
Coarse-graining of collective dynamics models

Metric vs topological interactions

Pierre Degond

Imperial College London

pdegond@imperial.ac.uk

<http://sites.google.com/site/degond/>

joint works with Adrien Blanchet

Toulouse School of Economics, Toulouse, France

1. Topological interactions
2. Smooth rank-based dynamics
3. Nearest neighbor
4. Conclusion

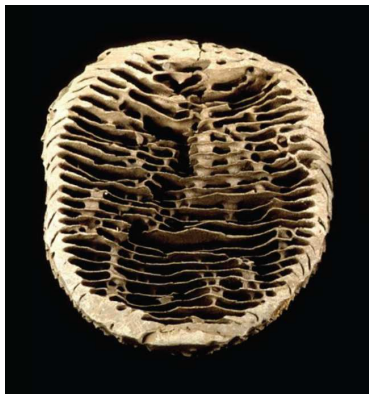
1. Topological interactions

System with locally interacting agents

Emergence of spatio-temporal coordination

Patterns, structures, correlations, synchronization

No leader



Mean-field (metric) interaction

Particle interact with all particles within a certain distance

Examples: Alignment (Vicsek ...)

Consensus (Cucker-Smale, Motsch-Tadmor, ...)

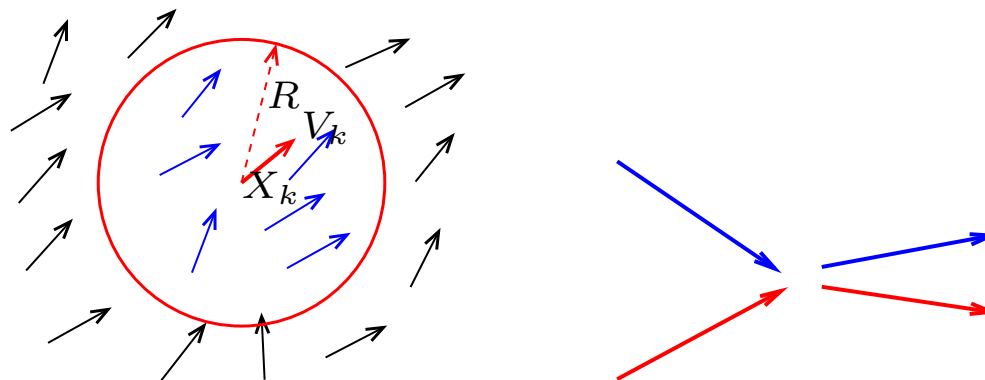
Attraction-repulsion (Bertozzi, Carrillo, ...)

Binary (ternary, ...)

Particle interacts with a partner at contact

Examples: Hard-sphere collisions (Boltzmann ...)

Alignment (Bertin, Droz & Grégoire, ...)



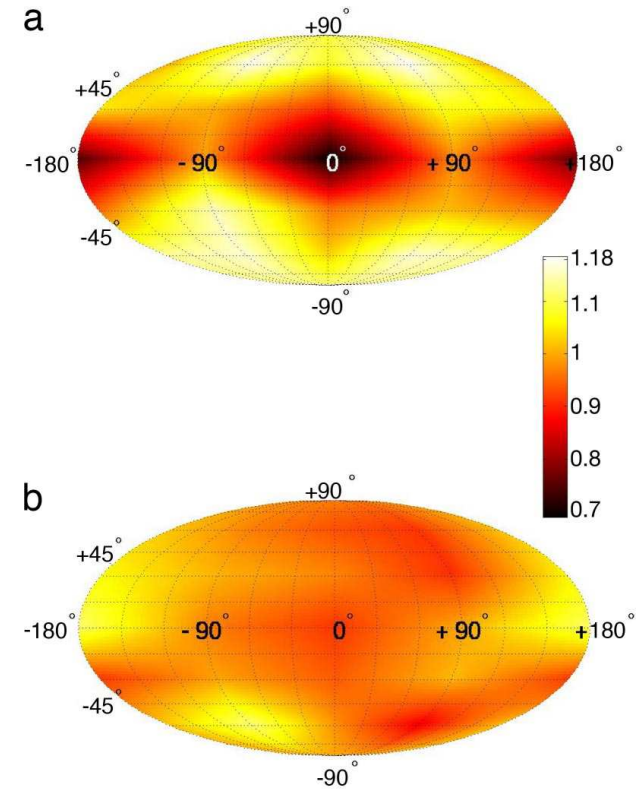
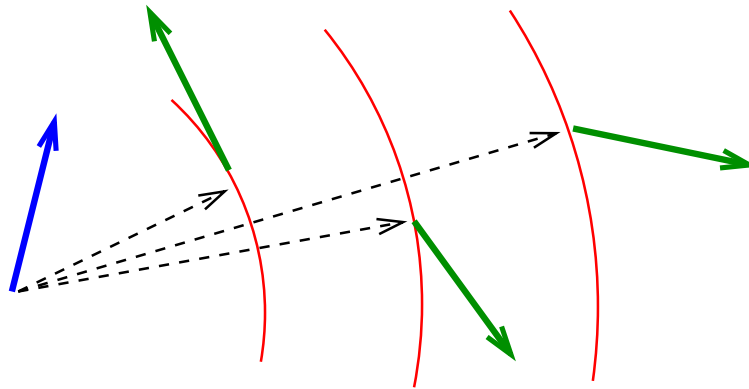
Particles interact with other particles
according to their rank
[Ballerini et al, PNAS 2008]

