Modeling and simulation of deformable bodies in low-Reynolds flows

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Fluid-Structure Interaction

Fluid inside or surrounding a solid. Fluid flow \rightleftharpoons Solid structure

- $\bullet~\mbox{Fluid}$ flow $\rightarrow~\mbox{deforms}$ the solid structure
- $\bullet~Structure~deformation \rightarrow impacts the flow$

Examples:

- Airflow around an aircraft wing
- Balloon inflating
- Wind action on a yacht sail







Incompressible Navier-Stokes equations

$$\rho\left(\frac{\partial u}{\partial t} + (u \cdot \nabla)u\right) = -\mu\Delta u + \nabla p$$
$$\nabla \cdot u = 0$$

- Balance equation for linear momentum and conservation of mass equation
- 2 Linear dependence of stresses on strain rate
- Son-linearity arising from the convective term
- Seglecting inertia leads to Stokes equations

Stokes equations

$$\mu \Delta u - \nabla p = 0$$

$$\nabla \cdot u = 0$$

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Elasticity equations

$$o \frac{\partial^2 \eta}{\partial t^2} - \nabla \cdot (F\Sigma) = f_s$$

- Balance equations
- **②** Formulated in terms of the displacement $\eta(t, x)$ between reference and current configuration
- $\textcircled{O} \text{ Non-linearity depends on the elasticity model, encoded in } \Sigma$

Coupling conditions on fluid-solid interface

$$\frac{\partial \eta}{\partial t} = u$$
$$\mathbb{T}n_f + F\Sigma n_s = 0$$

- Velocity continuity
- Ormal stress continuity

Warning

Fluid and elasticity equations are not defined on the same reference frame (Eulerian vs. Lagrangian).

There are several techniques to discretize the continuous coupling conditions:

- Dirichlet-Neumann: velocity continuity (fluid side) stress continuity (solid side)
- Robin-Robin: combination of Dirichlet and Neumann conditions (fluid & solid sides)

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The choice depends on the fluid/solid system and solution technique.

Finite Element Embedded Library in C++

- \bullet Open source library \rightarrow github.com/feelpp
- Galerkin methods
- Domain Specific Embedded Language (DSEL) in C++
- Scaling : from laptops to supercomputers
- Easy deployment \rightarrow Docker, Singularity
- (Multi)Physics toolboxes: Heat Transfer, Aerothermics, CFD, CSM, FSI, Maxwell
- docs.feelpp.org
- YouTube \rightarrow Feel++ channel

Project: Cells under flow in the zebrafish







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At the microscale the Reynolds number is small and inertia forces play no role in the motion of microorganisms.

 Swimmers that perform reciprocal strokes are not able to move.



 Microorganisms have adopted non-reciprocal swimming strategies.



Magneto-elastic microswimmer at ISIR - UPMC

Composite microrobot which can perform non-reciprocal strokes.



- Viscous fluid
- Elastic body
- Magnetic head
- Driven by external magnetic field

Micro-robots can be applied for targeted drug delivery and non-invasive surgical operations.

From Navier-Stokes to Stokes equations

Incompressible Navier-Stokes equations

$$ho\Big(rac{\partial u}{\partial t} + (u\cdot
abla)u\Big) = -\mu\Delta u +
abla p$$

 $abla \cdot u = 0$

- Identify typical length (L), velocity (U), timescale (T).
- 2 Adimensionalize

$$\frac{Re}{\partial u^*} + (u^* \cdot \nabla)u^* = -\Delta u^* + \nabla p^*$$
$$\nabla \cdot u^* = 0$$

where

$$Re = rac{LU
ho}{\mu}$$
 $u^* = rac{u}{U}$ $t^* = rac{t}{LU}$ $\rho^* = rac{\rho L}{\mu U}$

For microswimmers (spermatozoa, bacteria) in water

•
$$L \approx 10^{-6} m$$

- $U \approx 30 \cdot 10^{-6} m/s$
- $\mu/\rho\approx 10^{-6}~m^2/s$

Hence, $Re \approx 10^{-5}$

Stokes equations

$$\mu \Delta u - \nabla p = 0$$
$$\nabla \cdot u = 0$$

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In the swimmer frame one can decompose the Stokes problem

$$\mu \Delta u - \nabla p = 0 \quad \text{in } \Omega \setminus S$$
$$\nabla \cdot u = 0 \quad \text{in } \Omega \setminus S$$
$$u = \omega \times x + v + w_d \quad \text{on } \partial S$$
$$u = 0 \quad \text{on } \partial \Omega \setminus \partial S$$

in 7 (3D) or 4 (2D) independent Stokes problems to extract rigid body motion.

This is possible because Stokes equations are linear in (u, p).

