynamics.** Systems of Schrödinger time-evolution equations.
- ② **Scalar case.** How to trade solving high dimensional oscillating equations into solving ODEs ? Egorov theorem, thawed or frozen gaussians approximation, “non-hopping algorithm”.
- ③ **Systems.** Zoology of systems, revisiting the scalar methods, the need of “hopping” algorithms.

Quantum Dynamics

Schrödinger equations

- Born-Oppenheimer Approximation

$$i\varepsilon \partial_t \psi^\varepsilon = \hat{H} \psi^\varepsilon, \quad \psi^\varepsilon(0, x) = \psi_0^\varepsilon(x),$$

$$H = -\frac{\varepsilon^2}{2} |\xi|^2 + V(x),$$

V matrix-valued, self-adjoint, with subquadratic growth,

$$z = (x, \xi) \in \mathbb{R}^{2d}, \quad (t, x) \in \mathbb{R} \times \mathbb{R}^d, \quad \psi^\varepsilon \in L^2(\mathbb{R}^d, \mathbb{C}^N).$$

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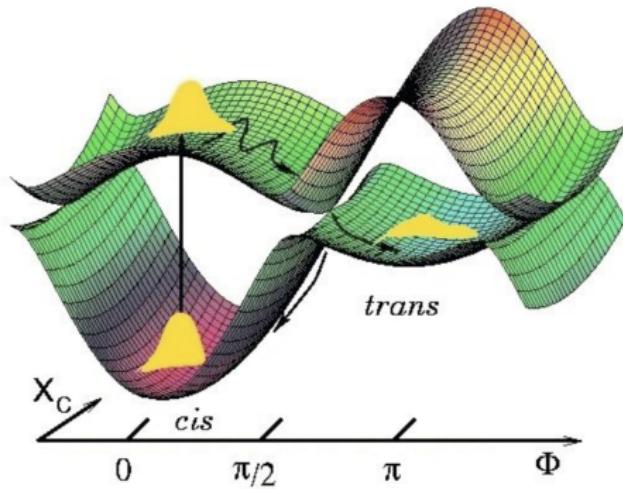
$$z = (x, \xi) \in \mathbb{R}^{2d}, \quad (t, x) \in \mathbb{R} \times \mathbb{R}^d, \quad \psi^\varepsilon \in L^2(\mathbb{R}^d, \mathbb{C}^N).$$

⇒ The aim: Describe

- ① $\psi^\varepsilon(t) = e^{-\frac{i}{\varepsilon} t \hat{H}} \psi_0^\varepsilon$
- ② $(\hat{a} \psi^\varepsilon(t), \psi^\varepsilon(t))_{L^2(\mathbb{R}^d, \mathbb{C}^N)}, \quad a \in \mathcal{C}_0^\infty(\mathbb{R}^{2d}, \mathbb{C}^{N,N})$.

What is known about H ? and ε

- Energy surfaces - Pyrazine



cf. Köppel and Yarkony

- Who is ε ?

For molecular dynamics, $\varepsilon = \sqrt{\frac{m_e}{m_N}}$.

Algorithms from the 70s

Underlying ideas:

- **Key point:** solving ODEs instead of PDEs by developing ansatz based on classical quantities.
- **Key tools:** Propagation of wave packets + Gaussian frame.

Algorithms from the 70s

Underlying ideas:

- Key point: solving ODEs instead of PDEs by developing ansatz based on classical quantities.
- Key tools: Propagation of wave packets + Gaussian frame.

A large range of methods:

- Frozen gaussians & Thawed gaussians approximations,
from chemistry [Heller 1981](#), [Herman-Kluk 1984](#), [Kay 1994](#), [Domcke & all 2019](#)
to mathematics [Rousse-Swart 2009](#), [Robert 2010](#), [Lasser-Sattleger 2017](#).
- Surface hopping semi-groups,
[Tully Preston 1971](#), [FK Lasser 2012](#), [Kube Lasser Weber 2009](#), [Lu 2018](#).
- Variational methods, MCTDH and G-MCTDH,
[Worth, Lasorne, Burghardt & Römer 2013](#), [Lubich 2015](#).

