# SQG in Bounded Domains

#### Peter Constantin

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### Collaborators

Mihaela Ignatova, (Princeton)

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Huy Nguyen (Princeton)



Active scalar

$$\partial_t \theta + \boldsymbol{u} \cdot \nabla \theta = \boldsymbol{0}$$

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 $\nabla^{\perp}\theta$  like vorticity in 3D Euler: level sets of theta are carried by the flow, tangent field stretched:

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Blow up problem open: 3D Euler, 2D SQG, 2D Boussinesq, 2D incompressible porous medium, 2D Oldroyd B. Similar.

# Numerical results

SQG- geophysical origin: Charney.



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Diego Cordoba: no blow up, under assumption of hyperbolic saddle. C-Lai-Sharma-Tseng-Wu. Parallel computation, cluster of 128 machines, well resolved for long time. Same initial data.



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$$\partial_t \theta + \boldsymbol{u} \cdot \nabla \theta = \boldsymbol{0}, \quad \boldsymbol{u} = \boldsymbol{R}^{\perp} \theta.$$

For periodic  $\theta = \sum_{j \in \mathbb{Z}^2} \widehat{\theta}(j) e^{i(j \cdot x)}$ , infinite ODE

$$\frac{d\theta}{dt} = N(\theta, \theta).$$

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Weak continuity:

 $\begin{array}{l} \| (-\Delta)^{-1} \left[ \textit{\textit{N}}(\theta_1, \theta_1) - \textit{\textit{N}}(\theta_2, \theta_2) \right] \|_{\textit{\textit{W}}} \leq \\ \textit{\textit{C}} \left\{ \| \theta_1 - \theta_2 \|_{\textit{\textit{W}}} \left( 1 + \log_+ \| \theta_1 - \theta_2 \|_{\textit{\textit{W}}} \right) \right\} (\| \theta_1 \|_{\textit{L}^2} + \| \theta_2 \|_{\textit{L}^2}) \end{array}$ 

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with  $||f||_w = \sup_{j \in \mathbb{Z}^2} |\hat{f}(j)|$ . Quasi-Lipschitz, with loss of two derivatives. A commutator structure.

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- ► quasilinear, critical in the sense of Goldilocks: easy for A<sup>s</sup>, s > 1, hard for s < 1.)</p>

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Regularity and uniqueness: with critical dissipation: Cordoba-Wu-C = small data in  $L^{\infty}$ .

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for  $q^2 = |\nabla^{\perp} \theta|^2$ , with

 $\boldsymbol{Q} = (\nabla \boldsymbol{u}) \nabla^{\perp} \boldsymbol{\theta} \cdot \nabla^{\perp} \boldsymbol{\theta} \leq |\nabla \boldsymbol{u}| \boldsymbol{q}^{2}.$ 

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 $|\nabla u| \sim q$ : *Q* is cubic. Nonlinear lower bound ! (Vicol, C)

$${\mathcal D}(q) = q \Lambda q - rac{1}{2} \Lambda \left( q^2 
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- Smallness of α: The term corresponding to Q in the finite difference version of the argument has a small (α) prefactor and it is dominated by the term corresponding to D(q)

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- Global existence of solutions for critical dissipative SQG: global interior Lipschitz bounds (Ignatova, C)

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- Commutator estimates (Ignatova, C, and H.Q. Nguen, C)
- Global existence of solutions for critical dissipative SQG: global interior Lipschitz bounds (Ignatova, C)

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▶ Global L<sup>2</sup> weak solutions for inviscid SQG (H.Q. Nguyen, C)

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#### Theorem

(*C*, Ignatova) Let  $\theta(x, t)$  be a smooth solution of critical SQG in the smooth bounded domain  $\Omega$ . There exists  $0 < \alpha < 1$  depending only on  $\|\theta_0\|_{L^{\infty}(\Omega)}$  and  $\Omega$ , and a constant  $\Gamma > 0$  depending only on the domain  $\Omega$  (in particular: not on T) such that

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 $\sup_{0 \leq t < T} \|\theta(t)\|_{\mathcal{C}^{\alpha}(\Omega)} \leq \Gamma \|\theta_0\|_{\mathcal{C}^{\alpha}(\Omega)}.$ 

Moreover,

$$\sup_{x \in \Omega, 0 \le t < T} d(x) |\nabla_x \theta(x, t)| \le \Gamma_1 \left[ \sup_{x \in \Omega} d(x) |\nabla_x \theta_0(x)| + \mathcal{P} \left( \|\theta_0\|_{L^{\infty}(\Omega)} \right) \right]$$

$$\begin{split} [f]_{\alpha} &= \sup_{x \in \Omega} (d(x))^{\alpha} \left( \sup_{h \neq 0, |h| < d(x)} \frac{|f(x+h) - f(x)|}{|h|^{\alpha}} \right) < \infty. \\ d(x) &= dist(x, \partial \Omega)). \text{ Norm in } C^{\alpha}(\Omega) \text{ (interior)} \\ &\|f\|_{C^{\alpha}} = \|f\|_{L^{\infty}(\Omega)} + [f]_{\alpha}. \end{split}$$

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 Gaussian bounds for heat kernel; cancellation due to translation invariance effective for small time.