PNAS

Interaction ruling animal collective behavior depends on topological rather than metric distance: Evidence from a field study

M. Ballerini^{*†}, N. Cabibbo^{‡§}, R. Candelier^{†¶}, A. Cavagna^{*||**}, E. Cisbani[†], I. Giardina^{*||}, V. Lecomte^{††‡‡}, A. Orlandi^{*}, G. Parisi^{*‡§**}, A. Procaccini^{*‡}, and M. Viale^{‡§§}, and V. Zdravkovic^{*}

^{*}Centre for Statistical Mechanics and Complexity (SMC), Consiglio Nazionale delle Ricerche-Istituto Nazionale per la Fisica della Materia, [†]Dipartimento di Fisica, and [‡]Sezione Istituto Nazionale di Fisica Nucleare, Università di Roma "La Sapienza," Piazzale Aldo Moro 2, 00185 Roma, Italy; [§]Istituto Superiore di Sanità, viale Regina Elena 299, 00161 Roma, Italy; ^{||}Istituto dei Sistemi Complessi (ISC), Consiglio Nazionale delle Ricerche, via dei Taurini 19, 00185 Roma, Italy; and ^{††}Laboratoire Matière et Systèmes Complexes, (Centre National de la Recherche Scientifique Unite Mixte de Recherche 7057), Université Paris VII, 10 rue Alice Domon et Léonie Duquet, 75205 Paris Cedex 13, France



Kinetic models

Largely investigated for mean-field metric or binary interactions

Almost inexistent for topological interactions

1st work is by J. Haskovec [Physica D 2013]
for Cucker-Smale type interactions

Also work by Y. Brenier
for a competition model

Our work: Boltzmann approach

Adapting 'Choose the Leader' model from
[Carlen, D., Wennberg, M3AS 2013],
[Carlen, Chatelin, D., Wennberg, Physica D 2013]

2. Smooth rank-based dynamics

A. Blanchet & P. D., J. Stat. Phys **163** (2016) 41-60

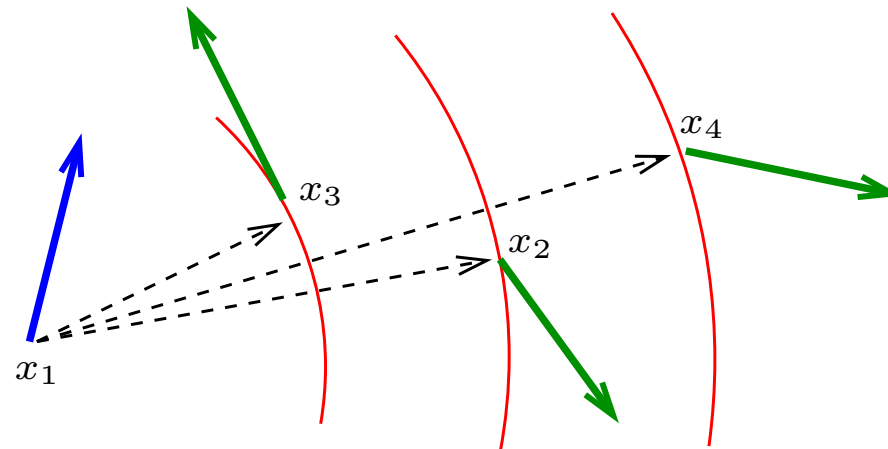
N particles $(x_i(t), v_i(t))_{i=1,\dots,N}$, $(x_i, v_i) \in \mathbb{R}^{2n}$