Scalar equations

$H = h \text{Id}$ for h real-valued.

Egorov theorem

- Egorov Theorem: for any observable a

$$(\widehat{a} \psi^\varepsilon(t), \psi^\varepsilon(t))_{L^2(\mathbb{R}^d, \mathbb{C})} = \left(\widehat{a \circ \Phi_h^t} \psi^\varepsilon(t), \psi^\varepsilon(t) \right)_{L^2(\mathbb{R}^d, \mathbb{C})} + O(\varepsilon^2)$$

where $t \mapsto \Phi_h^t(z_0)$ is the classical trajectory

$$\dot{\Phi}_h^t(z_0) = Jdh(\Phi_h^t(z_0)), \quad \Phi_h^0(z_0) = z_0.$$

$$h(x, \xi) = \frac{|\xi|^2}{2} + V(x), \quad J = \begin{pmatrix} 0 & \text{Id} \\ -\text{Id} & 0 \end{pmatrix}.$$

- Numerical realization

- ① Sample the initial Wigner transform $\implies (z_j, w_j)_{1 \leq j \leq N}$
- ② Propagate the weighted points:

$$(\Phi_h^t(z_j), w_j), \quad 1 \leq j \leq N.$$

- ③ Sum-up at finite time:

$$(\widehat{a} \psi^\varepsilon(t), \psi^\varepsilon(t))_{L^2(\mathbb{R}^d, \mathbb{C})} \sim \sum_{j=1}^N a(\Phi_h^t(z_j)) w_j.$$

Bargmann formula

- **Gaussian:** Let Γ in the positive half Siegel space,

$$g_0^\Gamma(x) = c_\Gamma e^{i\Gamma x \cdot x}, \quad \|g_0^\Gamma\|_{L^2(\mathbb{R}^d)} = 1, \quad {}^t\Gamma = \Gamma, \quad \text{Im } \Gamma > 0.$$

- **Gaussian wave packets:** Let $z = (q, p) \in \mathbb{R}^{2d}$,

$$g_z^{\Gamma, \varepsilon}(x) = \varepsilon^{-\frac{d}{4}} g_0^\Gamma \left(\frac{x - q}{\sqrt{\varepsilon}} \right) e^{\frac{i}{\varepsilon} p \cdot (x - q)}$$

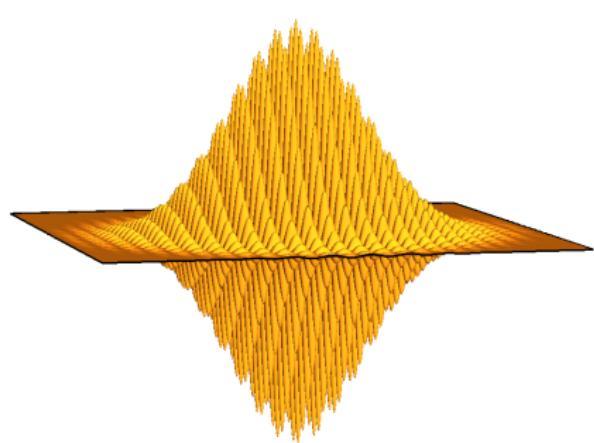
- **Bargmann transform:** Decomposition of $\psi \in L^2(\mathbb{R}^d)$ on Gaussians

$$\psi(x) = (2\pi\varepsilon)^{-d} \int_{z \in \mathbb{R}^{2d}} \langle g_z^\varepsilon, \psi \rangle g_z^\varepsilon(x) dz, \quad \psi \in L^2(\mathbb{R}^d)$$

where $g_z^{j\text{Id}, \varepsilon} = g_z^\varepsilon$.

