

Manifestations of localization in the random XXZ quantum spin chain

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We have $\sigma(H_\omega) = \{0\} \cup \left[1 - \frac{1}{\Delta}, \infty\right)$ almost surely.

XXZ chain Hamiltonian in finite intervals

Consider the finite interval $[-L, L] = [-L, L] \cap \mathbb{Z}$, $L \in \mathbb{N}$, and set

$$H_{\omega}^{(L)} = \sum_{i=-L}^{L-1} \left\{ \frac{1}{4} (I - \sigma_i^z \sigma_{i+1}^z) - \frac{1}{4\Delta} (\sigma_i^x \sigma_{i+1}^x + \sigma_i^y \sigma_{i+1}^y) \right\} + \lambda \sum_{i=-L}^L \omega_i \mathcal{N}_i$$

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- Unique ground state $\psi_0 = \psi_0^{(L)}$ determined by $\mathcal{N}_i \psi_0 = 0$ for all i .
- The spectrum of $H^{(L)} = H_{\omega}^{(L)}$ is almost surely simple, so that its normalized eigenvectors can be labeled as ψ_E , $E \in \sigma(H^{(L)})$.

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$$G_I = \{g : \mathbb{R} \rightarrow \mathbb{C} \text{ Borel measurable, } |g| \leq \chi_I\}.$$

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We will say that we have droplet localization in an interval I if the conclusions of the theorem hold in the interval I .

Best possible interval for droplet localization

We proved droplet localization on intervals

$$I_{1,\delta} = \left[1 - \frac{1}{\Delta}, (2 - \delta)\left(1 - \frac{1}{\Delta}\right)\right] \subset \left[1 - \frac{1}{\Delta}, 2\left(1 - \frac{1}{\Delta}\right)\right).$$

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- Results on the random XXZ quantum spin chain by Beaud and Warzel (2017).

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◆ $H = H_\omega$ will be a random XXZ spin chain satisfying droplet localization in the interval $I = I_{1,\delta} = \left[1 - \frac{1}{\Delta}, (2 - \delta)\left(1 - \frac{1}{\Delta}\right)\right]$.

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- ◆ Given a local observable X , we will generally specify a support for X , denoted by $\mathcal{S}_X = [s_X, r_X]$. We always assume $\emptyset \neq \mathcal{S}_X \subset [-L, L]$.

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- ◆ $H = H_\omega$ will be a random XXZ spin chain satisfying droplet localization in the interval $I = I_{1,\delta} = \left[1 - \frac{1}{\Delta}, (2 - \delta)(1 - \frac{1}{\Delta})\right]$.
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The time evolution of a local observable X under $H^{(L)}$ is given by

$$\tau_t(X) = \tau_t^{(L)}(X) = e^{itH^{(L)}} X e^{-itH^{(L)}} \quad \text{for } t \in \mathbb{R}.$$

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Thus, given an energy interval J , we consider the sub-Hilbert space $\text{Ran } P_J^{(L)}$, spanned by the the eigenstates of $H^{(L)}$ with energies in J , and localize an observable X in the energy interval J by considering its restriction to $\text{Ran } P_J^{(L)}$,

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Clearly

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$X_l = (X_{l_0})_l \implies$ the theorem holds with l substituted for l_0 .

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Moreover, the estimate (1) is not true without the counterterms.

Correlators

We define the truncated time evolution of an observable X in the energy window I by ($H = H_\omega^{(L)}$),

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We are interested in quantities of the form ($K \subset I$)

$$R_K(\tau_t^I(X), Y) = (\tau_t^I(X) Y)_K - (\tau_t^I(X))_K Y_K = (\tau_t^I(X) Y)_K - \tau_t(X_K) Y_K.$$

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While it is obvious where the first counterterm comes from, the same is not true of the second, where the time evolution seems to sit in the *wrong* place: it is $\tau_t^K(Y)$ and not $\tau_t^K(X)$. It turns out this term encodes information about the states above the energy window K , and the appearance of $\tau_t^K(Y)$ is related to the reduction of this data to P_0 .