At given time t , Rank $R^N(i, j)$ of j w.r.t. i :

Sort $(|x_k - x_i|)_{k=1,\dots,n, k \neq i}$ by increasing order

$R^N(i, j)$ is rank of $|x_j - x_i|$ in the list

$R^N(i, j) \in \{1, \dots, N - 1\}$



$$R(1, 3) = 1 \quad R(1, 2) = 2 \quad R(1, 4) = 3$$

Normalized rank: $r^N(i, j) = \frac{R^N(i, j)}{N-1} \in \bigcup_{k=1}^N \left\{ \frac{k}{N-1} \right\}$

Given: $K : [0, 1] \rightarrow [0, \infty)$ s.t. $\int_0^1 K(r) dr = 1$

$$\text{define } K^N(r) = \frac{K(r)}{\sum_{k=1}^{N-1} K\left(\frac{k}{N-1}\right)}$$

Interaction probabilities:

π_{ij}^N probability of i interacting with j : $\pi_{ij}^N = K^N(r^N(i, j))$

$$\text{Note: } \sum_{\substack{j=1 \\ j \neq i}}^{N-1} \pi_{ij}^N = K^N(r^N(i, j)) = 1$$

Free flights $\dot{x}_i = v_i$, $\dot{v}_i = 0$

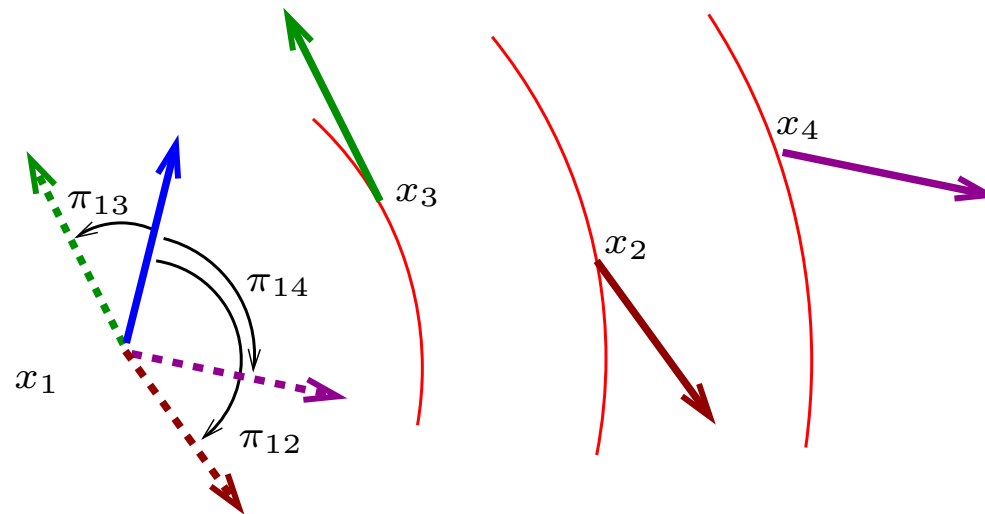
Collisions: at Poisson random times with rate N :

Pick $i \in \{1, \dots, N\}$ with uniform probability $1/N$

Pick $j \in \{1, \dots, N\}$, $j \neq i$ with probability π_{ij}^N

Perform: (x_i, x_j) remains unchanged

(v_i, v_j) changed into (v_j, v_i)



Properties

Rank is a function of positions $r^N(i, j) = r^N(i, j)(x_1, \dots, x_N)$
and so is π_{ij}^N : $\pi_{ij}^N = K^N(r^N(i, j)(x_1, \dots, x_N))$

Rank is permutation-invariant

$$r^N(\sigma(i) \sigma(j))(x_{\sigma(1)}, \dots, x_{\sigma(N)}) = r^N(i, j)(x_1, \dots, x_N)$$

