

Int'l Workshop on Geometric Quantization and Applications

Note Title

Deformation of the Prequantum Action

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- geometric quantisation  
background + joint work with W. Kirwin (2006)
- deformation quantisation  
background + joint work with A. Yoshioka (to appear)
- fermionic systems  
background + W. (2015 + to appear)

## Classical mechanics

$(M^{2n}, \omega)$  symplectic manifold  $\omega \in \Omega^2(M)$  non-degenerate,  $d\omega = 0$ .

Locally  $\omega = \sum_{i=1}^n dp_i \wedge dq^i$ , contains position and momentum variables

Poisson structure  $\Pi = \omega^{-1} \in \Gamma(M, \Lambda^2 TM)$

Hamiltonian vector field  $\iota_{H_f} \omega = -df$  or  $H_f = \Pi(df, \cdot)$

Poisson bracket  $\{f, g\} = \Pi(df \wedge dg) = -\omega(H_f, H_g)$ .

## Geometric quantisation

Kostant (1970) Souriau (1970)

classical phase space  $(M, \omega) \rightsquigarrow$  Hilbert space  $\mathcal{H}$  of quantum states

prequantisation.  $\begin{array}{c} \mathcal{L} \\ \downarrow \\ M \end{array}$  Hermitian line bundle with connection  $\nabla$ ,  
 $\text{curv}(\nabla) = \frac{\omega}{\sqrt{2\pi\hbar}}$ .  $\hbar > 0$  is a constant.

the prequantum action of  $C^\infty(M)$  on  $\Gamma(M, \mathcal{L})$

$$f \mapsto \hat{f} := -\sqrt{2\pi\hbar} \nabla_{H_f} + f$$

$$\{\hat{f}, \hat{g}\} = \frac{1}{\sqrt{2\pi\hbar}} [\hat{f}, \hat{g}], \quad \forall f, g \in C^\infty(M)$$

In physics,  $\Gamma(M, \mathcal{L})$  is too big as the quantum Hilbert space  
(the representation of  $C^\infty(M)$  on  $\Gamma(M, \mathcal{L})$  is highly reducible)

Choose a complex polarisation, an almost complex structure  $J$  on  $M$  compatible with  $\omega$ . Let  $\mathcal{J} =$  the set of such  $J$ .

Then  $T^{\mathbb{C}}M = T^{1,0}M \oplus T^{0,1}M$ ,  $\nabla = \nabla_J^{1,0} + \nabla_J^{0,1}$

$$\Gamma_J(M, \mathcal{L}) = \{ \psi \in \Gamma(M, \mathcal{L}) : \nabla_J^{0,1} \psi = 0 \}$$

$\mathcal{H}_J = L^2$ -completion of  $\Gamma_J(M, \mathcal{L})$ .

Problem How does  $\mathcal{H}_J$  depend on  $J \in \mathcal{J}$  ?

## Axelrod - Della Pietra - Witten (JDG, 1991)

Hilbert space bundle  $\mathcal{H}_{\mathcal{J}}$   $\mathcal{H} \subset \mathcal{J} \times \Gamma(M, \ell)$   $\rightsquigarrow$  connection  $\nabla^{\mathcal{H}}$   
 $\downarrow$   $\downarrow$   $\downarrow$  on  $\mathcal{H} \rightarrow \mathcal{J}$ .  
 $\mathcal{J} \in \mathcal{J} = \mathcal{J}$

If  $\nabla^{\mathcal{H}}$  is projectively flat, then the concept of quantum states does not depend on the choice of  $\mathcal{J}$ .

This is so if  $(M, \omega)$  is a symplectic vector space  $(V, \omega)$  and  $\mathcal{J}$  is the set of compatible linear complex structures.

Kirwin - W. (CMP, 2006)

$\mathcal{J}$  is identified as Siegel's upper-half space  $Sp(2n, \mathbb{R})/U(n)$

which is a non-compact Hermitian symmetric space,  
or a classical domain in  $\mathbb{C}^{\frac{1}{2}n(n+1)}$  (Siegel, 1943)

If  $n=1$ ,  $\mathcal{J}$  is the upper-half plane or the unit disc in  $\mathbb{C}$ .

$$\text{curv}(\nabla^{\mathcal{H}}) = -\frac{1}{4} \text{tr}_{V_{\mathcal{J}}^{\mathcal{H}}} (S\mathcal{J} \wedge S\mathcal{J}) \sim \text{standard Kähler form}$$