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- Finite difference bounds for Riesz transforms using the nonlinear max principle bound in its finite difference variant.
# Basics in bounded domains

- $\Omega \subset \mathbb{R}^d$  open, bounded, smooth boundary
- ► -△ Laplacian operator with homogeneous Dirichlet boundary conditions
- w<sub>j</sub> are L<sup>2</sup>(Ω) normalized eigenfunctions, λ<sub>j</sub> corresponding eigenvalues counted with their multiplicities

$$-\Delta w_j = \lambda_j w_j$$

 $\blacktriangleright 0 < \lambda_1 \leq \cdots \leq \lambda_j \to \infty$ 

►  $-\Delta$  positive self-adjoint operator in  $L^2$  with domain  $\mathcal{D}(-\Delta) = H^2(\Omega) \cap H_0^1(\Omega)$ 

The ground state is positive and

 $c_0d(x) \leq w_1(x) \leq C_0d(x)$ 

for all  $x \in \Omega$ , where

 $d(x) = dist(x, \partial \Omega)$ 

# Fractional powers in terms of heat kernel

$$(-\Delta)^{\alpha}f = \sum_{j=1}^{\infty} \lambda_j^{\alpha}f_j w_j$$

 $f_j = \int_{\Omega} f(y) w_j(y) \, dy$ 

$$\Lambda_D = (-\Delta)^{\frac{1}{2}}$$

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$$\min\left(\frac{w_{1}(x)}{|x-y|},1\right)\min\left(\frac{w_{1}(y)}{|x-y|},1\right)t^{-\frac{d}{2}}e^{-\frac{|x-y|^{2}}{Kt}} \leq H_{D}(t,x,y) \\ \leq C\min\left(\frac{w_{1}(x)}{|x-y|},1\right)\min\left(\frac{w_{1}(y)}{|x-y|},1\right)t^{-\frac{d}{2}}e^{-\frac{|x-y|^{2}}{Kt}}$$

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$$\frac{|\nabla_x H_D(t, x, y)|}{H_D(t, x, y)} \le C \begin{cases} \frac{1}{d(x)}, & \text{if } \sqrt{t} \ge d(x), \\ \frac{1}{\sqrt{t}} \left(1 + \frac{|x-y|}{\sqrt{t}}\right), & \text{if } \sqrt{t} \le d(x) \end{cases}$$

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holds for all  $0 \le t \le T$ . Interchange x and y:

$$\partial_1^{\beta} H_D(t, y, x) = \partial_2^{\beta} H_D(t, x, y) = \sum_{j=1}^{\infty} e^{-t\lambda_j} \partial_y^{\beta} w_j(y) w_j(x).$$

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$$|\nabla_x \nabla_x H_D(x, y, t)| \le C t^{-1-\frac{d}{2}} e^{-\frac{|x-y|^2}{Rt}}$$

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holds for  $t \leq cd(x)^2$  and  $0 < t \leq T$ .

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### Proposition

(*C*, Ignatova) Let  $\Omega$  be a bounded domain with smooth boundary, let 0 < s < 2. There exists a constant *C* depending on the domain and on *s* such that for every  $\Phi$ , a  $C^2$  convex function satisfying  $\Phi(0) = 0$ , and every  $f \in C_0^{\infty}(\Omega)$ 

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$$D(f) = f\Lambda_D f - \frac{1}{2}\Lambda_D f^2 \ge \frac{C}{d(x)}f^2(x)$$

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Dramatically different from  $\mathbb{R}^d$ !