Goal

Derive a Boltzmann operator
under Propagation of Chaos assumption

Note: BBGKY does not help
because all particles interact with all particles

Previous work using similar ideas:

[D. & Ringhofer, SIAM Appl. Math. 2007] (min interaction)

Notations: $\vec{x} = (x_1, \dots, x_N)$, $\vec{v} = (v_1, \dots, v_N)$
 $Z_j = (x_j, v_j)$, $\vec{Z} = (Z_1, \dots, Z_N)$, $d\vec{Z} = dZ_1 \dots dZ_N$

$f^N(\vec{Z}, t)$: N -particle distribution function

Find eq. for f^N (master equation)

Follow strategy of [Carlen, D., Wennberg, M3AS 2013]

Take $\Phi^N(\vec{Z})$ a test function

Drop drift term for simplicity

$$\frac{d}{dt} \int f^N \Phi^N d\vec{Z} =$$

$$N \int \left[\frac{1}{N} \sum_{i \neq j} \pi_{ij}^N(\vec{x}) \Phi^N(Z_1, \dots, x_i, v_j, \dots, x_j, v_j, \dots, Z_N) - \Phi^N(\vec{Z}) \right] f^N(\vec{Z}, t) d\vec{Z}$$

If $f^N|_{t=0}$ is permutation invariant

then $f^N(t)$ is permutation invariant

Marginal: $f_N^k(Z_1, \dots, Z_k, t) = \int f^N(\vec{Z}, t) dZ_{k+1} \dots dZ_N$

To compute eq. for f_N^1 , use $\Phi^N(\vec{Z}) = \phi(Z_1)$

Propagation of chaos

Assume $f^N(\vec{Z}, t) = \prod_{\ell=1}^N f_N^1(Z_\ell, t) + \text{negligible terms as } N \rightarrow \infty$

Define $\rho_N^1(x, t) = \int f_N^1(x, v, t) dv$

$$\frac{d}{dt} \int f_1^N(Z_1) \phi(Z_1) dZ_1 = \frac{1}{S^N(K)} \int (\phi(x_1, v_2) - \phi(x_1, v_1)) f_1^N(Z_1) f_1^N(Z_2) \left(\int K(r^N(1, 2)(\vec{x})) \prod_{\ell=3}^N \rho_N^1(x_\ell) dx_\ell \right) dZ_1 dZ_2$$

where

$$S^N(K) = \frac{1}{N-1} \sum_{k=1}^{N-1} K\left(\frac{k}{N-1}\right) \approx \int_0^1 K(r) dr = 1$$

Need to estimate the behavior as $N \rightarrow \infty$ of:

$$\int K(r^N(1, 2)(\vec{x})) \prod_{\ell=3}^N \rho_N^1(x_\ell) dx_\ell$$

$$\int K(r^N(1, 2)(\vec{x})) \prod_{\ell=3}^N \rho_N^1(x_\ell) dx_\ell = \mathbb{E} \left\{ K(r^N(1, 2)(x_1, x_2, x_3, \dots, x_N)) \right\}$$

when x_3, \dots, x_N are iid random variables following the law ρ_N^1

$$= \sum_{j=0}^{N-2} K\left(\frac{j+1}{N-1}\right) \binom{N-2}{j} \text{Prob}\left(\text{Card}\{k \in \{3, \dots, N\} \text{ s.t. } |x_k - x_1| \leq |x_2 - x_1|\} = j\right)$$

$$= \sum_{j=0}^{N-2} K\left(\frac{j+1}{N-1}\right) \binom{N-2}{j} m^j (1-m)^{(N-2)-j}$$

where $m = M_{\rho_N^1}(x_1, |x_2 - x_1|) = \int_{|x-x_1| \leq |x_2-x_1|} \rho_N^1(x) dx$

\approx Bernstein polynomial approximation of $K(m)$

$$= K(m) + \mathcal{O}\left(\frac{1}{N}\right)$$

In the limit $N \rightarrow \infty$, if $f_N^1(Z_1) \rightarrow f(Z_1)$, then

$$\frac{d}{dt} \int f(Z_1, t) \phi(Z_1) dZ_1 = \int (\phi(x_1, v_2) - \phi(x_1, v_1)) f(Z_1) f(Z_2) K(M_\rho(x_1, |x_2 - x_1|)) dZ_1 dZ_2$$

