#### The Causal Structure of Quantum Information

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Inria Grenoble - Rhône-Alpes

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#### **Quantum Foundations**

**Quantum foundations**: the study of the conceptual and mathematical underpinnings of quantum theory

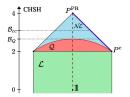


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Quantum entanglement







Quantum measurement



Quantum causality

#### Quantum correlations

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This lecture: Quantum causality



- Understand causal structure of quantum theory
- Fundamentally quantum causal structures?
- Exploit this for quantum information processing?

#### Outline

#### Causal structure of quantum information

Defining causal structure Causality and quantum circuits Quantum combs

#### Quantum control of causal structure

The Quantum Switch The Quantum *N*-Switch Application: Fourier Promise Problem

#### A general model of circuits with quantum control of causal structure Quantum circuits with quantum control

Other directions of study in quantum causality

What does it mean to say that X causes Y?

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- Causal (Bayesian) models: Framework to describe causal relations from observed correlations
  - J. Pearl, *Causality* (2000).

How to define causal relations between quantum events?

What do we take to be quantum events?

Classically: Events: random variables; Interventions: stochastic maps

How to define causal relations between quantum events?

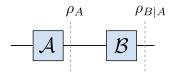
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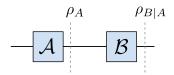
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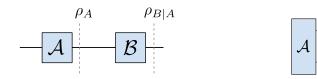


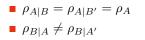
 $\rho_{A|B} = \rho_{A|B'} = \rho_A$  $\rho_{B|A} \neq \rho_{B|A'}$ 

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 $\rho_{BC|A} = \mathcal{B} \otimes \mathcal{C}(\rho_A)$ 

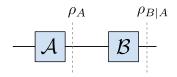
 $\rho_A$ 

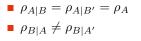
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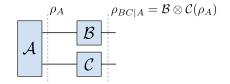
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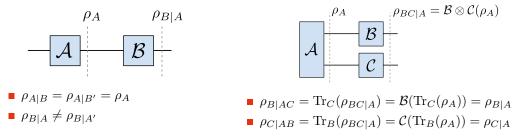


$$\rho_{B|AC} = \operatorname{Tr}_C(\rho_{BC|A}) = \mathcal{B}(\operatorname{Tr}_C(\rho_A)) = \rho_{B|A}$$
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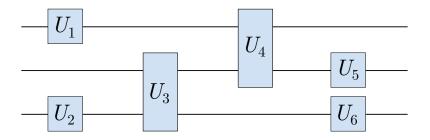
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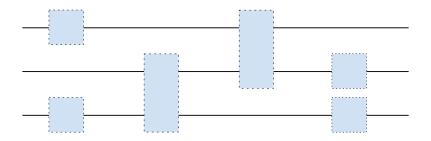


Causal structure defined by ability to influence or "signal" from one operation to another Quantum causal models [Barrett, Lorenz, Oreshkov, 2019]

Quantum circuits have a causal structure

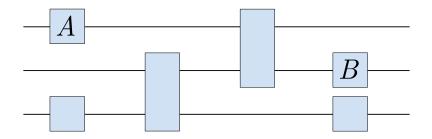


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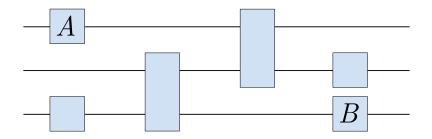
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Need to consider circuits with modifiable operations

• A circuit location A can influence another B if there is there is a path from A to B

## **Defining the Scenario**

Consider causal structure in a computational scenario

- Black-box "operations" (quantum channels)
  - $\mathcal{A}: \mathcal{L}(\mathcal{H}^{A^{I}}) \to \mathcal{L}(\mathcal{H}^{A^{O}})$

$$\rho - \mathcal{A} - \mathcal{A}(\rho)$$

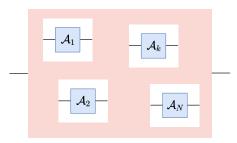
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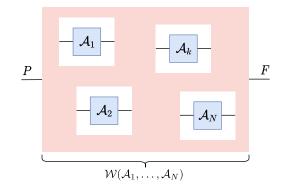
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- Consume N queries  $\mathcal{A}_1, \ldots, \mathcal{A}_N$  in some "computation"
  - May have all  $A_i \equiv A$  (N queries to A), but they could also be different operations



# Higher Order Operations (Quantum Supermaps)



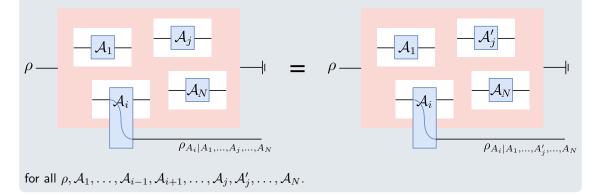
 $\text{Higher order operation: } (\mathcal{A}_1, \dots, \mathcal{A}_N) \mapsto \mathcal{W}(\mathcal{A}_1, \dots, \mathcal{A}_N) : \mathcal{L}(\mathcal{H}^P) \to \mathcal{L}(\mathcal{H}^F)$ 