Solve the Stokes problem with the following boundary conditions (in the swimmer frame):

- $e_1 \times x$ produces solution u_1 (3D)
- 2 $e_2 \times x$ produces solution u_2 (3D)
- $e_3 \times x$ produces solution u_3 (3D) u_1 (2D)
- e_1 produces solution u_4 (3D) u_2 (2D)
- e_2 produces solution u_5 (3D) u_3 (2D)
- e_3 produces solution u_6 (3D)
- w_d produces solution u_7 (3D) u_4 (2D)

$$u = \sum_{i=1}^{3} \omega_i u_i + \sum_{i=4}^{6} v_i u_i + u_7 \qquad p = \sum_{i=1}^{3} \omega_i p_i + \sum_{i=4}^{6} v_i p_i + p_7$$

Newtonian fluid - Stress tensor

$$T(u,p) = -p\mathbb{I} + 2\mu \frac{
abla u +
abla u^T}{2}$$

Newton equations

$$\int_{\partial S_t} T(u, p) = 0$$
$$\int_{\partial S_t} x \times T(u, p) = 0$$

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Stokes problem decomposition (3D)

$$Mz + N = 0$$
 where $z = (\omega, v)$

$$M_{ij} = \begin{cases} \int_{\partial S_t} (x \times e_i) \cdot T(u_j, p_j) n \, \mathrm{d}x & 1 \le i \le 3, 1 \le j \le 6\\ \int_{\partial S_t} e_{i-3} \cdot T(u_j, p_j) n \, \mathrm{d}x & 4 \le i \le 6, 1 \le j \le 6 \end{cases}$$
$$N_j = \begin{cases} \int_{\partial S_t} (x \times e_i) \cdot T(u_7, p_7) n \, \mathrm{d}x & 1 \le i \le 3\\ \int_{\partial S_t} e_{i-3} \cdot T(u_7, p_7) n \, \mathrm{d}x & 4 \le i \le 6 \end{cases}$$

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Stokes problem decomposition (2D)

$$Mz + N = 0$$
 where $z = (\omega, v)$

$$M_{ij} = \begin{cases} \int_{\partial S_t} (x \times e_3) \cdot T(u_j, p_j) n \, dx & 1 \le j \le 3 \\ \int_{\partial S_t} e_i \cdot T(u_j, p_j) n \, dx & 2 \le i \le 3, 1 \le j \le 3 \end{cases}$$
$$N_j = \begin{cases} \int_{\partial S_t} (x \times e_3) \cdot T(u_4, p_4) n \, dx \\ \int_{\partial S_t} e_j \cdot T(u_4, p_4) n \, dx & 2 \le j \le 3 \end{cases}$$

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The absence of inertial effects makes the problem independent of mass.

For massless swimmers it's necessary to account for these self-propulsion constraints on the deformation.

Self-propulsion constraints

$$\int_{\partial S_t} \partial_t \eta_d(t, x) \, \mathrm{d}x = 0$$
$$\int_{\partial S_t} \partial_t \eta_d(t, x) \times \eta_d(t, x) \, \mathrm{d}x = 0$$

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Scallop theorem

A body which deforms and follows the same deformation path back in time will not show a net motion.

An example of such body is the scallop (Coquille de Saint Jacques), which moves by repeatedly opening and closing its valves.







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In the absence of inertia forces, the rigid body motion computed this way is independent of the viscosity coefficient.



It is experimentally observed that micro-swimmers are attired by boundaries. In the case of a reciprocal swimmer, the proximity to the boundary produces a net motion towards it.

In our case the scallop approaches the left corner of the box and rotates.



The previous method could be applied to spermatozoa by prescribing the velocity w_d coming from deformation. A formula coming from experiments is proposed. The tangent angle is

$$\psi(s,t) = K_0 s + 2A_0 s \cos(\omega t - \frac{2\pi s}{\lambda})$$

- Mean curvature K₀
- Varying amplitude A₀s
- Travelling wave λ, ω



- Parallelize the solution of the Stokes subproblems
- Feed displacement velocity from a FSI code
- Couple the system with equations for magnetic field
- Improve meshing and remeshing strategies

Cancer cell metastasis



Cancer cells metastasis: cells going from primary to secondary site in the host body, using the vascular system.

The zebrafish embryo: a model organism





Zebrafish (ZF) features:

- Easy to observe: thin and transparent
- Cheap to maintain
- Endless supply: Hundreds of offspring weekly
- Extremely fast development (1 day ZF \rightarrow 1 month Human)
- Complete genome sequence is known: 70% of genes shared with humans

Two kinds of cells



Red Blood Cells (RBC) in... green

• Elliptical (in embryo: $\approx 5 \times 2.5 \mu m$)

Circulating Tumor Cell (CTC or TC) in red

- Softer
- Bigger (radius $pprox 5\mu m$)

Region of interest : Caudal plexus





Tumor cells preferentially stop in the caudal plexus.