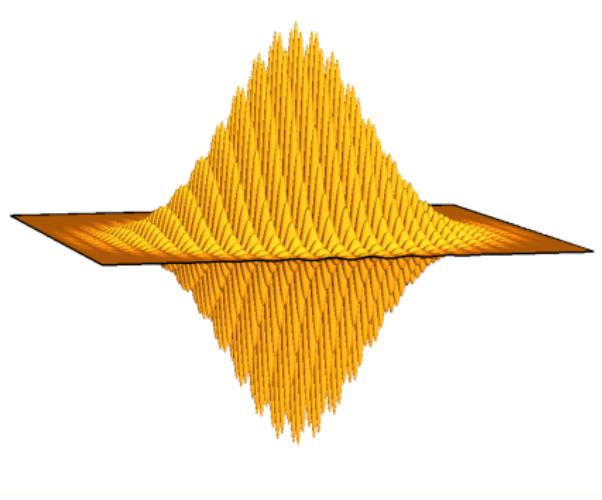
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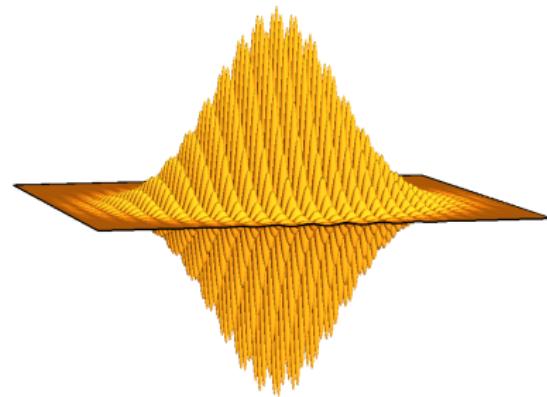
Bargmann formula

$$\psi_0^\varepsilon(x) = (2\pi\varepsilon)^{-d} \int_{z \in \mathbb{R}^{2d}} \langle g_z^\varepsilon, \psi_0^\varepsilon \rangle g_z^\varepsilon(x) dz, \quad \psi \in L^2(\mathbb{R}^d)$$



Bargmann formula

$$e^{-\frac{i}{\varepsilon}t\widehat{H}^\varepsilon}\psi_0^\varepsilon(x) = (2\pi\varepsilon)^{-d} \int_{z \in \mathbb{R}^{2d}} \langle g_z^\varepsilon, \psi_0^\varepsilon \rangle e^{-\frac{i}{\varepsilon}t\widehat{H}^\varepsilon} g_z^\varepsilon(x) dz.$$



Wave packet propagation:

$$e^{-\frac{i}{\varepsilon}t\widehat{H}^\varepsilon}g_z^\varepsilon(x) = e^{\frac{i}{\varepsilon}S(t,z)} g_{\Phi_H^t(z)}^{\Gamma(t,z), \varepsilon}(x) + O(\sqrt{\varepsilon})$$

Thawed Gaussian propagator for scalar equations

Theorem (Kay 2006, Rousse & Swart 2009, Robert 2010)

$$e^{-\frac{i}{\varepsilon}t\widehat{H}}\psi_0^\varepsilon = (2\pi\varepsilon)^{-d} \int_{\mathbb{R}^{2d}} \langle g_z^\varepsilon, \psi_0^\varepsilon \rangle e^{\frac{i}{\varepsilon}S(t,z)} g_{\Phi_h^t(z)}^{\Gamma(t,z), \varepsilon} dz + O(\varepsilon) \text{ in } L^2(\mathbb{R}^d),$$

Classical quantities

- $\Phi_h^t(z) = (q(t), p(t))$ is the classical trajectory arising from z .
- $S(t, z)$ is the classical action, $\dot{S} = p \cdot \dot{q} - h(q, p)$, $S(0) = 0$.
- $\Gamma(t, z) = (C(t, z) + iD(t, z))(A(t, z) + iB(t, z))^{-1}$
where $F(t, z) = \begin{pmatrix} A(t, z) & B(t, z) \\ C(t, z) & D(t, z) \end{pmatrix}$ satisfies

$$\partial_t F = J \operatorname{Hess}_z h(t, \Phi_h^t(z)) F(t), \quad F(0) = \operatorname{Id}_{\mathbb{R}^{2d}}.$$

