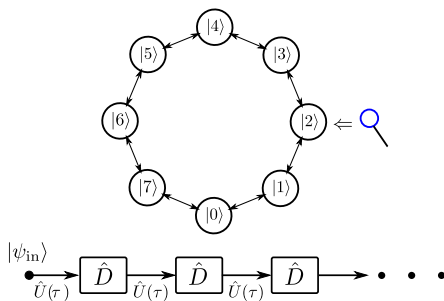


Large fluctuations of the first detected quantum return time

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20th Jan. 2020

Repeated strong measurements until the first success:



$$\hat{U}(\tau) = \exp(-iH\tau) \quad (\hbar = 1), \quad \hat{D} = |\psi_d\rangle\langle\psi_d|$$

Outcome: a string of length n : “0, 0, 0, ..., 1” (“0” for non-detection, and “1” for detection)

If

- First-detected “return” time (FDRt): $|\psi_d\rangle = |\psi_{in}\rangle$
- The system’s energy spectrum is discrete

Known results: (Grünbaum, Velázquez et al. 2013)

- The FDR amplitude, the generating function(Z Transform):
$$\phi_n = \langle \psi_{in} | [\hat{U}(\tau)(1 - \hat{D})]^{n-1} \hat{U}(\tau) | \psi_{in} \rangle, \tilde{\phi}(z) = \frac{\langle \psi_{in} | \sum_{n=1}^{\infty} z^n \hat{U}(n\tau) | \psi_{in} \rangle}{1 + \langle \psi_{in} | \sum_{n=1}^{\infty} z^n \hat{U}(n\tau) | \psi_{in} \rangle}$$
- The system is recurrent
- The mean: $\langle n \rangle = \sum_{n=1}^{\infty} n |\phi_n|^2 = w \in \mathbb{Z}$, = the number of zeros $\{z_i\}$ of $\tilde{\phi}(z)$ [or the phases $\exp(iE_k\tau)$] =winding number of $\tilde{\phi}(e^{i\theta})$
- The variance:

$$\text{Var}(n) = \langle n^2 \rangle - \langle n \rangle^2 = \sum_{i,j=1}^w \frac{2z_i z_j^*}{1 - z_i z_j^*}$$

- Charge Theory: $\{z_i\}$ (excluding $z_0 = 0$)

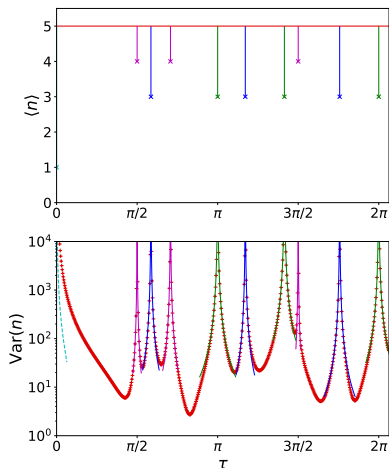
$$V(z) = \sum_{k=1}^w p_k \ln |e^{iE_k\tau} - z| \Rightarrow \text{stationary points}$$

$$F(z) = \sum_{k=1}^w \frac{p_k}{e^{iE_k\tau} - z} \Rightarrow \text{equilibrium points}$$

where $p_k = \sum_{l=1}^{g_k} |\langle \psi_{in} | E_{kl} \rangle|^2 \implies$ charges on the unit circle.

Example

$$H = -\gamma \sum_{x=0}^7 [|x\rangle\langle x+1| + |x+1\rangle\langle x| - 2|x\rangle\langle x|].$$
$$E_k/\gamma = 2 - 2 \cos(\pi k/4)$$



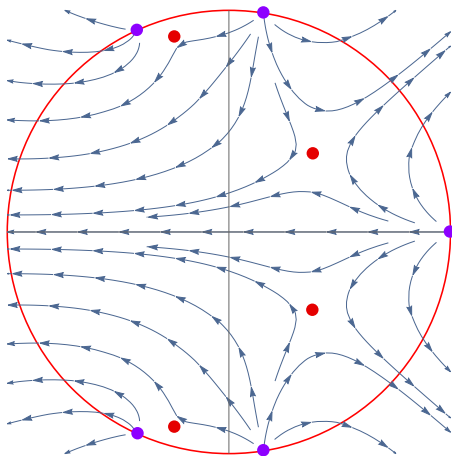


Figure: Eight-site ring. $\tau = 1$.

