

On the decay rate of the singular values of bivariate functions

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Applications: PDEs with random coefficient

Example

Given $0 < \kappa(\cdot, y) \in L^\infty(D)$, a.e., $y \in \Omega$,

$$\begin{cases} -\nabla \cdot (\kappa(\cdot, y) \nabla u(\cdot, y)) = f & \text{in } D \\ u(\cdot, y) = g & \text{on } \partial D \end{cases}$$

- ▶ Popularly used numerical methods: Stochastic Galerkin \ Collocation method
- ▶ Requirement: finite noise approximation to $\kappa(x, y)$ by

$$\kappa_M(x; Y) \approx \kappa(x, y) \text{ with } Y \text{ being a parameter of dimension } M.$$

- ▶ The number of terms M shows the complexity of the corresponding parameterized PDE.

Applications: Multi-dimensional magnetic particle imaging

Example

The functions $\{\kappa_\ell\}_{\ell=1}^L$ can be expressed using the receive coil sensitivities $\{\rho_\ell\}_{\ell=1}^L$ and the particles' mean magnetic moment vector $\bar{m} : \Omega \times I \rightarrow \mathbb{R}^3$ as $\kappa_\ell = \mu_0 \rho_\ell^t \bar{m}$. This relation follows from Faraday's law and the law of reciprocity. Then the inverse problem is to find the concentration $c : \Omega \rightarrow \mathbb{R}^+ \cup \{0\}$ from $\{v_\ell\}_{\ell=1}^L$:

$$v_\ell(t) = \int_{\Omega} c(x) \underbrace{\int_I a_\ell(t-t') \kappa_\ell(x, t') dt'}_{:=\kappa(x, t)} dx, \quad \text{with } \kappa_\ell = \mu_0 \rho_\ell^t \bar{m}.$$

- Singular value decay rate of $\kappa(x, t)$ determines the degree of ill-posedness of the inverse problem above.

Karhunen-Loevè expansion

Let $\kappa(x, \omega) \in L^2(D \times \Omega)$, we aim at utilizing as fewer terms as possible to represent $\kappa(x, \omega)$ as a finite sum of terms.

- ▶ Covariance kernel:

$$R(x, x') := \int_{\Omega} \kappa(x, \omega) \kappa(x', \omega) d\omega$$

- ▶ Covariance operator:

$$\mathcal{R} : L^2(D) \rightarrow L^2(D), \quad \mathcal{R}v = \int_D R(x, x') v(x') dx'.$$

- ▶ Eigenpairs of \mathcal{R} : seeking $\{\lambda_n, \phi_n\} \in \mathbb{R} \times L^2(D)$, s.t.,

$$\mathcal{R}\phi_n = \lambda_n \phi_n.$$

Reorder the eigenvalues by $\lambda_1 \geq \lambda_2 \geq \dots \lambda_n \geq$

Note that the operator \mathcal{R} is a Hilbert-Schmidt operator, therefore, $\lambda_n \rightarrow 0$ as $n \rightarrow \infty$.

Karhunen-Loevè expansion

- ▶ For all $\lambda_n \neq 0$, define

$$\psi_n(\omega) := \lambda_n^{-1/2} \int_D \kappa(x, \omega) \phi_n(x) dx.$$

- ▶ Karhunen-Loevè expansion:

$$\kappa(x, \omega) = \sum_{n=1}^{\infty} \sqrt{\lambda_n} \phi_n(x) \psi_n(\omega)$$

- ▶ M -term truncation:

$$\kappa_M(x, \omega) = \sum_{n=1}^M \sqrt{\lambda_n} \phi_n(x) \psi_n(\omega).$$

- ▶ Truncation estimate: $\kappa_M(x, \omega)$ is optimal in mean square error merely relying on the decay rate of $\{\lambda_n\}_{n=1}^{\infty}$

$$\|\kappa(x, \omega) - \kappa_M(x, \omega)\|_{L^2(\Omega \times D)} = \left(\sum_{n>M} \lambda_n \right)^{1/2}.$$

State of art

Table: State of the art of eigenvalue decay results and their corresponding truncation error estimates.

reference	$\kappa(x, \omega)$	λ_n	M -term truncation error	
			Condition	Rate
a	$L^2(\Omega, H^s(D))$	$\mathcal{O}(n^{-\frac{s}{d}})$	$s > d$	$\mathcal{O}(M^{\frac{1}{2} - \frac{s}{2d}})$
b	$L^2(\Omega, H^s(D))$	$\mathcal{O}(n^{-\frac{2s}{d} - \frac{1}{2}})$	$s > \frac{d}{4}$	$\mathcal{O}(M^{\frac{1}{4} - \frac{s}{d}})$
c	$H^s(\Omega \times D)$	$\mathcal{O}(n^{-\frac{2s}{d^*}})$	$s > \frac{d^*}{2}$	$\mathcal{O}(M^{\frac{1}{2} - \frac{s}{d^*}})$
d	$L^2(\Omega, H^s(I))$	-	$s > 0$	$\mathcal{O}(M^{-s})$
results of this article	$L^2(\Omega, \dot{H}^s(D))$	$\mathcal{O}(n^{-1 - \frac{2s}{d}})$	$s > 0$	$\mathcal{O}(M^{-\frac{s}{d}})$
	$L^2(\Omega, H^s(D))$			

- a Agmon, Princeton, N.J.-Toronto-London, 1965.
- b Pietsch, Cambridge University Press, Cambridge, 1987.
- c Griebel and Harbrecht, IMA J. Numer. Anal., 34(1):28-54, 2014.
- d Azaïez and Belgacem. Comput. Methods Appl. Mech. Engrg., 290:57-72, 2015.