Frozen Gaussian propagator for scalar equations

Theorem (Kay 2006, Rousse & Swart 2009, Robert 2010)

$$e^{-\frac{i}{\varepsilon}t\widehat{H}}\psi_0^\varepsilon = (2\pi\varepsilon)^{-d} \int_{\mathbb{R}^{2d}} \langle g_z^\varepsilon, \psi_0^\varepsilon \rangle u(t, z) e^{\frac{i}{\varepsilon}S(t, z)} g_{\Phi_h^t(z)}^\varepsilon dz + O(\varepsilon) \text{ in } L^2(\mathbb{R}^d),$$

$$\text{with } u(t, z) = 2^{-d/2} \det^{1/2} (A(t, z) + D(t, z) + i(C(t, z) - B(t, z))).$$

Key points of the proof [Robert 2010]:

- ① Wave packet propagation

$$e^{-\frac{i}{\varepsilon}t\widehat{H}^\varepsilon} g_z^\varepsilon(x) = e^{\frac{i}{\varepsilon}S(t, z)} g_{\Phi_H^t(z)}^{\Gamma(t, z), \varepsilon}(x) + O(\sqrt{\varepsilon})$$

- ② Treat the remainder in $O(\sqrt{\varepsilon})$ in order to obtain $O(\varepsilon)$ (integration by parts).
- ③ Turn the width $\Gamma(t, z)$ of the Gaussian into the prefactor $u(t, z)$ by an evolution argument drawing a path from $\Gamma_0 = i\text{Id}$ to $\Gamma_1 = \Gamma(t, 0, z)$.

Frozen Gaussian propagator for scalar equations

Algorithmic realization [Lasser & Sattlegger 2017]

- ① Initial Sampling of the data by a Monte Carlo procedure

$$\psi_0^\varepsilon(x) \sim (2\pi\varepsilon)^{-d} \sum_{1 \leq j \leq N} r_0^\varepsilon(z_j) g_{z_j}^\varepsilon$$

with z_1, \dots, z_N i.i.d. according to the probability measure

$$\mu_0^\varepsilon(dz) = \left(\int |\langle g_z, \psi_0^\varepsilon \rangle| dz \right)^{-1} |\langle g_z, \psi_0^\varepsilon \rangle| dz.$$

- ② Transport of the sample points + classical quantities

$$z_j(t) = \Phi_h^t(z_j), \quad S(t, z_j), \quad u(t, z_j), \quad 1 \leq j \leq N.$$

- ③ Quadrature formula

$$\psi^\varepsilon(t, x) \sim (2\pi\varepsilon)^{-d} \sum_{1 \leq j \leq N} r_0^\varepsilon(z_j) e^{\frac{i}{\varepsilon} S(t, z_j)} u(t, z_j) g_{\Phi_h^t z_j}^\varepsilon.$$

What about systems ?

Which systems ? which data ?

- The data is chosen on an energy surface:

$$\psi_0^\varepsilon = \widehat{\vec{V}_0} \phi_0^\varepsilon$$

with $\phi_0^\varepsilon \in L^2(\mathbb{R}^d, \mathbb{C})$ and $H(z) \vec{V}_0(z) = h(z) \vec{V}_0(z)$.

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- How is h ?

① **Gapped eigenvalue**: there exists $\delta_0 > 0$ such that

$$d(h(z), \sigma(H(z)) \setminus \{h(z)\}) > \delta_0.$$

② **Crossings**: let $h =: h_1$ and h_2 , eigenvalues of H ,

$$\Upsilon := \{z \in \mathbb{R}^{2d}, \quad h_1(z) = h_2(z)\} \neq \emptyset.$$

Classification according to the dimension of Υ ([Hagedorn 1994](#)), its geometric properties ([Colin de Verdière 2003](#), [FK & Gérard 2002](#))
⇒ normal forms and transition formulas.