Mechanism of Jumps of $\langle n \rangle$ and Diverging $\text{Var}(n)$

Some zero(s) approach the unit circle:

- Mean: when the zero(s) “reach” the unit circle, it(they) does(do) not count into $\langle n \rangle$ (canceled out)
- $\text{Var}(n)$: when the zero(s) is(are) approaching to the unit circle \implies vanishing denominator $(1 - z_i z_j^*) \implies$ diverging $\text{var}(n)$
“Signature” to jumps of the topological number

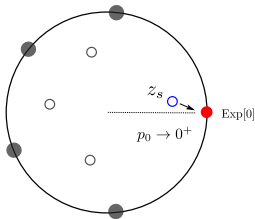
From basic knowledge of electrostatics:

- some “weak” charge(s) (“Shrinking”, Grünbaum, Velázquez et al. 2013): L_1 point in sun-earth system
- charges merging (“Fusion”, Grünbaum, Velázquez et al. 2013): neighboring stars

Perturbation Method

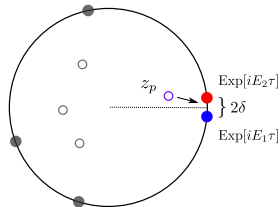
$$F(z) = \sum_{k=1}^w \frac{p_k}{e^{iE_k\tau} - z} = 0 \Rightarrow \text{critical zero } z_c$$

$$\text{Var}(n) \sim \frac{2|z_c|^2}{1 - |z_c|^2}$$



$$z_s \sim 1 - \frac{p_0}{\sum_{j \neq 0} p_j / [1 - \exp(iE_j\tau)]}$$

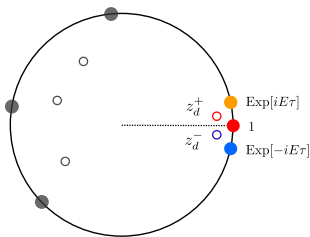
$$\text{Var}(n) \sim \frac{1}{2p_0} \left\{ 1 + \left[\sum_{j \neq 0} p_j \cot [(E_j - E_0)\tau/2] \right]^2 \right\}$$



$$z_p \sim 1 + i \frac{p_1 - p_2}{p_1 + p_2} \delta + \left[\frac{4p_1 p_2}{(p_1 + p_2)^3} \sum_{j \neq 1, 2} \frac{p_j}{e^{iE_j\tau} - 1} - \frac{1}{2} \right] \delta^2$$

$$\text{Var}(n) \sim 2 \frac{(p_1 + p_2)^3}{p_1 p_2} \frac{1}{\tau^2 (\bar{E}_2 - \bar{E}_1)^2}$$

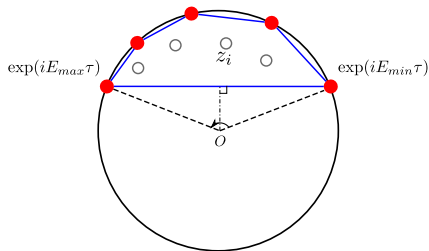
Triple-Charge Theory: Symmetric Scenario



$$\text{Var}(n) \sim \sum_{\sigma=\pm} \frac{2|z_d^\sigma|^2}{1-|z_d^\sigma|^2} + \underbrace{\left(\frac{2z_d^+(z_d^-)^*}{1-z_d^+(z_d^-)^*} + \text{c.c.} \right)}_{\mathcal{M}},$$

$$\text{Var}(n) \sim 16 \frac{(p_0 + 2p)^2}{p} \frac{1}{\tau^2 (\bar{E}_+ - \bar{E}_-)^2}, \text{ the mixing term } \mathcal{M} \text{ is finite.}$$

Zeno Regime–Lower Bound



$$\text{Var}(n) \geq (w - 1) \left[2 \cot^2 (\Delta E_m \tau / 2) - w + 2 \right]$$

Time-Energy Uncertainty Principle

Assume $\Delta E_m \tau \ll \hbar$

$$(\Delta E_m)^2 (\Delta t_{\text{det}})^2 \gtrsim 8(\langle n \rangle - 1) \hbar^2,$$

where $\Delta E_m = E_{\text{max}} - E_{\text{min}}$,

$$\Delta t_{\text{det}} = \sqrt{\langle (n\tau)^2 \rangle - \langle n\tau \rangle^2} = \tau \sqrt{\text{Var}(n)}$$

