

Additive Combinatorics methods in Fractal Geometry II

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Review: dynamical self-similarity

Definition

- G is a compact Abelian group, and $h \in G$ is such that the orbit $\{nh : n \in \mathbb{Z}\}$ is dense. We let $T(g) = g + h$.
- $\lambda \in (0, 1)$ is a contraction parameter.
- $\Delta(x) : G \rightarrow \mathcal{A}_{\mathbb{C}}^d$ is a map taking values in purely atomic measures in \mathbb{R}^d with at most C atoms.

We call (G, T, λ, Δ) a **model**. The measures

$$\mu_x = *_{n=1}^{\infty} \mathcal{S}_{\lambda^n} \Delta(T^n x)$$

are called **dynamical self-similar measures** generated by the model.

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Further examples of dynamical self-similar measures

- 1 **Self-homothetic measures on \mathbb{R}^d** : they correspond to $G = \{0\}$, λ the (common) contraction of maps in the IFS and $\Delta = \sum_i p_i \delta_{t_i}$ (where $t_i \in \mathbb{R}^d$ are translations) is built from the translations and the probabilities of the IFS.
- 2 If μ, ν are two measures as above with contractions λ_1, λ_2 , then $\mu * \mathcal{S}_{e^x} \nu$ are DSSM where G is a finite group if $\log \lambda_2 / \log \lambda_1 \in \mathbb{Q}$, and the circle otherwise. This extends to

$$\mu_1 * \mathcal{S}_{e^{x_2}} \mu_2 * \cdots * \mathcal{S}_{e^{x_m}} \mu_m.$$

- 3 **A homogeneous self-similar measure in dimension d** is, by definition, a measure of the form

$$\mu = *_{n=1}^{\infty} \mathcal{S}_{\lambda^n} O^n \Delta, \quad O \in \mathbb{O}_d, \lambda \in (0, 1), \Delta \in \mathcal{A}.$$

It can be realized as a DSSM where $G = \overline{\langle O \rangle}$, $h = O$, $\Delta(g) = g\Delta$.

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It can be realized as a DSSM where $G = \overline{\langle O \rangle}$, $h = O$, $\Delta(g) = g\Delta$.

Discrete approximations and shifted self-similarity

$$\mu_X = *_{n=0}^{\infty} \mathcal{S}_{\lambda^n} \Delta(T^n X).$$

We define the **discrete step- n approximations**

$$\mu_{X,n} = *_{j=0}^{n-1} \mathcal{S}_{\lambda^j} \Delta(T^j X).$$

Note that $\mu_{X,n}$ is purely atomic with $\leq \prod_{j=0}^{n-1} |\text{supp}(\Delta(T^j X))| \leq C^n$ atoms.

We then have the following crucial **shifted self-similarity relationship**:

$$\mu_X = \mu_{X,n} * \mathcal{S}_{\lambda^n} \mu_{T^n X}.$$

This says that μ_X is a convex combination of scaled down (by a factor λ^n) translated copies of $\mu_{T^n X}$.

Review: Frostman exponent

Definition

Let μ be a measure on \mathbb{R}^d . The **Frostman exponent** $\dim_\infty(\mu)$ is the supremum of all s such that

$$\mu(B(x, r)) \leq C_s r^s \quad \text{for all } r \in (0, 1] \text{ and all } x.$$

Review: exponential separation

Definition

We say that a model (G, T, λ, Δ) has **exponential separation** if for Haar-**almost all** $x \in G$ there is $R > 0$ such that following holds for infinitely many n : **the atoms of the discrete approximation**

$$\mu_{n,x} = *_{j=1}^n S_{\lambda^j} \Delta(T^j x)$$

are (distinct and) e^{-Rn} -separated.

Main Theorem: Frostman exponents of DSSM

Theorem (P.S.)

Let (G, T, λ, Δ) be a model with exponential separation on \mathbb{R} . We also assume that the maps $x \mapsto \Delta(x)$ and $x \mapsto \mu_x$ are continuous a.e., and that μ_x is supported on $[0, 1]$. Let

$$s = \min \left(\frac{\int \log \|\Delta(x)\|_\infty dx}{\log \lambda}, 1 \right),$$

where $\|\Delta\|_\infty = \max_y \Delta\{y\}$.