- Multilinear in its arguments
- $\mathcal{W}(\mathcal{A}_1,\ldots,\mathcal{A}_N)$  a quantum channel whenever  $\mathcal{A}_1,\ldots,\mathcal{A}_N$  are quantum channels

## Formalising Causal Order

#### Compatibility with causal order

 $\mathcal{W}$  is compatible with  $A_1 \prec A_2 \prec \cdots \prec A_N$  if, for all i < j,  $A_j$  cannot signal to  $A_i$ :



Note:  $\ensuremath{\mathcal{W}}$  can be consistent with several causal orders

## Quantum Combs

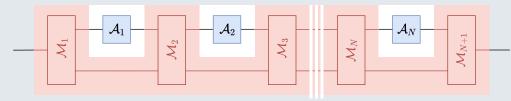
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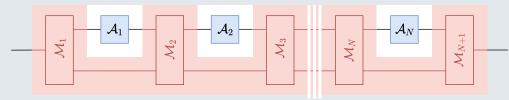
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Alternatively: process operator, channel with memory, process tensor, ...

Quantum circuits are the most general computation with a fixed causal structure.

# Formalising Quantum Combs

Higher order operations can be nicely formulated in the Choi picture

#### Choi-Jamiołkowski isomorphism

 $\mathsf{CP} \text{ maps } \mathcal{C}: \mathcal{L}(\mathcal{H}^X) \to \mathcal{L}(\mathcal{H}^Y) \text{ are in a bijection with PSD operators } C \in \mathcal{L}(\mathcal{H}^X \otimes \mathcal{H}^Y)$ 

$$C = \mathcal{I} \otimes \mathcal{C}(|\mathbb{1}\rangle\!\rangle\!\langle\!\langle \mathbb{1}|), \quad \text{where } |\mathbb{1}\rangle\!\rangle = \sum_{i} |i\rangle \otimes |i\rangle$$

• TP condition: 
$$\operatorname{Tr}_Y[C] = \mathbb{1}^X$$

Inverse: 
$$\mathcal{C}(\rho) = \operatorname{Tr}_X[(\rho^T \otimes \mathbb{1})C] = \rho * C$$

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Higher order maps:

 $\blacksquare \mathcal{W} \leftrightarrow W \in \mathcal{L}(\otimes_i(\mathcal{H}^{A_i^I} \otimes \mathcal{H}^{A_i^O}))$ 

Quantum combs: Choi operator  $\boldsymbol{W}$  has nice additional structure

Can be characterised with semidefinite programming

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Other directions of study in quantum causality

# **Beyond Fixed Causal Structures**

Fundamental questions:

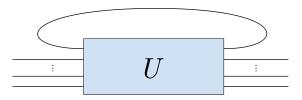
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# **Beyond Fixed Causal Structures**

Fundamental questions:

- What are the physical limits of information processing?
- Does nature allow us to process information in noncausal ways?

One idea: Quantum circuits with closed-timelike curves (CTCs) [Deutsch, 1991]:



- Nature "magically" finds a consistent fixed point solution
- Compatible with general relativity
- Computationally (too?) powerful (P<sub>CTC</sub> = BQP<sub>CTC</sub> = PSPACE) [Aaronson & Watrous, 2008]
- Nonlinear, ...

#### **Quantum Causal Order**

Quantum Causal Structure: Can we have intrinsically quantum causal relations?

- For example, superposition of cause and effect relations?
- Should be linear and well-behaved: quantum supermap (like quantum combs)

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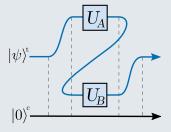
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Use a quantum system to coherently control order that two quantum operations (channels/unitaries) are applied.

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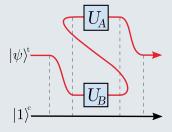
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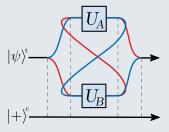
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 $U_B U_A \otimes \left| 0 
ight
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# The Quantum Switch

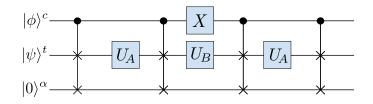
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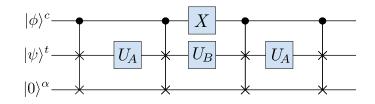
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The Quantum Switch is physically meaningful

- Several experimental realisations with quantum optics
  - Vienna, Brisbane, Shanghai, Concepción, ...