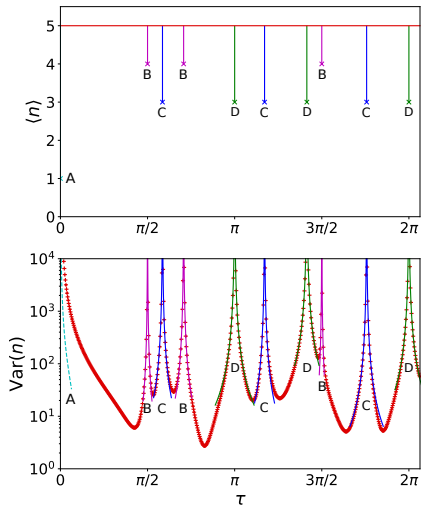
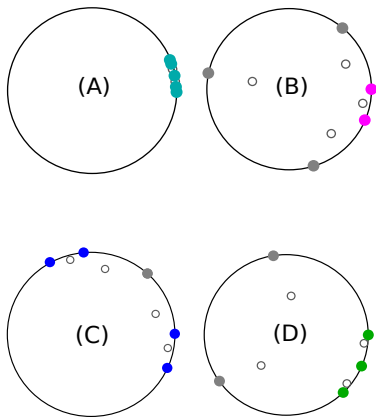
In the mathematical limit $\langle n \rangle = w = 1$, the particle is detected at the first attempt for any sampling time τ .

$$\Delta t_{\text{det}} = 0$$

The presence of the factor $\langle n \rangle - 1$ is physically reasonable.

Example

Eight-Site Ring



Interacting-Boson Model

$$H = -\frac{J}{2}(\hat{a}_1^\dagger \hat{a}_r + \hat{a}_r^\dagger \hat{a}_l) + U(\hat{n}_l^2 + \hat{n}_r^2), \quad \hat{n}_{l,r} = \hat{a}_{l,r}^\dagger \hat{a}_{l,r}$$

$$E_0 = 3U - \sqrt{U^2 + J^2}, E_1 = 4U, E_2 = 3U + \sqrt{U^2 + J^2}.$$

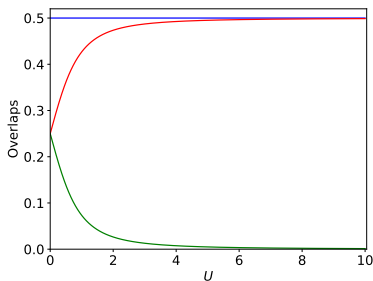


Figure: $|\psi_{\text{in}}\rangle = |2, 0\rangle$, and J is set as 1. The green curve represents p_0 , and the blue/red is p_1/p_2 . U is large, the ground state is almost $|1, 1\rangle$.

Another Example

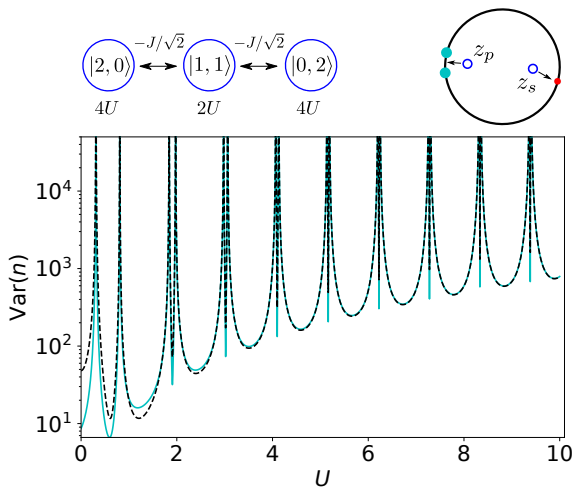


Figure: $J = 1$, $\tau = 3$

- Diverging $\text{Var}(n)$ accompany isolated jumps of $\langle n \rangle$
- We quantify the magnitude of large fluctuations of first detected return time based on four scenarios: [single-weak charge](#), [two-charge merging](#), [three-charge merging](#), [Zeno regime](#)
- Topology-dependent time-energy uncertainty relation

Thank You!

R. Yin, K. Ziegler, F. Thiel and E. Barkai,
Phys. Rev. Research 1, 033086(Editors' Suggestion)