# The nonlinear bound for derivatives

Theorem (*C*, *I*) Let  $f \in L^{\infty}(\Omega) \cap \mathcal{D}(\Lambda_D^s)$ ,  $0 \le s < 2$ . Assume that  $f = \partial q$  with  $q \in L^{\infty}(\Omega)$  and  $\partial$  a first order derivative. Then there exist constants *c*, *C* depending on  $\Omega$  and *s* such that

$$f\Lambda_D^s f - \frac{1}{2}\Lambda_D^s f^2 \ge c \|q\|_{L^\infty}^{-s} |f_d|^{2+s}$$

holds pointwise in  $\Omega$ , with

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Proof: nontrivial, uses precise bounds on the heat kernel and

$$f\Lambda_D^s f - \frac{1}{2}\Lambda_D^s f^2 \ge \frac{c_s}{2}\int_0^\infty t^{-1-\frac{s}{2}} dt \int_\Omega H_D(t,x,y)(f(x)-f(y))^2 dy$$

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## Good cutoff

#### Lemma

(C,I) Let  $\Omega$  be a bounded domain with  $C^2$  boundary. For  $\ell > 0$  small enough (depending on  $\Omega$ ) there exist cutoff functions  $\chi_{\ell} = \chi$  with the properties:  $0 \le \chi \le 1$ ,  $\chi(y) = 0$  if  $d(y) \le \frac{\ell}{4}$ ,  $\chi(y) = 1$  for  $d(y) \ge \frac{\ell}{2}$ ,  $|\nabla^k \chi| \le C\ell^{-k}$  with C independent of  $\ell$  and

$$\int_\Omega rac{(1-\chi(y))}{|x-y|^{d+j}} dy \leq C rac{1}{d(x)^j}$$

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hold for  $j \ge 0$  and  $d(x) \ge \ell$ . We will refer to such  $\chi$  as a "good cutoff".

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$$\int_{\Omega} |\nabla \chi(y)| \frac{1}{|x-y|^d} \leq C \frac{1}{d(x)}$$

hold for  $j \ge 0$  and  $d(x) \ge \ell$ . We will refer to such  $\chi$  as a "good cutoff". Useful because of the Gaussian bounds on the heat kernel. Makes work in  $\Omega$  look like work in half-space, where  $\chi_{\ell} = \chi_1(\frac{X_d}{\ell})$ , without changing coordinates.

## Nonlinear bound, finite differences

Theorem (C,I) Let  $\Omega$  be a bounded domain with smooth boundary. Let  $\chi \in C_0^{\infty}(\Omega)$  be a good cutoff with scale  $\ell > 0$  and let

$$f(x) = \chi(x)(\delta_h q(x)) = \chi(x)(q(x+h) - q(x))$$

with  $q \in L^{\infty}(\Omega) \cap H^1_0(\Omega)$ . Then

$$D(f)(x) = (f\Lambda_D f)(x) - \frac{1}{2}(\Lambda_D f^2)(x) \ge \gamma_1 |h|^{-1} \frac{|f_d(x)|^3}{\|q\|_{L^{\infty}}} + \gamma_1 \frac{f^2(x)}{d(x)}$$

holds a.e. pointwise in  $\Omega$  when  $|h| \leq \frac{\ell}{16}$ , and  $d(x) \geq \ell$  with

$$|f_d(x)| = |f(x)|,$$
 if  $|f(x)| \ge M ||q||_{L^{\infty}(\Omega)} \frac{|h|}{d(x)}.$ 

## Commutator

Let  $\chi$  be a good cutoff.

### Lemma

(C,I) There exists a constant  $\Gamma_0$  such that the commutator

 $C_h(\theta) = \chi \delta_h \Lambda_D \theta - \Lambda_D(\chi \delta_h \theta)$ 

obeys

$$|C_h( heta)(x)| \leq \Gamma_0 rac{|h|}{d(x)^2} \| heta\|_{L^\infty(\Omega)}$$

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for  $d(x) \geq \ell$ ,  $|h| \leq \frac{\ell}{16}$ .

# Finite difference of Riesz transform

Lemma (C,I) Let  $\chi$  be a good cutoff, and let u be defined by

 $u = R_D^{\perp} \theta.$ 

Then

$$|\delta_h u(x)| \leq C \left( \sqrt{\rho D(f)(x)} + \|\theta\|_{L^{\infty}} \left( \frac{|h|}{d(x)} + \frac{|h|}{\rho} \right) + |\delta_h \theta(x)| \right)$$

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holds for  $d(x) \ge \ell$ ,  $\rho \le cd(x)$ ,  $f = \chi \delta_h \theta$  and with *C* a constant depending on  $\Omega$ .

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This gives a bound on  $|h|^{-1}|\delta_h u(x)|$  which costs D(f).