# Application: Commuting/Anticommuting Unitaries

#### Commuting vs. Anticommuting Unitary Problem (Chiribella [2012])

Input: Unitaries  $U_A, U_B$  (oracle access)

**Promise**:  $U_A$  and  $U_B$  either:

- Commute:  $[U_A, U_B] = U_A U_B U_B U_A = 0$ ,
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Recall:  $\mathcal{W}_{switch} : (U_A, U_B) \mapsto U_B U_A \otimes |0\rangle \langle 0|^c + U_A U_B \otimes |1\rangle \langle 1|^c$ 

$$\begin{split} \left|\psi\right\rangle^{t} \otimes \frac{1}{\sqrt{2}}(\left|0\right\rangle^{c} + \left|1\right\rangle^{c}) \xrightarrow{\mathcal{W}_{\text{switch}}} \frac{1}{\sqrt{2}}(U_{B}U_{A}\left|\psi\right\rangle^{t} \otimes \left|0\right\rangle^{c} + U_{A}U_{B}\left|\psi\right\rangle^{t} \otimes \left|1\right\rangle^{c}) \\ &= \frac{1}{2}\{U_{A}, U_{B}\}\left|\psi\right\rangle^{t} \otimes \left|+\right\rangle^{c} - \frac{1}{2}[U_{A}, U_{B}]\left|\psi\right\rangle^{t} \otimes \left|-\right\rangle^{c} \end{split}$$

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Impossible with a quantum comb

# **Quantum** *N*-Switch

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### **Quantum** *N*-Switch

How powerful is quantum control?

First step: Generalise the quantum switch

#### Quantum N-Switch

The quantum N-switch is a supermap  $W_N$ :

$$(U_1,\ldots,U_N)\mapsto \sum_{\pi}U_{\pi(N)}\cdots U_{\pi(1)}\otimes |\pi\rangle\langle\pi|^c,$$

where  $\pi$  is a permutation of  $(1, \ldots, N)$ .

■ Coherent control of all *N*! orders of gates

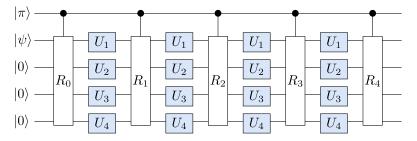
If we initialise the control to  $|\phi\rangle^c = \frac{1}{\sqrt{N!}} \sum_{\pi} |\pi\rangle_c$ , we apply the gates in a superposition of all possible orders:

$$|\psi\rangle^t \otimes |\phi\rangle^c \to \frac{1}{\sqrt{N!}} \sum_{\pi} U_{\pi(N)} \cdots U_{\pi(1)} |\psi\rangle^t \otimes |\pi\rangle^c.$$

# Simulating the *N*-Switch

How much overhead would simulating this with a quantum comb require?

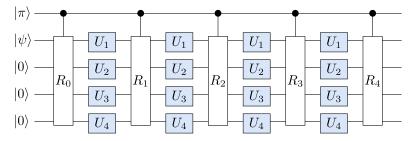
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#### Theorem (Facchini and Perdrix (2014))

Any circuit simulating  $\mathcal{W}_N$  requires at least  $N^2 - o(n^{7/4+\epsilon})$  queries to  $\{U_1, \ldots, U_N\}$ .

### **Fourier Promise Problem**

#### Fourier Promise Problem (Araújo, Costa, Brukner [2014])

Input: Unitaries  $U_1, \ldots, U_N$  (oracle access).

Promise: Let x = 0, ..., N! - 1 be a labelling of permutations  $\pi_x$  and

$$\Pi_x = U_{\pi_x(N-1)} \cdots U_{\pi_x(1)} U_{\pi_x(0)}.$$

Then the unitaries satisfy

$$\forall x, \quad \Pi_x = \omega^{xy} \Pi_0, \quad (\text{where } \omega = e^{i \frac{2\pi}{N!}})$$

for some  $y \in \{0, ..., N! - 1\}$ .

Problem: Find y.

Quantum N-Switch: Solves perfectly (N total queries)

### **Quantum** *N*-**Switch**: **Discussion**

What do we learn from the N-Switch?

- There are physically meaningful computations beyond the circuit model
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  - Other variants: Hadamard promise problem, ....
- Potentially useful for fundamentally quantum problems:
  - Quantum metrology (parameter estimation), ...

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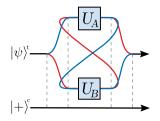
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Not a general model of computation with quantum control!

### **Towards a General Model**

We can have more general "switch-like" computations

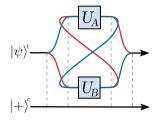
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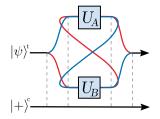
Can we imagine using quantum control in more general ways?

- In quantum switch, control is fixed initially: static control
- How can we have a dynamical, adaptive quantum control structure?

### **Towards a General Model**

We can have more general "switch-like" computations

- Transformations on target between unitaries
- Allow a quantum memory between queries



Can we imagine using quantum control in more general ways?