Time-dependent eigenvectors

Adiabatic decoupling: If $\psi_0^\varepsilon = \widehat{\vec{V}_0} \phi_0^\varepsilon + O(\varepsilon)$ with $H\vec{V}_0 = h\vec{V}_0$, then

$$\psi^\varepsilon(t) = \widehat{\vec{V}(t)} e^{-\frac{i}{\varepsilon} t \widehat{h}} \phi_0^\varepsilon + O(\varepsilon)$$

where $H\vec{V} = h\vec{V}$ is given by parallel transport:

$$\partial_t \vec{V} + \xi \cdot \nabla_x \vec{V} - \nabla h \cdot \vec{V} = \Omega \vec{V}, \quad \vec{V}(0) = \vec{V}_0.$$

$$\Omega(x, \xi) = (\text{Id} - \Pi(x)) \xi \cdot \nabla \Pi(x).$$

⇒

- Egorov theorem for matrix-valued symbols.
- Thawed/Frozen Gaussian approximation of the propagator.

Gapped systems - Thawed Gaussian propagator

- Propagation of the Gaussian WP:

$$e^{-i\frac{t}{\varepsilon}\hat{H}}(\vec{V}_0 g_z^\varepsilon) = e^{\frac{i}{\varepsilon}S(t,z)} g_{\Phi_h^t(z)}^{\varepsilon,\Gamma(t)} \vec{V}(t) + O(\sqrt{\varepsilon})..$$

Proposition (FLR 2019)

$$e^{-i\frac{t}{\varepsilon}\hat{H}}\psi_0^\varepsilon = (2\pi\varepsilon)^{-d} \int_{\mathbb{R}^{2d}} \langle g_z^\varepsilon, v_0^\varepsilon \rangle u(t, z) \vec{V}(t, \Phi_h^t(z)) e^{\frac{i}{\varepsilon}S(t,z)} g_{\Phi_h^t(z)}^\varepsilon dz + O(\varepsilon).$$

- A vector-valued prefactor:

$$\vec{U}(t, z) = u(t, z) \vec{V}(t, \Phi_h^t(z)).$$

What about systems with crossings ?

Systems with crossings - surface hoppings

Key ideas:

- Far from the crossing set: adiabatic regime.
- Close to the crossing set: generate new trajectories when the trajectories reach a **hopping hypersurface**

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Good news:

- For “generic” conical crossings, the time-dependent eigenvector exists up to the crossing.

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Two questions:

- Which hopping hypersurface ?
 \Rightarrow minimal gap condition (**FK Lasser** 2012 - expectation values).
- Which weight on the new trajectories ?
 \Rightarrow Transition coefficients arising from the analysis of the normal forms model problems (starting from **Landau and Zener** 1930's).

Systems with codimension 1 crossings

- $N \geq 2$, $H = h_1 \Pi_1 + h_2 \Pi_2$.
- The functions h_1 , h_2 , Π_1 and Π_2 are smooth

$$h_1 = v + f, \quad h_2 = v - f, \quad v = \frac{1}{2} \operatorname{Tr} H, \quad f \text{ gap.}$$

- Classical quantities associated with h_j , $j \in \{1, 2\}$

$$\Phi_j^{t,t_0}(z), \quad S_j(t, t_0, z), \quad F_j(t, t_0, z), \quad \Gamma_j(t, t_0, z), \quad u_j(t, t_0, z), \quad \vec{U}_j(t, t_0, z), \dots$$

Systems with codimension 1 crossings

- Generic codimension 1 crossing:

$$\mu(z) := \frac{1}{2}(\partial_t f + \{v, f\}(z)) \neq 0, \quad \forall z \in \Upsilon.$$