Particle number conservation

An important property of the XXZ chain is particle number conservation:

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It follows that

$$H^{(L)} = \bigoplus_{N=0}^{2L+1} H_N^{(L)} \quad \text{with respect to} \quad \mathcal{H}^{(L)} = \bigoplus_{N=0}^{2L+1} \mathcal{H}_N^{(L)}.$$

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$\mathcal{H}_N^{(L)} \cong \ell^2(\mathcal{X}_N^{(L)})$ ($x_1 < x_2 < \dots < x_N$ are the sites with down spins)

$$H_N^{(L)} \cong -\frac{1}{2\Delta} \mathcal{L}_N^{(L)} + \left(1 - \frac{1}{\Delta}\right) \widetilde{W} + \lambda V_\omega + \left(\beta - \frac{1}{2}\left(1 - \frac{1}{\Delta}\right)\right) \chi^{(L)}.$$

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- $\chi^{(L)} = \chi_{-L} + \chi_L$, the left and right boundary terms.

Local number operators

Recall $\mathcal{N}_i = \begin{cases} 1 & \text{if the spin at site } i \text{ is down} \\ 0 & \text{otherwise} \end{cases}$, and $\mathcal{N}_i = \bigoplus_{N=0}^{2L+1} \mathcal{N}_i^{(N)}$.

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Thus $\mathcal{N}_i^{(N)} \cong Q_i^{(N)}$, where $Q_i^{(N)}$ is the characteristic function of the set

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It follows that

$$\sum_{E \in \sigma(H^{(L)}) \cap I} \|\mathcal{N}_i \psi_E\| \|\mathcal{N}_j \psi_E\| = \sum_{N=1}^{\infty} Q_N^{(L)}(i, j; I) \quad \text{almost surely.}$$

Reformulation of droplet localization

Theorem

Fix $0 < \delta < 1$, and let $I_{1,\delta} = \left[1 - \frac{1}{\Delta}, (2 - \delta)\left(1 - \frac{1}{\Delta}\right)\right]$.

There exists a constant $K > 0$ with the following property: If

$$\lambda\sqrt{\Delta - 1} \min\{1, (\Delta - 1)\} \geq K,$$

there exist constants $C < \infty$ and $m > 0$ such that

$$\sum_{N=1}^{\infty} \mathbb{E}(Q_N^{(L)}(i, j; I_{1,\delta})) \leq Ce^{-m|i-j|} \quad \text{for all } -L \leq i, j \leq L,$$

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This reformulation reduces the proof of droplet localization in the droplet spectrum to establishing decay properties of the Green's functions associated with the random Schrödinger operators $H_N^{(L)}$

Strategy for the proof of the theorem

The analysis is first done separately along the edge

$$\mathcal{X}_{N,1} = \{x \in \mathcal{X}_N : \widetilde{W}(x) = 1\} = \{x = (x_1, x_1+1, \dots, x_1+N-1) : x_1 \in \mathbb{Z}\}$$

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- In the bulk we use (purely deterministic) **Combes-Thomas-type estimates**.
- Along the edge we establish a **fractional moment estimate**.
- These estimates are combined to derive **localization on a pair of “boxes”**, as in an energy interval multiscale analysis, from which we derive droplet localization.

Combes-Thomas-type estimates in the bulk

Theorem

Let $z \notin \sigma(H_N^{(L)})$ and let

$$\left\| \widetilde{W}^{\frac{1}{2}} (H_N^{(L)} - z)^{-1} \widetilde{W}^{\frac{1}{2}} \right\| \leq \frac{1}{\eta_z}.$$

Then for all $\Phi, \Psi \in \mathcal{X}_N^{(L)}$ we have

$$\left\| \chi_\Phi \widetilde{W}^{\frac{1}{2}} (H_N^{(L)} - z)^{-1} \widetilde{W}^{\frac{1}{2}} \chi_\Psi \right\| \leq \frac{2}{\eta_z} e^{-\log\left(1 + \frac{\eta_z \Delta}{2}\right) \text{dist}_1(\Phi, \Psi)}.$$

(dist_1 is the distance in the $\|\cdot\|_1$ norm.)