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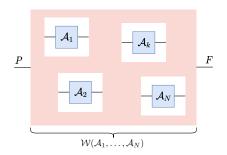
Goal: A generalised model of computation incorporating quantum control

# **Reminder of Scenario**

#### Goal

A generalised model of computation incorporating quantum control:

- A physically well-defined (linear) quantum supermap  $\mathcal{W}$
- Composition of  $\mathcal{A}_1,\ldots,\mathcal{A}_N$  not necessarily in a well-defined, classical order



A subtle requirement:

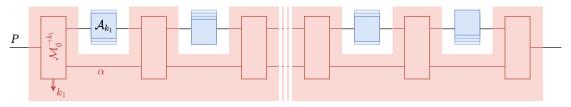
- At the end, every operation should have been applied exactly once
- $\blacksquare$  Necessary for  ${\mathcal W}$  to be linear and well-defined

Simpler situation: quantum circuits with classical control of causal order

Dynamical control structure: determine at each step which operation to apply

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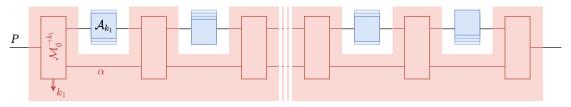
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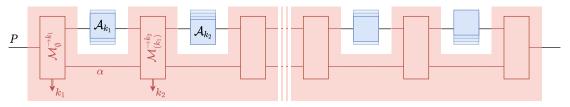
- 1) Input to circuit: a state  $\rho \in \mathcal{L}(\mathcal{H}_P)$
- 2) Perform a quantum instrument  $\{\mathcal{M}_{\emptyset}^{\rightarrow k_1}\}_{k_1}$ . Apply  $\mathcal{A}_{k_1}$  to the target subsystem of  $\mathcal{M}_{\emptyset}^{\rightarrow k_1}(\rho) \in \mathcal{L}(\mathcal{H}^t \otimes \mathcal{H}^{\alpha})$

#### Quantum instruments (generalised quantum measurements)

A quantum instrument is a set  $\{\mathcal{M}_a\}_a$  of CP maps such that  $\mathcal{M} = \sum_a \mathcal{M}_a$  is CPTP. Obtain outcome a with probability  $\operatorname{Tr}[\mathcal{M}_a(\rho)]$  and state becomes  $\mathcal{M}_a(\rho)/\operatorname{Tr}[\mathcal{M}_a(\rho)]$ .

Simpler situation: quantum circuits with classical control of causal order

Dynamical control structure: determine at each step which operation to apply

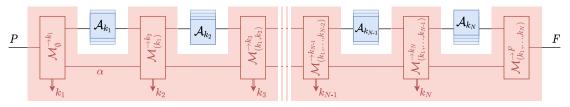


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- 4) Etc.

[Wechs, Dourdent, Abbott, Branciard (2021)]

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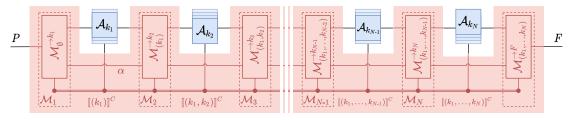
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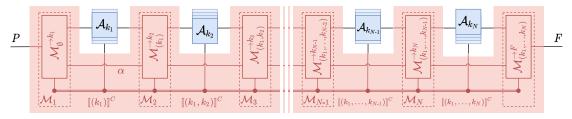
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- Problem: For quantum control, we don't want to destroy coherence by measuring
- Intermediate step: Reformulate classical control with explicit control system



"Classical" control register [[(k<sub>1</sub>,...,k<sub>n</sub>)]] := |(k<sub>1</sub>,...,k<sub>n</sub>))\((k<sub>1</sub>,...,k<sub>n</sub>)|
 Classically controlled operations:

$$\mathcal{M}_1 = \sum_{k_1} \mathcal{M}_{\emptyset}^{\to k_1} \otimes \llbracket (k_1) \rrbracket$$
$$\mathcal{M}_2 = \sum_{k_1, k_2} \mathcal{M}_{\{k_1\}}^{\to k_2} \otimes \Pi_{(k_1), k_2}$$

### From Classical to Quantum Control

Turning the classical into quantum control requires a few tweaks:

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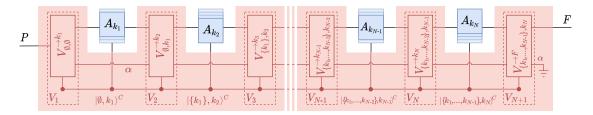
- Quantum control system:  $|\{k_1, \ldots, k_{n-1}\}, k_n\rangle$ 
  - $k_n$ : The operation to apply at slot n
  - $\{k_1, \ldots, k_{n-1}\}$ : History recording which operations have already been used
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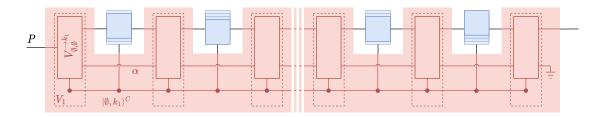
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Circuit evolves coherently, exploring causal structures in a quantum superposition



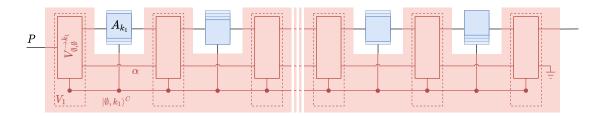
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# Circuits with Quantum Control of Causal Order



•  $V_1: |\psi\rangle^P \to \sum_{k_1} (V_{\emptyset,\emptyset}^{\to k_1} |\psi\rangle)^{t\alpha} \otimes |\emptyset, k_1\rangle^C$ 