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Goals:

- Model and simulate cells in the zebrafish
- Model cell-cell and cell-wall contacts

The plan:

- Create a 2D mesh of the ZF
- Model and simulate a single cell in the flow using the level set method
- Extend to multiple cells and manage contacts

Blood flow - Stokes model

Blood in a ZF blood vessel :

- incompressible newtonian fluid
- density $ho = 1000 kg/m^3$, dynamic viscosity $\mu = 2 imes 10^{-3} Pa.s$
- mean velocity $U = 1 \times 10^{-3} m/s$
- characteristic dimension $L = 20 \times 10^{-6} m$ (ZF blood vessel)
- Reynolds number $Re = \frac{\rho UL}{\mu} = 10^{-2}$
- \Rightarrow We use the Stokes model for incompressible newtonian fluids

$$\begin{cases} -\mu\Delta u + \nabla p = F \\ \nabla \cdot u = 0 \end{cases}$$

Where u is the velocity, p is the pressure and F are the external forces.

 $\begin{array}{l} \mbox{Cell model} \rightarrow \mbox{vesicle: an inextensible membrane filled with fluid.} \\ \Rightarrow \mbox{ A cell in blood flow: an interface } \Gamma \mbox{ separating two fluids.} \end{array}$

inextensibility of the interface: $\nabla_s \cdot u = \nabla \cdot u - (\nabla u \cdot n) \cdot n = 0$ on Γ with *n* the outward normal vector.

The level set method \rightarrow Track the interface implicitly.

- \bullet define a level set function ϕ on the fluid domain Ω
- ϕ : signed distance to the interface (negative inside the cell)
- advect it using the fluid velocity field

$$\begin{cases} \frac{\partial \phi}{\partial t} + u \cdot \nabla \phi = \mathbf{0} \\ \nabla \cdot u = \mathbf{0} \end{cases}$$

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Interface related quantities

Smoothed Heaviside function

$$H_{\epsilon}(\phi) = \begin{cases} 0, & \phi \leq -\epsilon \\ \frac{1}{2} \left(1 + \frac{\phi}{\epsilon} + \frac{\sin(\frac{\pi\phi}{\epsilon})}{\pi} \right), & -\epsilon \leq \phi \leq \epsilon \\ 1, & \phi \geq \epsilon \end{cases}$$

 \rightarrow define quantities in each fluid : $\rho(\phi(x)) = \rho_2 + (\rho_1 - \rho_2)H_{\epsilon}(\phi)$. Smoothed delta function

$$\delta_{\epsilon}(\phi) = \left\{ egin{array}{cc} 0, & \phi \leq -\epsilon \ rac{1}{2\epsilon} \left(1 + \cos(rac{\pi \phi}{\epsilon})
ight), & -\epsilon \leq \phi \leq \epsilon \ 0, & \phi \geq \epsilon \end{array}
ight.$$

ightarrow define quantities on interface : $\int_{\Gamma} 1 \approx \int_{\Omega} \delta_{\epsilon}(\phi)$

Helfrich model: bending energy proportional to the square of the curvature of the membrane. In 2D :

$$E_b = \int_{\Gamma} \frac{k_B}{2} \kappa^2$$

with $k_B \approx 10^{-19} J$ for a vesicle.

Corresponding force:

$$F_{b} = \int_{\Omega} k_{B} \nabla \cdot \left[\frac{-\kappa^{2}}{2} \frac{\nabla \phi}{|\nabla \phi|} + \frac{1}{|\nabla \phi|} \left(Id - \frac{\nabla \phi \otimes \nabla \phi}{|\nabla \phi|^{2}} \right) \nabla \{ |\nabla \phi| \} \kappa \right] \delta_{\epsilon}$$

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$$\frac{D(\rho(\phi(x))u)}{Dt} - \nabla \cdot \left(\left[\mu(\phi(x))(\nabla u + (\nabla u)^{T}) \right] \right) + \nabla p = F \quad \text{in } \Omega$$
$$\nabla \cdot u = 0 \quad \text{in } \Omega$$
$$\nabla_{s} \cdot u = 0 \quad \text{on } \Gamma$$
$$\frac{\partial \phi}{\partial t} + u \cdot \nabla \phi = 0 \quad \text{in } \Omega$$
$$u = g \quad \text{on } \partial \Omega$$

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Mesh from a 2D image



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Mesh from a 2D image



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Octave script to generate geometry file:

- extract contours
- filter out duplicate points
- write geometry file for GMSH

Mesh should be:

- fine enough to preserve geometry
- fine enough to define a cell with level set method
- coarse enough to save computing time

Mesh from a 2D image



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Mesh from a 2D image



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The point: solve a physics problem without rewriting everything

How it works:

- choose your toolbox : Multifluid with built-in level set method.
- give it a mesh or a geometry
- pick and adjust the physical models
- select and tune the solver
- have fun