Kirwin-W. (2006) continued

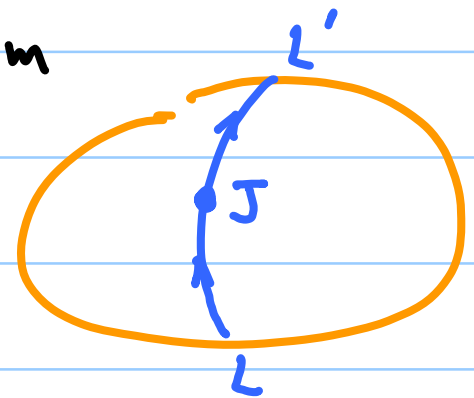
the Shiho boundary of  $\mathcal{J}$  is  $\{L: \text{real Lagrangian subspace}\}$

$\mathcal{H}_2 = \{L^\circ\text{-functions on } V/L\}$  in the real polarisation  $L$ .

parallel transport along the geodesics in  $\mathcal{J}$

from  $L$  to  $\mathcal{J}$  : Segal-Bargmann transform

from  $L$  to  $L'$  : Fourier transform



Generalisations: toric manifold

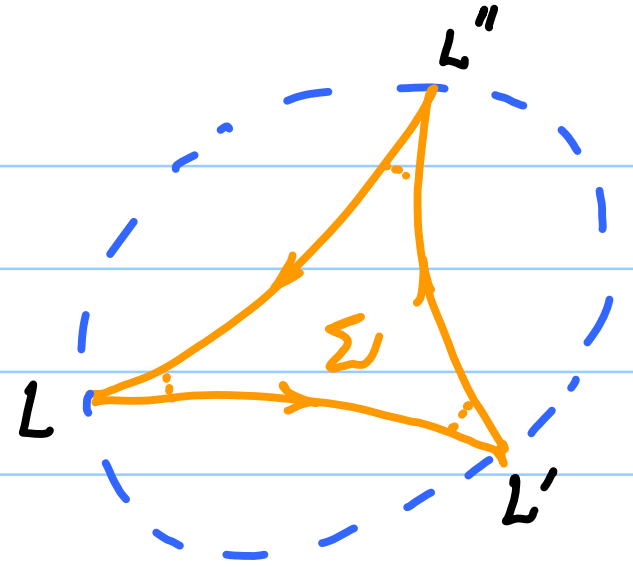
Kirwin-Mourão-Nunes (2013)



Given  $L, L', L'' \in \mathcal{L}$ , mutually transverse.

$$F_{LL''} \circ F_{L''L'} \circ F_{L'L} = e^{\frac{i}{2} \int_{\Sigma} \sigma} \text{id}_{\mathcal{H}_L}$$

(holonomy of proj flat conn.)



Compare **Lion-Vergne (1980)**

$$F_{LL''} \circ F_{L''L'} \circ F_{L'L} = e^{\frac{i}{4} \pi \tau(L, L', L'')} \text{id}_{\mathcal{H}_L}$$

$$\tau(L, L', L'') = \frac{1}{2\pi} \int_{\Sigma} \sigma$$

W. (2011)

Deformation quantisation Bayen-Flato-Frønsdal-Lichnerowicz - Sternheimer (1978)

non-commutative but associative star product

$$f * g = fg + c_1(f, g)\hbar + c_2(f, g)\hbar^2 + \dots$$

each  $c_k(f, g)$  is a differential polynomial of  $f, g$ ,  $c_1(f, g) - c_1(g, f) = \{f, g\}$ .

$1 * f = f * 1 = f$  and  $(f * g) * h = f * (g * h)$ .

existence of formal star product on symplectic manifolds (de Wilde-Lecante, 1983)

Poisson manifolds (Kontsevich, 2003)

Problem Is the formal power series in  $\hbar$  convergent?

Yes if  $(M, \omega)$  is a symplectic vector space  $(V, \omega)$

Moyal product (1949)  $f *_0 g = f e^{\frac{\hbar}{2} \pi^{\mu\nu} \overleftarrow{\partial}_\mu \overrightarrow{\partial}_\nu} g$

converges for "reasonable" functions  $f, g$ .

e.g. Omori-Maeda-Miyazaki-Yoshioka (1999)

$$\Sigma_p(V) = \left\{ f : \|f\|_{p,s} = \sup_{x \in V} |f(x)| e^{-s|x|^2} < \infty \right\} \quad \text{Fréchet space}$$

If  $0 < p \leq 2$ .  $*_0 : \Sigma_p(V) \times \Sigma_p(V) \rightarrow \Sigma_p(V)$  is non-formal

They also introduced  $f *_K g = f e^{\frac{\hbar K}{2} (\pi + K)^{\mu\nu} \overleftarrow{\partial}_\mu \overrightarrow{\partial}_\nu} g$ ,  $\forall K \in S^2 V^{\otimes 2}$ .

associative, non-formal

$$\mathcal{O} = S^2 V^E \times \Sigma_p(V)$$

$$\downarrow$$

$$S^2 V^E$$

$$\begin{array}{ccc} \Sigma_p(V) & \xrightarrow{I_{K,K'}} & \Sigma_p(V) \\ \downarrow & & \downarrow \\ K & & K' \end{array}$$

intertwining operator  $I_{K,K'} f = e^{\int_K^{K'} (K'-K)^\mu \partial_\mu} f$ .

$$I_{K'}(f *_K g) = (I_{K'} f) *_K (I_{K'} g).$$

$\leadsto \exists$  flat connection  $\nabla^{\mathcal{O}}$  on  $\mathcal{O} \rightarrow S^2 V^E$  such that the parallel transport from  $K$  to  $K'$  is  $I_{K,K'}$ .