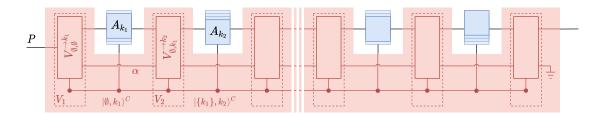
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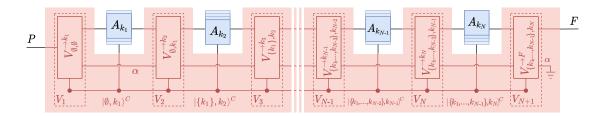
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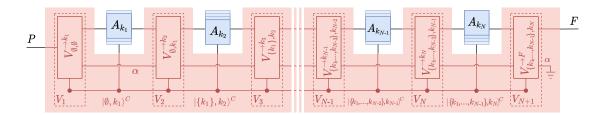


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$$V_{N+1} \colon \to \left( \sum_{(k_1,\dots,k_N)} V_{\{k_1,\dots,k_{N-1}\},k_N}^{\to F} (A_{k_N} \otimes \mathbb{1}^{\alpha}) \cdots (V_{\emptyset,\emptyset}^{\to k_1} |\psi\rangle)^{t\alpha} \right) \otimes |\{1,\dots,N\},F\rangle$$

# Beyond the Quantum Switch

The quantum N-switch can be represented as quantum circuit with quantum control of causal order

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Can we do anything else with this model?

- Yes! We can have quantum dynamical causal structure.
- $\blacksquare$  Before slot n+1 we apply

$$V_{n+1} = \sum_{\mathcal{K}_{n-1}, k_n, k_{n+1}} V_{\mathcal{K}_{n-1}, k_n}^{\to k_{n+1}} \otimes |\mathcal{K}_{n-1} \cup \{k_n\}, k_{n+1} \rangle \langle \mathcal{K}_{n-1}, k_n |$$

•  $k_{n+1}$  can depend on outcome of previous operations

- Examples that exploit this have very recently been devised
- Computations making use of this are active research topic

[Wechs, Dourdent, Abbott, Branciard (2021)]

What can we do with quantum circuits with quantum control of causal order (QC-QCs)?

- $\blacksquare$  Can still simulate with  ${\cal O}(N^2)$  queries
- What types of problems can we solve "efficiently" with QC-QCs?
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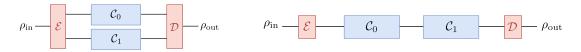
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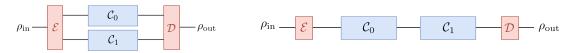
# Quantum Shannon Theory without Causal Order

Quantum Shannon theory: quantum states, channels, but classical, causal trajectories for information carriers

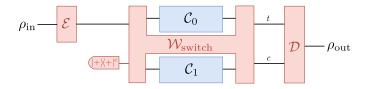


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Quantum Shannon theory: quantum states, channels, but classical, causal trajectories for information carriers



What if we allow for superposition, or quantum control, of trajectories?



[Chiribella & Kristjánsson, Quantum Shannon theory with superpositions of trajectories (2019)]

## **Causal Activation of Capacity**

Recall (fully) depolarising channel:

$$\mathcal{N}: \rho \mapsto \operatorname{Tr}[\rho] \frac{1}{d}$$

•  $\mathcal{N} = \mathcal{N} \circ \mathcal{N}$  has zero capacity (classical and quantum)

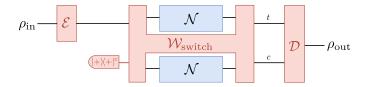
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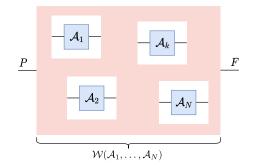
Causal activation:  $\mathcal{W}_{|+\chi|+}(\mathcal{N},\mathcal{N})$  has nonzero classical capacity



- Similar results for quantum capacity with other channels
- Meaning of these results still debated

[Salek, Ebler, Chiribella (2018)]

## **Causal Structure of Quantum Supermaps**



Higher order operations, a.k.a. quantum supermaps:

- Quantum combs
- Quantum circuits with quantum control (QC-QCs)

???

#### **Process Matrices**

Quantum supermaps beyond QC-QCs exist

- Process matrix framework: most general operations compatible with local causality, but without any global causality constraint
- Introduced independently to study quantum gravity
- Beyond QC-QCs, not clear if they have a physical interpretation

[Oreshkov, Costa, Brukner (2012)]

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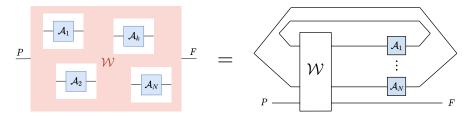
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 $\blacksquare \mathsf{BQP}_{\ell\mathsf{CTCs}} \subseteq \mathsf{PP}$ 

[Araújo, Guérin, Baumeler (2017)]

Need new tools to analyse causal structure of quantum information processing

- Quantum combs (higher order quantum maps)
- Quantum causal models

By exploiting quantum control, can process information in an indefinite causal order

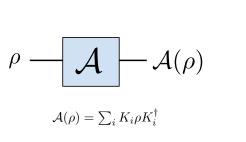
- Quantum switch
- Generalised circuit model incorporating quantum control
- Provides some advantages in quantum information
- Only beginning to understand its potential and limits
- Fundamental questions
  - What are the limits of physical information processing?