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Note that $|u - v|_\infty = |u_1 - v_1|$ for $u, v \in \mathcal{X}_{N,1}^{(L)}$.

Decomposition of local observables

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Given a local observable X , we define projections $P_{\pm}^{(X)}$ by

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Note that $P_0 = \bigotimes_{j \in [-L, L]} (1 - \mathcal{N}_j)$, so $P_{-}^{(X)} P_0 = P_0 P_{-}^{(X)} = 0$, and $P_{-}^{(X)} \leq \sum_{i \in \mathcal{S}_X} \mathcal{N}_i$.

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Moreover, since $P_{+}^{(X)}$ is a rank one projection on $\mathcal{H}_{\mathcal{S}_X}$, we must have

$$X^{+, +} = \zeta_X P_{+}^{(X)}, \quad \text{where} \quad \zeta_X \in \mathbb{C}, \quad |\zeta_X| \leq \|X\|.$$

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In particular,

$$(X - \zeta_X)^{+, +} = 0 \quad \text{and} \quad \|X - \zeta_X\| \leq 2 \|X\|,$$

so we can assume $X^{+, +} = 0$ in the proofs.

Droplet localization for general local observables

Droplet localization is defined in terms of the local number operators \mathcal{N}_i . For proving the theorems we need to apply it to general local observables.

Lemma

Let X, Y be local observables, $\ell \geq 1$. Then

$$\mathbb{E} \left(\sup_{g \in G_{I_0}} \left\| P_-^{(X)} g(H) P_-^{(Y)} \right\|_1 \right) \leq C e^{-m \text{dist}(X, Y)}$$

$$\mathbb{E} \left(\left\| P_-^{(Y)} P_-^{(X)} P_{I_0} \right\|_1 \right) \leq C e^{-\frac{1}{2} m \text{dist}(X, Y)}$$

$$\mathbb{E} \left(\sup_{I \in G_I} \left\| P_-^{(X)} g(H) P_+^{(S_{X, \ell})} \right\|_1 \right) \leq C e^{-m \ell}$$

$$\mathbb{E} \left(\sup_{g \in G_I} \left\| P_+^{(S_{Y, \ell}^c)} g(H) P_+^{(S_{X, \ell}^c)} \right\|_1 \right) \leq C e^{-m(\text{dist}(X, Y) - 2\ell)}$$

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The commutator is estimated by the Lieb-Robinson bound for small t .



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For $E, E' \in K$ we have

$$\begin{aligned} P_E \left(\int_{\mathbb{R}} e^{-irH} Y \tau_r(X) \hat{f}(r) dr \right) P_{E'} &= P_E Y f(E + E' - H) X P_{E'} \\ &= P_E Y P_{K_f} f(E + E' - H) X P_{E'} = P_E \left(\int_{\mathbb{R}} e^{-irH} Y \{P_{K_f}\} \tau_r(X) \hat{f}(r) dr \right) P_{E'}. \end{aligned}$$

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To prove: Droplet localization in $I = \left[1 - \frac{1}{\Delta}, \Theta_1\right] \implies \Theta_1 \leq 2\left(1 - \frac{1}{\Delta}\right)$.

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Let X, Y be local observables with $X^{+,+} = Y^{+,+} = 0$. The Lemmas yield

$$\begin{aligned} \|(XP_0Y)_K\| &= \|(Xh(H)Y)_K\| \\ &\leq C \|X\| \|Y\| e^{-m_1(\text{dist}(X,Y))^{\frac{1}{2}}} + C' \sup_{r \in \mathbb{R}} \|(YP_{K_h} \tau_r(X))_K\|, \end{aligned}$$

where $K_h \subset [2\Theta_0 - \varepsilon, 2\Theta_2 + \varepsilon] \subset [\Theta_0, \Theta_1] = I$.