Then $\dim_\infty(\mu_x) = s$ for every $x \in G$.

Moreover, for every $\varepsilon > 0$ there is a constant C_ε such that, for all $x \in G$, $y \in \mathbb{R}$ and $r \in (0, 1]$,

$$\mu_x(B(y, r)) \leq C_\varepsilon r^{s-\varepsilon}$$

Remarks on main theorem

Remarks

- *In the self-similar case (corresponding to constant Δ) a version of this result was obtained by M. Hochman but his version is for Hausdorff dimension rather than Frostman exponents.*
- *While exponential separation has to be checked for **almost all** x , the conclusion holds for **all** x*
- *The transitive translation on a compact Abelian group can be replaced by a uniquely ergodic transformation on a compact metric space.*

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Application 1: Furstenberg's Intersection Conjecture

Theorem (P.S./Meng Wu 2016)

Suppose $\log p / \log q \notin \mathbb{Q}$. If A, B are closed and T_p, T_q -invariant, then

$$\dim_{\mathbb{H}}(A \cap f(B)) \leq \overline{\dim}_{\mathbb{B}}(A \cap f(B)) \leq \max(\dim_{\mathbb{H}}(A) + \dim_{\mathbb{H}}(B) - 1, 0)$$

for all affine bijections $f : \mathbb{R} \rightarrow \mathbb{R}$.

Remark

A T_p -invariant set A can be embedded in a T_{p^n} -Cantor set of dimension $\dim_{\mathbb{H}}(A) + \varepsilon$ (the allowed digits are the length- n sequences appearing in A). So it is enough to prove the theorem under the assumption that A, B are p, q -Cantor sets.

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Review: convolutions of Cantor measures as DSSM

Let A, B be p, q -Cantor sets, and let μ, ν be the natural measures on them. We saw yesterday that

$$\mu * \mathcal{S}_{e^x} \nu = *_{n=1}^{\infty} \mathcal{S}_{p^{-n}} \Delta(T^n x),$$

where

$$T(x) = x + \log p \pmod{\log q}$$

and

$$\Delta(x) = \begin{cases} \Delta_A * \mathcal{S}_{e^x} \Delta_B & \text{if } x \in [0, \log p) \\ \Delta_A & \text{if } x \in [\log p, \log q) \end{cases} .$$

A corollary of the main result

Corollary

For all $u \in \mathbb{R} \setminus \{0\}$ it holds that

$$\dim_{\infty}(\mu * S_u \nu) = \min(\dim_{\text{H}}(A) + \dim_{\text{H}}(B), 1) =: s.$$

Remarks

- All the assumptions in the main theorem are clear except exponential separation, which is a simple lemma.*
- A small calculation shows that indeed the value of s given by the main theorem is the RHS above.*
- The main theorem gives the corollary for $u \in [1, \log q]$. Using self-similarity it is easy to expand this to all non-zero u .*

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Some additional remarks

- Up to a similarity map, $A \cap (rB + t)$ is the same as $(A \times B) \cap \{y = rx + t\}$.
- $\mu(B(x, r)) \approx r^{\dim_{\text{H}}(A)}$ for $x \in A$, and likewise for B , so

$$(\mu \times \nu)(B(x, r)) \approx r^{\dim_{\text{H}}(A) + \dim_{\text{H}}(B)}$$

for $(x, y) \in A \times B$.

- The convolution $\mu * S_u \nu$ is the push-forward of $\mu \times \nu$ under the projection $\Pi_u(x, y) = x + uy$.

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Conclusion of the proof

- Let $\{B(x_j, \delta)\}_{j=1}^M$ be disjoint collection of balls intersecting $(A \times B) \cap y = -x/u + t$. We need to bound M from above.

$$(\mu \times \nu) \left(\bigcup_{j=1}^M B(x_j, \delta) \right) \gtrsim M \delta^{\dim_{\text{H}}(A) + \dim_{\text{H}}(B)}.$$