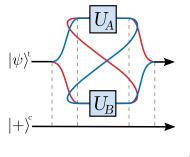
# The Quantum Switch as a Supermap

The Quantum Switch is a quantum supermap:  $(\mathcal{A}, \mathcal{B}) \mapsto \mathcal{W}_{\mathsf{switch}}(\mathcal{A}, \mathcal{B})$ 

• If  $\mathcal{A}, \mathcal{B}$  are quantum channels with Kraus operators  $\{K_i\}_i, \{L_j\}_j$ , then  $\mathcal{W}_{switch}(\mathcal{A}, \mathcal{B})$  has Kraus operators

$$S_{ij} = L_j K_i \otimes |0\rangle \langle 0|^c + K_i L_j \otimes |1\rangle \langle 1|^c.$$





$$[\mathcal{W}_{\mathsf{switch}}(\mathcal{A},\mathcal{B})](\rho^{tc}) = \sum_{i,j} S_{ij} \rho^{tc} S_{ij}^{\dagger}$$

## **Fourier Promise Problem Working**

Quantum N-Switch: Solves perfectly (N total queries)

1) Initial state: 
$$\frac{1}{\sqrt{N!}} \sum_{x=0}^{N!-1} |\psi\rangle^t \otimes |x\rangle^c$$

2) Apply N-switch: 
$$ightarrow rac{1}{\sqrt{N!}} \sum_{x=0}^{N!-1} \Pi_x \ket{\psi}^t \otimes \ket{x}^c$$

3) Apply QFT to control:  $\rightarrow \frac{1}{N!} \sum_{x,s=0}^{N!-1} \omega^{-xs} \Pi_x |\psi\rangle^t \otimes |s\rangle^c = \sum_{x,s=0}^{N!-1} \omega^{x(y-s)} \Pi_0 |\psi\rangle^t \otimes |s\rangle^c$ 

4) Measure the control: 
$$p(s) = \frac{1}{(N!)^2} \|\sum_{x=0}^{N!-1} \omega^{x(y-s)} \Pi_0 |\psi\rangle^t \otimes |s\rangle^c \|^2 = \delta_{s,y}$$

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With a quantum comb?

- The simulating the switch requires  $\Omega(N^2)$  queries.
- Other approaches: need to determine phase to accuracy  $2\pi/N!$
- Rigorous error-robust analysis lacking thus far...

$$\begin{split} \|\psi\rangle^{i} & \bigvee_{|+\rangle^{c}} \\ \bullet V_{1} = \sum_{k_{1}} V_{\emptyset,\emptyset}^{\to k_{1}} \otimes |\emptyset, k_{1}\rangle^{C} = \sum_{k_{1}} \mathbb{1}^{P_{t} \to t} \otimes \langle k_{1}|^{P_{c}} \otimes |\emptyset, k_{1}\rangle^{C} \\ \bullet V_{2} = \sum_{(k_{1},k_{2})} V_{\emptyset,k_{1}}^{\to k_{2}} \otimes |\{k_{1}\}, k_{2}\rangle\langle \emptyset, k_{1}|^{C} = \sum_{k_{1},k_{2}} \mathbb{1} \otimes |\{k_{1}\}, k_{2}\rangle\langle \emptyset, k_{1}|^{C} \\ \bullet V_{3} = \sum_{(k_{1},k_{2})} V_{\{k_{1}\},k_{2}}^{\to F} \otimes \langle \{k_{1}\}, k_{2}| = \sum_{(k_{1},k_{2})} \mathbb{1}^{t \to F_{t}} \otimes |k_{1}\rangle^{F_{c}} \otimes \langle \{k_{1}\}, k_{2}| \\ \end{split}$$

ττ

$$\begin{aligned} \|\psi^{i}-\psi^{i}-\psi^{i}\| \\ \|\psi^{i}-\psi^{i}-\psi^{i}\| \\ \|\psi^{i}-\psi^{i}\| \\ \|\psi^{i}\| \\ \|\psi^{i}-\psi^{i}\| \\ \|\psi^{i}\| \\ \|\psi^{i}\|$$