We can prove

$$\mathbb{E} \left(\sup_{r \in \mathbb{R}} \| (Y P_{K_h} \tau_r (X))_K \| \right) \leq C \|X\| \|Y\| e^{-\frac{1}{8} m \text{dist}(X, Y)},$$

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so we conclude that

$$\mathbb{E} (\|(X P_0 Y)_K\|) \leq C \|X\| \|Y\| e^{-m_2 (\text{dist}(X, Y))^{\frac{1}{2}}}.$$

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In particular, it follows that we have, uniformly in L ,

$$\mathbb{E} \left(\left\| \left(\sigma_i^x P_0^{(L)} \sigma_j^x \right)_K \right\| \right) \leq C e^{-m_2 (|i-j|)^{\frac{1}{2}}} \quad \text{for all } i, j \in [-L, L]. \quad (2)$$

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$$\mathbb{E} \left(\liminf_{L \rightarrow \infty} \left\| \left(\sigma_i^x P_0^{(L)} \sigma_j^x \right)_K \right\| \right) \geq \gamma_K > 0. \quad (3)$$

We can prove

$$\mathbb{E} \left(\sup_{r \in \mathbb{R}} \|(YP_{K_h} \tau_r(X))_K\| \right) \leq C \|X\| \|Y\| e^{-\frac{1}{8}m \text{dist}(X,Y)},$$

so we conclude that

$$\mathbb{E} (\|(XP_0 Y)_K\|) \leq C \|X\| \|Y\| e^{-m_2(\text{dist}(X,Y))^{\frac{1}{2}}}.$$

In particular, it follows that we have, uniformly in L ,

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(2) and (3) give a contradiction $\implies \Theta_1 \leq 2\Theta_0$.

Non-spreading of information- Sketch of proof

To prove: Given a local observables X , $t \in \mathbb{R}$ and $\ell > 0$, there is a local observable $X_\ell(t) = (X_\ell(t))_\omega$ with support $\mathcal{S}_{X,\ell}$ satisfying

$$\mathbb{E} \left(\sup_{t \in \mathbb{R}} \left\| (X_\ell(t) - \tau_t(X))_{I_0} \right\|_1 \right) \leq C \|X\| e^{-\frac{1}{16} m \ell}.$$

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We first prove that

$$\mathbb{E} \left(\sup_{t \in \mathbb{R}} \left\| \left(P_+^{(\mathcal{O})} \tau_t(X_{I_0}) P_+^{(\mathcal{O})} - \tau_t(X) \right)_{I_0} \right\|_1 \right) \leq C \|X\| e^{-\frac{1}{16} m \ell}.$$

We now observe that for all observables Z we have

$$P_+^{(\mathcal{O})} Z P_+^{(\mathcal{O})} = \tilde{Z} P_+^{(\mathcal{O})} = P_+^{(\mathcal{O})} \tilde{Z},$$

where \tilde{Z} is an observable with $\mathcal{S}_{\tilde{Z}} = \mathcal{S}_{X, \frac{\ell}{2}}$ and $\|\tilde{Z}\| \leq \|Z\|$.

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We conclude that

$$\mathbb{E} \left(\sup_{t \in \mathbb{R}} \left\| \left(P_+^{(\mathcal{O})} \widetilde{\tau_t(X)_{I_0}} - \tau_t(X) \right)_{I_0} \right\|_1 \right) \leq C \|X\| e^{-\frac{1}{16} m \ell}.$$

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Since $P_+^{(\mathcal{O})} \widetilde{\tau_t(X_{I_0})}$ does not have support in $\mathcal{S}_{X, \ell}$, we now define

$$X_\ell(t) = P_+^{(\mathcal{T})} \widetilde{\tau_t(X_{I_0})} \quad \text{for } t \in \mathbb{R},$$

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We conclude that




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Since $P_+^{(\mathcal{O})} \widetilde{\tau_t(X_{I_0})}$ does not have support in $\mathcal{S}_{X, \ell}$, we now define

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an observable with support in $\mathcal{S}_{X, \frac{\ell}{2}} \cup \mathcal{T} = \mathcal{S}_{X, \ell}$, and prove

$$\mathbb{E} \left(\sup_{t \in \mathbb{R}} \left\| \left(P_+^{(\mathcal{O})} \widetilde{\tau_t(X_{I_0})} - X_\ell(t) \right)_{I_0} \right\|_1 \right) \leq C \|X\| e^{-\frac{1}{4} m \ell}.$$

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