- But (since Π_u is Lipschitz and the line $y = -x/u + t$ is the fiber $\Pi_u^{-1}(tu)$)

$$\Pi_u \left(\bigcup_{j=1}^M B(x_j, \delta) \right) \subset B(tu, C\delta).$$

- Since $\dim_{\infty}(\Pi_u(\mu \times \nu)) = \min(\dim_{\text{H}}(A) + \dim_{\text{H}}(B), 1)$, we conclude

$$\begin{aligned} (\mu \times \nu) \left(\bigcup_{j=1}^M B(x_j, \delta) \right) &\leq \Pi_u(\mu \times \nu) \Pi_u \left(\bigcup_{j=1}^M B(x_j, \delta) \right) \\ &\lesssim_{\varepsilon} \delta^{\min(\dim_{\text{H}}(A) + \dim_{\text{H}}(B), 1) - \varepsilon}. \end{aligned}$$

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Exponential separation in the deterministic case I

Fix $\Delta \in \mathcal{A}$, $\lambda \in (0, 1)$. Let

$$\nu = \nu_{\Delta, \lambda} = *_{n=0}^{\infty} \mathbf{S}_{\lambda^n} \Delta.$$

The atoms of

$$\nu_n = *_{j=0}^{n-1} \mathbf{S}_{\lambda^j} \Delta$$

are of the form $P(\lambda)$, for a polynomial of degree $< n$ with coefficients in $D := \text{supp}(\Delta)$.

Therefore **exponential separation holds if and only if there are $R > 0$ and infinitely many n such that**

$$|Q(\lambda)| > e^{-Rn}$$

for all polynomials Q of degree $< n$ with coefficients in $D - D$.

In the Bernoulli convolution setting, $D - D = \{-1, 0, 1\}$.

Exponential separation in the deterministic case II

Lemma (M. Hochman 2014)

- *If $D = \text{supp}\Delta$ is algebraic, then there is exponential separation if and only if λ is not a root of a polynomial with coefficients in $D - D$.*
- *For any fixed Δ , there is exponential separation for all λ outside of a set of zero Hausdorff dimension.*

Corollary

- *$\dim_\infty(\nu_\lambda) = 1$ for all algebraic numbers in $(1/2, 1)$ which are not roots of a $\{-1, 0, 1\}$ -polynomial.*
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Application 2: Densities of Bernoulli convolutions

Theorem (P.S. 2016)

The BC ν_λ has a density in every L^q for $\lambda \in (1/2, 1)$ outside of a set of exceptions of zero Hausdorff dimension.

Separating ν_λ into large and pseudorandom parts

$$\nu_\lambda = *_{n=0}^{\infty} \mathcal{S}_{\lambda^n} \Delta = \left(*_{k|n}^{\infty} \mathcal{S}_{\lambda^n} \Delta \right) * \left(*_{k \nmid n}^{\infty} \mathcal{S}_{\lambda^n} \Delta \right) =: \nu_{\lambda^k} * \eta_{\lambda,k}.$$

- 1 Erdős-Kahane: for all $\lambda \in (0, 1)$ outside of a set of zero Hausdorff dimension, $|\hat{\nu}_\lambda(\xi)| \leq C_\lambda |\xi|^{-\delta(\lambda)}$: **polynomial Fourier decay**.
- 2 The measures η_λ are also homogeneous self-similar measures: the contraction ratio is λ^k and the atomic measure is

$$*_{j=0}^{k-1} \mathcal{S}_{\lambda^j} \Delta.$$

If exponential separation holds for ν_λ then it also holds for $\eta_{\lambda,k}$ (fewer atoms to consider, we are skipping some digits). So from the main theorem we have

$$\dim_\infty(\eta_{\lambda,k}) = \min \left(\frac{(k-1) \log 2}{k \log(1/\lambda)}, 1 \right).$$

for all $\lambda \in (1/2, 1)$ outside of a zero-dimensional set of exceptions, provided k is taken large enough.

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Conclusion of the proof

By taking the union of the exceptional sets over all k , we get that for $\lambda \in (1/2, 1)$ outside of a set of zero Hausdorff dimension, we can split

$$\nu_\lambda = \nu'_\lambda * \eta_\lambda,$$

where

- ν'_λ has power Fourier decay,
- $\dim_\infty(\eta_\lambda) = 1$.