$$\begin{split} V_{1} &= \sum_{k_{1}} V_{\emptyset,\emptyset}^{\rightarrow k_{1}} \otimes |\emptyset, k_{1}\rangle^{C} = \sum_{k_{1}} \mathbb{1}^{P_{t} \rightarrow t} \otimes \langle k_{1}|^{P_{c}} \otimes |\emptyset, k_{1}\rangle^{C} \\ V_{2} &= \sum_{(k_{1},k_{2})} V_{\emptyset,k_{1}}^{\rightarrow k_{2}} \otimes |\{k_{1}\}, k_{2}\rangle\langle \emptyset, k_{1}|^{C} = \sum_{k_{1},k_{2}} \mathbb{1} \otimes |\{k_{1}\}, k_{2}\rangle\langle \emptyset, k_{1}|^{C} \\ V_{3} &= \sum_{(k_{1},k_{2})} V_{\{k_{1}\},k_{2}}^{\rightarrow F_{c}} \otimes \langle \{k_{1}\}, k_{2}| = \sum_{(k_{1},k_{2})} \mathbb{1}^{t \rightarrow F_{t}} \otimes |k_{1}\rangle^{F_{c}} \otimes \langle \{k_{1}\}, k_{2}| \\ &|\psi\rangle^{P_{t}} \otimes |\phi\rangle^{P_{c}} \xrightarrow{V_{1}} |\psi\rangle^{t} \otimes (\langle A|\phi\rangle |\emptyset, A\rangle + \langle B|\phi\rangle |\emptyset, B\rangle) \\ &\stackrel{cU}{\longrightarrow} U_{A} |\psi\rangle^{t} \otimes \langle A|\phi\rangle |\emptyset, A\rangle + U_{B} |\psi\rangle^{t} \otimes \langle B|\phi\rangle |\emptyset, B\rangle \\ &\stackrel{V_{2}}{\longrightarrow} U_{A} |\psi\rangle^{t} \otimes \langle A|\phi\rangle |\{A\}, B\rangle + U_{B} |\psi\rangle^{t} \otimes \langle B|\phi\rangle |\{B\}, A\rangle \\ &\stackrel{cU}{\longrightarrow} U_{B}U_{A} |\psi\rangle^{t} \otimes \langle A|\phi\rangle |\{A\}, B\rangle + U_{A}U_{B} |\psi\rangle^{t} \otimes \langle B|\phi\rangle |\{B\}, A\rangle \\ &\stackrel{V_{3}}{\longrightarrow} U_{B}U_{A} |\psi\rangle^{F_{t}} \otimes \langle A|\phi\rangle |A\rangle + U_{A}U_{B} |\psi\rangle^{F_{t}} \otimes \langle B|\phi\rangle |B\rangle \end{split}$$

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$$\begin{split} \|\psi^{i}-\psi^{i}\|_{l+\gamma} & = \sum_{k_{1}} V_{\emptyset,\emptyset}^{\to k_{1}} \otimes |\emptyset,k_{1}\rangle^{C} = \sum_{k_{1}} \mathbb{1}^{P_{t} \to t} \otimes \langle k_{1}|^{P_{c}} \otimes |\emptyset,k_{1}\rangle^{C} \\ \mathbb{V}_{2} &= \sum_{(k_{1},k_{2})} V_{\emptyset,k_{1}}^{\to k_{1}} \otimes |\{k_{1}\},k_{2}\rangle\langle\emptyset,k_{1}|^{C} = \sum_{k_{1},k_{2}} \mathbb{1} \otimes |\{k_{1}\},k_{2}\rangle\langle\emptyset,k_{1}|^{C} \\ \mathbb{V}_{3} &= \sum_{(k_{1},k_{2})} V_{\{k_{1}\},k_{2}}^{\to F} \otimes \langle\{k_{1}\},k_{2}| = \sum_{(k_{1},k_{2})} \mathbb{1}^{t \to F_{t}} \otimes |k_{1}\rangle^{F_{c}} \otimes \langle\{k_{1}\},k_{2}| \\ & |\psi\rangle^{P_{t}} \otimes |\phi\rangle^{P_{c}} \xrightarrow{V_{1}} |\psi\rangle^{t} \otimes (\langle A|\phi\rangle |\emptyset,A\rangle + \langle B|\phi\rangle |\emptyset,B\rangle) \\ & \xrightarrow{cU} U_{A} |\psi\rangle^{t} \otimes \langle A|\phi\rangle |\{A\},B\rangle + U_{B} |\psi\rangle^{t} \otimes \langle B|\phi\rangle |\emptyset,B\rangle \\ & \xrightarrow{V_{2}} U_{A} |\psi\rangle^{t} \otimes \langle A|\phi\rangle |\{A\},B\rangle + U_{B} |\psi\rangle^{t} \otimes \langle B|\phi\rangle |\{B\},A\rangle \\ & \xrightarrow{CU} U_{B}U_{A} |\psi\rangle^{F_{t}} \otimes \langle A|\phi\rangle |A\rangle + U_{A}U_{B} |\psi\rangle^{F_{t}} \otimes \langle B|\phi\rangle |B\rangle \end{split}$$

$$\begin{aligned} \|\psi^{i}-\psi^{i}-\psi^{i}\| \\ \|\psi^{i}-\psi^{i}-\psi^{i}\| \\ \|\psi^{i}-\psi^{i}\| \\ \|\psi^{i}-\psi^{i}\| \\ \|\psi^{i}-\psi^{i}\| \\ \|\psi^{i}-\psi^{i}\| \\ \|\psi^{i}-\psi^{i}\| \\ \|\psi^{i}\| \\ \|\psi^{i}-\psi^{i}\| \\ \|\psi^{i}\| \\ \|\psi$$