In strong form (adding back the drift term):

$$\partial_t f + v \cdot \nabla_x f = Q(f)$$

$$Q(f)(x, v) = \rho(x) \int f(x', v) K(M_\rho(x, |x' - x|)) dx' - f(x, v)$$

$$\text{with } M_\rho(x, s) = \int_{|x' - x| \leq s} \rho(x') dx'$$

$$\begin{aligned}\int Q(f) dv &= \rho(x) \left(\int \rho(x') K(M_\rho(x, |x' - x|)) dx' - 1 \right) \\ &= \rho(x) \left(\int K(m) dm - 1 \right) = 0\end{aligned}$$

So, the continuity eq. is satisfied

$$\partial_t \rho + \nabla_x \cdot j = 0, \quad j(x, t) = \int_{\mathbb{R}^n} f(x, v, t) v dv$$

3. Nearest neighbor

A. Blanchet & P. D., arXiv:1703.05131

At each collision event:

A particle is selected randomly with uniform probability

When selected, particle i follows its nearest neighbor

i.e. the probability π_{ij}^N of i interacting with j is

$$\pi_{ij}^N = \delta\left(R^N(i, j) - 1\right)$$

here: rank is un-normalized: $R^N(i, j) \in \{1, \dots, N - 1\}$, so:

$\pi_{ij}^N = 0$ except if $j =$ nearest neighbor in which case $\pi_{ij}^N = 1$

Rate of collision events: $\lambda(N) N$ with $\lambda(N) \rightarrow \infty$ TBD

Master eq. unchanged: - multiply by $\lambda(N)$

- replace $K^N(r^N(i, j))$ by $\delta(R^N(i, j) - 1)$

$$\frac{d}{dt} \int f_1^N(Z_1) \phi(Z_1) dZ_1 = \int (\phi(x_1, v_2) - \phi(x_1, v_1)) f_1^N(Z_1) f_1^N(Z_2) \left(\lambda(N) (N-1) \int \delta(R^N(1, 2)(\vec{x}) - 1) \prod_{\ell=3}^N \rho_N^1(x_\ell) dx_\ell \right) dZ_1 dZ_2$$

Need to estimate the behavior as $N \rightarrow \infty$ of:

$$\lambda(N) (N-1) \int \delta(R^N(1, 2)(\vec{x}) - 1) \prod_{\ell=3}^N \rho_N^1(x_\ell) dx_\ell$$

Drop scripts N and 1 for clarity i.e. $f_1^N \rightarrow f$, $\rho_1^N \rightarrow \rho$

$$\int \delta(R^N(1, 2)(\vec{x}) - 1) \prod_{\ell=3}^N \rho(x_\ell) dx_\ell = \mathbb{E} \left\{ \delta(R^N(1, 2)(\vec{x}) - 1) \right\}$$

when x_3, \dots, x_N are iid random variables following the law ρ

$$= \text{Prob} \left(\text{Card} \{ k \in \{3, \dots, N\} \text{ s.t. } |x_k - x_1| \leq |x_2 - x_1| \} = 0 \right)$$

$$= (1 - m)^{(N-2)}$$

where $m = M_\rho(x_1, |x_2 - x_1|) = \int_{|x-x_1| \leq |x_2-x_1|} \rho(x) dx$

So: $\lambda(N) (N - 1) \int \delta(R^N(1, 2)(\vec{x}) - 1) \prod_{\ell=3}^N \rho(x_\ell) dx_\ell$

$$= \lambda(N) (N - 1) (1 - M_\rho(x_1, |x_2 - x_1|))^{(N-2)}$$

$$\frac{d}{dt} \int f(Z_1) \phi(Z_1) dZ_1 = \int (\phi(x_1, v_2) - \phi(x_1, v_1)) f(x_1, v_1) f(x_2, v_2) \lambda(N) (N-1) (1 - M_\rho(x_1, |x_2 - x_1|))^{(N-2)} dx_1 dx_2 dv_1 dv_2$$