W. - Yoshioka (to appear)

$$J \mapsto S^2 V^e$$

$$J \mapsto K_J := \frac{E}{2} (J \otimes 1 - 1 \otimes J) \pi$$

$*_J = *_{K_J}$  is the normal order star product w.r.t.  $J$ .

$$\frac{1}{2}(\pi + K_J) = \left( \frac{1 + FJ}{2} \otimes \frac{1 - FJ}{2} \right) \pi$$

$$f *_J g = f e^{Fh \pi^{ij} \overleftarrow{\partial}_i \overrightarrow{\partial}_j} g, \text{ where } i, j \text{ are indices in } V_J^{\leq 0}.$$

Now  $f \in \mathcal{O}_J$ ,  $\psi \in \mathcal{H}_J$ , define

$$f *_J \psi = f e^{\frac{Fh}{2} (\pi + K_J)^{uv} \overleftarrow{\partial}_u \overrightarrow{\partial}_v} \psi = f e^{Fh \pi^{ij} \overleftarrow{\partial}_i \overrightarrow{\partial}_j} \psi$$

Note:  $\nabla_{\overleftarrow{i}} \psi = 0$ ,  $[\nabla_j, \nabla_k] = 0$ .

## Properties

prequantum action

1. non-formal,  $f *_J \psi = f \psi - \sqrt{\hbar} \nabla_{H_f} \psi + O(\hbar^2)$

2. If  $\nabla_J^{\circ,1} \psi = 0$ , then  $\nabla_J^{\circ,1} (f *_J \psi) = 0$

3. associativity:  $(f *_J g) *_J \psi = f *_J (g *_J \psi)$ .

4. for any path  $\{J_t\}_{t \in [0,1]}$  from  $J_0$  to  $J_1$  in  $\mathcal{J}$ ,

$$U_{\{J_t\}} (f *_J \psi) = (I_{J_0}^{J_1} f) *_J (U_{\{J_t\}} \psi),$$

where  $U_{\{J_t\}}$  is the parallel transport under  $\nabla^{\mathcal{H}}$ .

proof of 2:

$$\nabla_{\vec{e}_i} (f *_J \psi) = \dots$$

$$= f *_J \left( \cancel{\vec{e}_i} + \sqrt{\hbar} \pi^{jk} \cancel{\vec{e}_j} \underbrace{[\vec{e}_i, \vec{e}_k]}_{\omega_{ik}/\sqrt{\hbar}} + \vec{e}_i \right) \psi$$

$$= 0$$

## Fermionic systems

phase space: Euclidean space  $V$   
 $C^\infty(V)$  or  $S(V^*) \rightsquigarrow \wedge V^*$

Weyl algebra  $\rightsquigarrow$  Clifford algebra

Geometric quantisation  $\left\{ \begin{array}{l} \text{prequantisation} \\ \text{quantisation} \\ \text{projectively flat connection} \end{array} \right. \begin{array}{l} \text{Kostant (1975)} \\ \text{Woodhouse (1981)} \\ \text{W. (2015)} \end{array}$

$$\mathcal{H} \downarrow \mathcal{J} \cong \text{SO}(m) / \text{U}\left(\left[\frac{m}{2}\right]\right)$$

Compact Hermitian symmetric space  
L-K Hua (1944)

$\forall J \in \mathcal{J}$ , (pre) quantisation  $\Rightarrow$

$\mathcal{H}_J = (\text{Spinor repr.}) \text{ fermionic gaussian} \subset \Lambda^0 V_{\mathbb{C}}^*$ .

They form a vect bundle  $\mathcal{H} \rightarrow \mathcal{J}$  with a projectively flat connection W. (2011, 2012)

- holonomy  $\sim$  fermionic analog of triple Maslov index introduced by Magneron (1984, 1988)
- the connection becomes flat upon metaplectic correction



## Deformation quantisation

- fermionic star product on  $\Lambda^* V_{\mathbb{C}}$

Hirshfeld, Henselder (2003, 2005)

- star product of  $\Lambda^* V_{\mathbb{C}}$  on  $\mathbb{H}_J$

- flat and proj'ly flat connections over  $\mathcal{J}$

W. (to appear)