⇒ The trajectories are transverse to the **crossing hypersurface**

$$\Upsilon = \{f(q) = 0\}.$$

- If $z \in \mathbb{R}^d$, $t^\flat(z)$ is the crossing time

$$\Phi_1^{t^\flat(z), 0}(z) = z^\flat(z) \in \Upsilon.$$

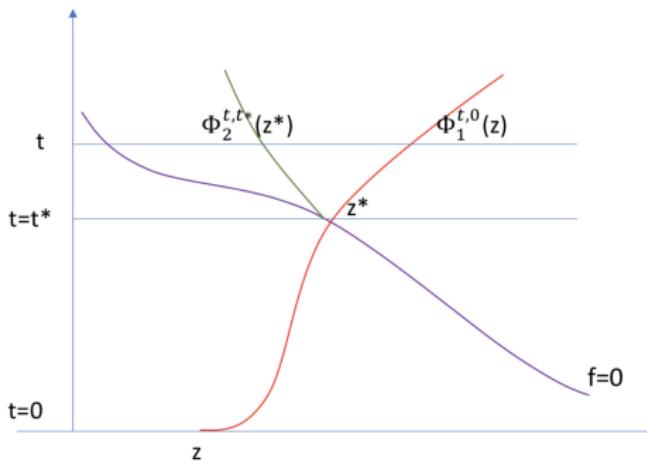
- With $\vec{V}_1 = \Pi_1 \vec{V}_1$, we associate $\vec{V}_1(t, 0, z)$ and $\vec{V}_2(t, t^\flat, z)$ with

$$\vec{V}_2(t^\flat, t^\flat, z) = \vec{V}_1(t^\flat, 0, z)^\perp.$$

- We associate with $z \in \mathbb{R}^{2d}$

$$S^\flat(z) = S_1(t^\flat, 0, z), \quad \alpha^\flat(z) = \|(\partial_t \Pi_1 + \{v, \Pi_1\}) \vec{V}_1(t^\flat, z^\flat)\|_{\mathbb{C}^N}.$$

Systems with codimension 1 crossings



Propagation of a Wave Packet

Propagation of Gaussian wave packets

Theorem (Hagedorn 94, Watson & Weinstein 18 , CFK Lasser Robert 19)

Assume $\psi_0^\varepsilon = \widehat{V}_1 g_{z_0}^\varepsilon$, then in $L^2(\mathbb{R}^d)$

$$e^{-\frac{i}{\varepsilon}t\widehat{H}}\psi_0^\varepsilon = \widehat{\vec{V}}_1(t)v_1^\varepsilon(t) + \sqrt{\varepsilon}\mathbf{1}_{t>t^\flat}\widehat{\vec{V}}_2(t)v_2^\varepsilon(t) + O(\varepsilon^{2/3})$$

- $v_1^\varepsilon(t)$ solves $i\varepsilon\partial_t v_1^\varepsilon = \widehat{h}_1 v_1^\varepsilon$, $v_1^\varepsilon(0) = g_{z_0}^\varepsilon$.
- $v_2^\varepsilon(t)$ solves $i\varepsilon\partial_t v_2^\varepsilon = \widehat{h}_2 v_2^\varepsilon$, $v_2^\varepsilon(t^\flat) = \alpha^\flat e^{iS^\flat/\varepsilon} g_{z^\flat}^{\Gamma^\flat, \varepsilon}$,
where α^\flat , S^\flat , Γ^\flat are classical quantities associated with the crossing point z^\flat .