Use polar coordinates $x_2 = x_1 + r\omega$, $r \in [0, \infty)$, $\omega \in \mathbb{S}^{n-1}$

$$\frac{d}{dt} \int f(Z_1) \phi(Z_1) dZ_1 = \int (\phi(x_1, v_2) - \phi(x_1, v_1)) f(x_1, v_1) F_f^N(x_1, v_2) dx_1 dv_1 dv_2$$

$$F_f^N(x_1, v_2) = \lambda(N) (N-1) \int f(x_1 + r\omega, v_2) (1 - M_\rho(x_1, r))^{(N-2)} r^{n-1} dr d\omega$$

Change of variables $m = M_\rho(x_1, r) \Leftrightarrow r = R_\rho(x_1, m)$, $m \in [0, 1)$

$$F_f^N(x_1, v_2) = \lambda(N) (N - 1) \int_0^1 G_f(x_1, v_2, m) (1 - m)^{N-2} dm$$

$$G_f(x_1, v_2, m) = \frac{\int f(x_1 + R_\rho(x_1, m)\omega, v_2) d\omega}{\int \rho(x_1 + R_\rho(x_1, m)\omega) d\omega}$$

As $m \rightarrow 0$

$$G_f(x_1, v_2, m) = \frac{f(x_1, v_2)}{\rho(x_1)} + \frac{1}{2n} \left(\frac{m}{\rho(x_1)} \right)^{\frac{2}{n-1}} D(\rho, f)(x_1, v_2) + o\left(m^{\frac{2}{n-1}}\right)$$

$$D(\rho, f)(x_1, v_2) = \frac{1}{\rho(x_1)} \left(\Delta_x f(x_1, v_2) - \frac{f(x_1, v_2)}{\rho(x_1)} \Delta_x \rho(x_1) \right)$$

Denote:

$$C_{N,n} = (N-1) \frac{n^{\frac{2}{n}-1}}{2} \int_0^1 m^{\frac{2}{n}} (1-m)^{N-2} dm$$

$C_{N,n} \rightarrow 0$ as $N \rightarrow \infty$. For instance: $C_{N,2} = \frac{1}{2N}$

Choose: $\lambda(N) = 1/C_{N,n}$

larger frequency: $\lambda(N) \rightarrow \infty$ as $N \rightarrow \infty$

Then, as $N \rightarrow \infty$:

$$F_f^N(x_1, v_2) - \lambda(N) \frac{f(x_1, v_2)}{\rho(x_1)} \longrightarrow \left(\frac{1}{\rho(x_1)} \right)^{\frac{2}{n-1}} D(\rho, f)(x_1, v_2)$$

Probability $p_{N,n}(dm) \sim m^{\frac{2}{n-1}} (1-m)^{N-2} dm \rightarrow \delta_0$ as $N \rightarrow \infty$

In the limit $N \rightarrow \infty$, if $f_N^1(Z_1) \rightarrow f(Z_1)$, then

$$\frac{d}{dt'} \int f(x_1, v_1, t') \phi(x_1, v_1) dx_1 dv_1 = \int (\phi(x_1, v_2) - \phi(x_1, v_1)) f(x_1, v_1) \Delta_x f(x_1, v_2) \frac{1}{\rho(x_1)^{\frac{n+1}{n-1}}} dx_1 dx_2 dv_2$$

In strong form (adding back the drift term):

$$\partial_t f + v \cdot \nabla_x f = Q(f)$$

$$Q(f)(x, v) = \frac{1}{\rho(x)^{\frac{2}{n-1}}} \left(\Delta_x f(x, v) - \frac{f(x, v)}{\rho(x)} \Delta_x \rho(x) \right)$$

$$\int Q(f) dv = 0 \quad (\text{obvious})$$

So, the continuity eq. is satisfied

$$\partial_t \rho + \nabla_x \cdot j = 0, \quad j(x, t) = \int_{\mathbb{R}^n} f(x, v, t) v dv$$

4. Conclusion

Smooth rank-based topological interaction

Derivation of spatially non-local Boltzmann model

Highly nonlinear new collision operator

Nearest-neighbor interaction

gives rise to spatial nonlinear diffusion

Still preserves continuity eq.

Perspectives

Well-posedness of the resulting eqs.

Proof of propagation of chaos and convergence

Hydrodynamic limits

Extension to more complex interaction rules