$$\begin{aligned} \|\psi^{i}-\psi^{i}\|_{l+\gamma} & = \sum_{k_{1}} V_{\emptyset,\emptyset}^{\to k_{1}} \otimes |\emptyset,k_{1}\rangle^{C} = \sum_{k_{1}} \mathbb{1}^{P_{t} \to t} \otimes \langle k_{1}|^{P_{c}} \otimes |\emptyset,k_{1}\rangle^{C} \\ \mathbb{V}_{2} &= \sum_{(k_{1},k_{2})} V_{\emptyset,k_{1}}^{\to k_{2}} \otimes |\{k_{1}\},k_{2}\rangle\langle\emptyset,k_{1}|^{C} = \sum_{k_{1},k_{2}} \mathbb{1} \otimes |\{k_{1}\},k_{2}\rangle\langle\emptyset,k_{1}|^{C} \\ \mathbb{V}_{3} &= \sum_{(k_{1},k_{2})} V_{\{k_{1}\},k_{2}}^{\to F_{c}} \otimes \langle\{k_{1}\},k_{2}| = \sum_{(k_{1},k_{2})} \mathbb{1}^{t \to F_{t}} \otimes |k_{1}\rangle^{F_{c}} \otimes \langle\{k_{1}\},k_{2}| \\ &|\psi\rangle^{P_{t}} \otimes |\phi\rangle^{P_{c}} \xrightarrow{V_{1}} |\psi\rangle^{t} \otimes (\langle A|\phi\rangle |\emptyset,A\rangle + \langle B|\phi\rangle |\emptyset,B\rangle) \\ &\stackrel{cU}{\to} U_{A} |\psi\rangle^{t} \otimes \langle A|\phi\rangle |\{A\},B\rangle + U_{B} |\psi\rangle^{t} \otimes \langle B|\phi\rangle |\emptyset,B\rangle \\ &\stackrel{V_{2}}{\to} U_{A} |\psi\rangle^{t} \otimes \langle A|\phi\rangle |\{A\},B\rangle + U_{B} |\psi\rangle^{t} \otimes \langle B|\phi\rangle |\{B\},A\rangle \\ &\stackrel{U_{3}}{\to} U_{B}U_{A} |\psi\rangle^{F_{t}} \otimes \langle A|\phi\rangle |A\rangle + U_{A}U_{B} |\psi\rangle^{F_{t}} \otimes \langle B|\phi\rangle |B\rangle \end{aligned}$$

$$\begin{split} \|\psi^{i}-\psi^{i}\|_{l+\gamma} & = \sum_{k_{1}} V_{\emptyset,\emptyset}^{\to k_{1}} \otimes |\emptyset,k_{1}\rangle^{C} = \sum_{k_{1}} \mathbb{1}^{P_{t} \to t} \otimes \langle k_{1}|^{P_{c}} \otimes |\emptyset,k_{1}\rangle^{C} \\ \mathbb{1} V_{2} &= \sum_{(k_{1},k_{2})} V_{\emptyset,k_{1}}^{\to k_{1}} \otimes |\{k_{1}\},k_{2}\rangle\langle\emptyset,k_{1}|^{C} = \sum_{k_{1},k_{2}} \mathbb{1} \otimes |\{k_{1}\},k_{2}\rangle\langle\emptyset,k_{1}|^{C} \\ \mathbb{1} V_{3} &= \sum_{(k_{1},k_{2})} V_{\{k_{1}\},k_{2}}^{\to F_{c}} \otimes \langle \{k_{1}\},k_{2}| = \sum_{(k_{1},k_{2})} \mathbb{1}^{t \to F_{t}} \otimes |k_{1}\rangle^{F_{c}} \otimes \langle \{k_{1}\},k_{2}| \\ &|\psi\rangle^{P_{t}} \otimes |\phi\rangle^{P_{c}} \xrightarrow{V_{1}} |\psi\rangle^{t} \otimes (\langle A|\phi\rangle |\emptyset,A\rangle + \langle B|\phi\rangle |\emptyset,B\rangle) \\ &\stackrel{cU}{\to} U_{A} |\psi\rangle^{t} \otimes \langle A|\phi\rangle |\{A\},B\rangle + U_{B} |\psi\rangle^{t} \otimes \langle B|\phi\rangle |\emptyset,B\rangle \\ &\stackrel{V_{2}}{\to} U_{A} |\psi\rangle^{t} \otimes \langle A|\phi\rangle |\{A\},B\rangle + U_{B} |\psi\rangle^{t} \otimes \langle B|\phi\rangle |\{B\},A\rangle \\ &\stackrel{CU}{\to} U_{B}U_{A} |\psi\rangle^{F_{t}} \otimes \langle A|\phi\rangle |A\rangle + U_{A}U_{B} |\psi\rangle^{F_{t}} \otimes \langle B|\phi\rangle |B\rangle \end{split}$$