Proposition (P.S.-B. Solomyak 2016)

*If ρ has power Fourier decay and η has full Frostman exponent, then the convolution $\rho * \eta$ is absolutely continuous and the density is in L^q for all finite q .*

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L^q dimensions

Definition

Given a probability μ on \mathbb{R}^d and $q \in (1, \infty)$, we let

$$S_n(\mu, q) = \sum_{I \in \mathcal{D}_n} \mu(I)^q,$$

$$\dim_q(\mu) = \liminf_{n \rightarrow \infty} \frac{\log S_n(\mu, q)}{n(1 - q)} \in [0, d].$$

- $q \mapsto \dim_q(\mu)$ is non-increasing and $\dim_q(\mu) \rightarrow \dim_\infty(\mu)$ as $q \rightarrow \infty$.
- The main theorem holds not only for Frostman exponents but also for L^q dimensions.
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Main Theorem: L^q dimensions of DSSM

Theorem (P.S.)

Let (G, T, λ, Δ) be a model with exponential separation on \mathbb{R} . We also assume that the maps $x \mapsto \Delta(x)$ and $x \mapsto \mu_x$ are continuous a.e., and that μ_x is supported on $[0, 1]$. Let

$$s(q) = \min \left(\frac{\int \log \|\Delta(x)\|_q^q dx}{(q-1) \log \lambda}, 1 \right),$$

where $\|\Delta\|_q^q = \sum_y \Delta(y)^q$.

Then

$$\dim_q(\mu_x) = s(q)$$

for **every** $x \in G$ and $q > 1$.

Tools involved in the proof

- 1 **Additive combinatorics**: an inverse theorem for the L^q norm of the convolution of two finitely supported measures (Balog-Szemerédi-Gowers Theorem, Bourgain's additive part of discretized sum-product results).
- 2 **Ergodic theory**: key role played by subadditive cocycle over a uniquely ergodic transformation (cocycle borrowed from Nazarov-Peres-S. 2012, uses the proof of the subadditive ergodic theorem given by Katznelson-Weiss).
- 3 **Multifractal analysis** (L^q spectrum, regularity at points of differentiability).
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A submultiplicative cocycle

$$\dim_q(\mu_x) = \lim_{m \rightarrow \infty} \frac{\log \left(\sum_{I \in \mathcal{D}_m} \mu(I)^q \right)}{(1-q)m}.$$

Let

$$\psi_n(x) = \sum_{I \in \mathcal{D}_{m(n)}} \mu(I)^q,$$

where $2^{m(n)} \approx \lambda^n$ (so that $|I| \approx \lambda^n$). Then

$$\dim_q(\mu_x) = \lim_{n \rightarrow \infty} \frac{\log \psi_n(x)}{(q-1)(\log \lambda)n}.$$

Lemma

$$\psi_{n+k}(x) \leq C_q \psi_n(x) \psi_k(T^n x).$$

Use of unique ergodicity

$$\psi_{n+k}(x) \leq C_q \psi_n(x) \psi_k(T^n x).$$

By the subadditive ergodic theorem, there exists $D(q)$ such that

$$\lim_{n \rightarrow \infty} \frac{\log \psi_n(x)}{(q-1)(\log \lambda)n} = D(q) \quad \text{for a.e. } x \in G.$$

Lemma (Furman; follows from the Katznelson-Weiss proof of the subadditive ergodic theorem)

$$\liminf_{n \rightarrow \infty} \frac{\log \psi_n(x)}{(q-1)(\log \lambda)n} \geq D(q) \quad \text{uniformly in } x \in G.$$

Reduction from “everywhere” to “almost everywhere”

We need to show that

$$\dim_q(\mu_x) = \lim_{n \rightarrow \infty} \frac{\log \psi_n(x)}{(q-1)(\log \lambda)n} = \min \left(\frac{\int \log \|\Delta(x)\|_q^q dx}{(q-1) \log \lambda}, 1 \right) =: s(q).$$

The upper bound $\dim_q(\mu, x) \leq s(q)$ for all x is easy. But we saw that

$$\liminf_{n \rightarrow \infty} \frac{\log \psi_n(x)}{(q-1)(\log \lambda)n} \geq D(q) \text{ for all } x.$$

So it is enough to show that $D(q) = s(q)$. In other words, it is enough to prove that

$$\lim_{m \rightarrow \infty} \frac{\log \sum_{I \in \mathcal{D}_m} \mu(I)^q}{(1-q)m} = s(q) \text{ for almost all } x \in G.$$

End of part II

Merci beaucoup!