# Idea of proof of Hölder bound

Good cutoff, and equation for  $\delta_h \theta$  imply:

$$\frac{1}{2}L_{\chi}(\delta_{h}\theta)^{2}+D(f)+(\delta_{h}\theta)C_{h}(\theta)=0$$

with

$$\mathcal{L}_{\chi}g = \partial_t g + u \cdot \nabla_x g + \delta_h u \cdot \nabla_h g + \Lambda_D(\chi^2 g).$$

and

$$D(f) \ge \gamma_1 |h|^{-1} ||\theta||_{L^{\infty}}^{-1} |(\delta_h \theta)_d|^3 + \gamma_1 (d(x))^{-1} |\delta_h \theta|^2$$

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### Idea of proof of Hölder bound

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$$L_{\chi}\left(\frac{\delta_{h}\theta(x)^{2}}{|h|^{2\alpha}}\right) + \frac{\gamma_{1}}{4d(x)}\left(\frac{\delta_{h}\theta(x)^{2}}{|h|^{2\alpha}} - \Gamma_{1}\ell^{-2\alpha}\|\theta\|_{L^{\infty}}^{2}\right) \leq 0$$

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# Inviscid global weak solutions, bounded domains

Theorem (*C*, *Q.H.* Nguyen.) Let  $\theta_0 \in L^2(\Omega)$ . There exists a weak solution of inviscid SQG  $\partial_t \theta + R_D^{\perp} \theta \cdot \nabla \theta = 0$ with  $\psi = \Lambda_D^{-1} \theta \in C([0,\infty), H_0^{1-\epsilon}(\Omega))$  for any  $0 < \epsilon < 1$ . The Hamiltonian  $\int_{\Omega} \theta(t) \Lambda_D^{-1} \theta(t) dx$ 

is conserved in time, and the  $L^2(\Omega)$  norm of  $\theta(t)$  is nonincreasing in time.

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#### Theorem

(*C*, Ignatova, Nguyen) Let T > 0 and let  $\theta_k(x, t)$ ,  $0 \le t \le T$  be a sequence of solutions of critical SQG with "viscosities"  $\nu_k \to 0$  and initial data uniformly bounded in  $L^2(\Omega)$ . Then the limit of any weakly  $L^2$  convergent subsequence is a weak solution of inviscid SQG.

### **Elements of Proof**

Weak continuity from commutator structure (adapted for bounded domains):  $\phi$  test function,  $\psi = \Lambda_D^{-1} \theta$ :

 $\int_{\Omega} (\mathbf{R}_{D}^{\perp} \theta \cdot \nabla \theta) \phi dx$  $= -\frac{1}{2} \int_{\Omega} \psi[\Lambda_{D}, \nabla^{\perp}] \psi \cdot \nabla \phi dx + \frac{1}{2} \int_{\Omega} \nabla^{\perp} \psi \cdot [\Lambda_{D}, \nabla \phi] \psi dx$ 

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Together with commutator estimates

### Theorem

(Ignatova, C) Let  $\chi \in B(\Omega)$  with  $B(\Omega) = W^{2,\infty}(\Omega) \cap W^{1,\infty}(\Omega)$  if  $d \ge 3$ , and  $B(\Omega) = W^{2,p}(\Omega)$  with p > 2 if d = 2. There exists a constant  $C = C(d, p, \Omega)$  such that

 $\|\boldsymbol{\Lambda}_D^{\frac{1}{2}}[\boldsymbol{\Lambda}_D, \boldsymbol{\chi}]\boldsymbol{\psi}\|_{L^2(\Omega)} \leq \boldsymbol{C} \|\boldsymbol{\chi}\|_{\boldsymbol{B}(\Omega)} \|\boldsymbol{\Lambda}_D^{\frac{1}{2}}\boldsymbol{\psi}\|_{L^2(\Omega)}.$ 

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#### Theorem

(Ignatova, Nguyen, C.) For  $1 \le p \le \infty, \, 0 < s < 2,$  there exists C such that for all  $x \in \Omega$ 

$$\|[\Lambda^s_D, 
abla]\psi(x)\| \leq Cd(x)^{-1-s-rac{d}{p}}\|\psi\|_{L^p(\Omega)}$$

# **Conclusions and Outlook**

Global interior regularity a priori bounds for critical SQG
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- Inviscid limit: Any L<sup>2</sup> weak limit of decent regularizations of SQG converge to weak solutions of SQG.

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 Construction of global interior Lipschitz solutions: good approximations

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$$u_{ au}(x,t) = 
abla^{\perp} \int_{ au}^{\infty} s^{-rac{1}{2}} e^{s\Delta} heta(x,t) ds, \quad au > 0.$$

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In preparation.

- Uniqueness ?
- Uniform gradient bounds up to the boundary

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$$u_{\tau}(x,t) = 
abla^{\perp} \int_{ au}^{\infty} s^{-rac{1}{2}} e^{s\Delta} heta(x,t) ds, \quad au > 0.$$

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