The result extends to non gaussian wave packets data:

$$\varepsilon^{-\frac{d}{4}} a \left(\frac{x - q}{\sqrt{\varepsilon}} \right) e^{\frac{i}{\varepsilon} p \cdot (x - q)}, \quad a \in \mathcal{S}(\mathbb{R}^d).$$

Frozen Gaussian app. for codim 1 crossings

Theorem (CFK. Lasser Robert, *still writing...*)

Assume $\psi_0^\varepsilon = \widehat{\vec{V}}_1 v_0^\varepsilon$. then

$$\begin{aligned} \psi^\varepsilon(t, x) &= (2\pi\varepsilon)^{-d} \int_{\mathbb{R}^{2d}} e^{\frac{i}{\varepsilon} S_1(t, 0, z)} \vec{U}_1^\varepsilon(t, 0, z) \langle g_z^\varepsilon, v_0^\varepsilon \rangle g_{\Phi_1^{t, 0}(z)}^\varepsilon(x) dz \\ &+ \sqrt{\varepsilon} (2\pi\varepsilon)^{-d} \int_{\mathbb{R}^{2d}} \mathbf{1}_{t^\flat(z) < t} \alpha^\flat(z) e^{\frac{i}{\varepsilon} S^\flat(z)} e^{\frac{i}{\varepsilon} S_2(t, t^\flat, z^\flat)} \vec{U}_2^\varepsilon(t, t^\flat, z^\flat) \\ &\quad \times \langle g_z^\varepsilon, v_0^\varepsilon \rangle g_{\Phi_2^{t, t^\flat}(z^\flat)}^\varepsilon(x) dz + o(\sqrt{\varepsilon}) \end{aligned}$$

with vector-valued Herman-Kluk prefactors

$$\vec{U}_1^\varepsilon(t, 0, z) = \vec{V}_1 \left(t, 0, \Phi_1^{t, 0}(z) \right) u_1(t, 0, z),$$

$$\vec{U}_2^\varepsilon(t, t^\flat, z^\flat) = \vec{V}_2 \left(t, t^\flat, \Phi_2^{t, t^\flat}(z^\flat)(z^\flat) \right) u_2(t, t^\flat, z^\flat).$$

Frozen Gaussian solver for codimension 1 crossings

Assume $\psi_0^\varepsilon = \widehat{V}_1 v_0^\varepsilon$.

- ① Initial sampling of v_0^ε : $\Rightarrow N$ weighted sample points

$$(z_1, r_1), \dots, (z_N, r_N).$$

- ② Transport along Φ_1^t : $\Rightarrow z_1^{(1)}(t), \dots, z_N^{(1)}(t)$,
the actions $S_1(t, 0, z_j)$ and prefactors $\vec{U}_1(t, 0, z_j)$.

- ③ Branching process: If $f(z_j(t))$ changes of sign at time t_j^\flat .

Generate $z_j^{(2)}(t)$ on the mode h_2 with starting point $z_j^\flat := z_j^{(1)}(t_j^\flat)$.
Set $r_j^\flat := \alpha(z_j^\flat) e^{\frac{i}{\varepsilon} S_1(t_j^\flat, 0, z_j)} r_j$.

- ④ Conclusion:

$$\begin{aligned} \psi^\varepsilon(t, x) \sim & (2\pi\varepsilon)^{-d} \sum_{1 \leq j \leq N} r_j e^{\frac{i}{\varepsilon} S_1(t, 0, z_j)} \vec{U}_1(t, 0, z_j) g_{z_j^{(1)}(t)}^\varepsilon \\ & + \sqrt{\varepsilon} \sum_{t_j^\flat < t} r_j^\flat e^{\frac{i}{\varepsilon} S_2(t, t_j^\flat, z_j^\flat)} \vec{U}_2(t, t_j^\flat, z_j^{(2)}(t)) g_{z_j^{(2)}(t)}^\varepsilon \end{aligned}$$

Conclusion

- Better estimates for expectation value than for wave function .
- For systems, parallel transport + hopping trajectories with deterministic branching process.
- Codimension 2 and 3 crossing (including Dirac points):
 - ① single switch surface hopping for observables (transitions at leading order)
 - ② thawed & frozen gaussian approximations (work in progress with Lysianne Hari and Stephanie Gamble).