Motivation

If $\kappa(x, \omega) \in L^2(\Omega, L^2(D))$, then $\{\lambda_n\}_{n=1}^\infty \in \ell_1$.

In view that

$$\kappa(x, \omega) = \sum_{n=1}^{\infty} \sqrt{\lambda_n} \phi_n(x) \psi_n(\omega),$$

we can obtain

$$\int_{\Omega} \left(\int_D |\kappa(x, \omega)|^2 dx \right) d\omega = \sum_{n=1}^{\infty} \lambda_n < \infty$$

Therefore, the M -term KL truncation κ_M converges to κ in $L^2(\Omega, L^2(D))$.

However, none of the previous results on the decay rate of the eigenvalues guarantee this simple fact.

Main result

Theorem (Eigenvalue estimate for the general space $L^2(\Omega, H^s(D))$)

Assume that D satisfies the strong local Lipschitz condition. Let $\kappa(x, \omega) \in L^2(\Omega, H^s(D))$. Then $\{\lambda_n\}_{n=1}^\infty \in \ell_{\frac{d}{d+2s}, 1}$. In specific, there holds

$$\lambda_n \leq \text{diam}(D)^{2s} C_{\text{em}}(d, s) C_{\text{ext}}(D, s) \|\kappa\|_{L^2(\Omega, H^s(D))}^2 n^{-1-\frac{2s}{d}},$$

where $C_{\text{ext}}(D, s)$ is a constant depending only on D and s .

Proof.

Step 1. Prove this result when $\kappa(x, \omega) \in L^2(\Omega, \dot{H}^s(D))$, where

$$\dot{H}^s(D) = \left\{ v \in L^2(D) : \sum_{j=1}^{\infty} \nu_j^s \cdot (v, \xi_j)^2 < \infty \right\}.$$

Here, $\{\nu_j, \xi_j\}_{j=1}^\infty$ denotes the eigenpairs of $A := -\Delta|_{H_0^1(D)}$ with nondecreasing eigenvalues.

Step 2. Prove this result in $L^2(\Omega, H^s(I^d))$, where $I := [0, 1]$.

Step 3. Then utilizing the extension Theorem to prove the result.

Proof to Step 1(1/3).

The Hilbert space $\dot{H}^s(D)$ is endowed with the inner product

$$(v, w)_s = \sum_{j=1}^{\infty} \nu_j^s(v, \xi_j)(w, \xi_j), \text{ for } v, w \in \dot{H}^s(D).$$

For any $s > 0$, one can now define the fractional power operator $T = A^{s/2}$ on $\dot{H}^s(D)$ by

$$Tv = \sum_{j=1}^{\infty} \nu_j^{\frac{s}{2}} \cdot (v, \xi_j) \cdot \xi_j.$$

Consequently, there holds

$$\|Tv\|_{L^2(D)}^2 = \sum_{j=1}^{\infty} \nu_j^s (v, \xi_j)^2 = |v|_s^2 \quad \text{for all } v \in \dot{H}^s(D).$$

Let $\{\mu_j\}_{j=1}^{\infty}$ be the eigenvalues of T in nondecreasing order. Then we can obtain

$$\mu_j \geq C_{\text{weyl}}(d)^{\frac{s}{2}} \text{diam}(D)^{-s} j^{s/d}.$$

Proof to Step 1(2/3).

Let

$$R_T(x, x') := \int_{\Omega} T\kappa(x', \omega) T\kappa(x, \omega) d\omega \in L^2(D \times D),$$

and denote by $\mathcal{R}_T : L^2(D) \rightarrow L^2(D)$ the Hilbert-Schmidt operator associated with the kernel function $R_T(x, x')$. Let

$$\mathcal{R}_1 = T\mathcal{R}_T,$$

with its domain $\mathcal{D}(\mathcal{R}_1) = \dot{H}^s(D)$.

The following statements are valid:

- (i) $\mathcal{R}_1 \in \mathcal{B}(\dot{H}^s(D), L^2(D))$.
- (ii) $\mathcal{R}_T|_{\dot{H}^s(D)} = \mathcal{R}_1$. Hence, \mathcal{R}_1 is a symmetric operator with \mathcal{R}_T as its self-adjoint extension operator.
- (iii) There holds the identity

$$\|T\kappa\|_{L^2(\Omega \times D)}^2 = \sum_{n=1}^{\infty} \lambda_n \|T\phi_n\|_{L^2(D)}^2.$$

Proof to Step 1(3/3).

For any $L^2(D)$ -orthonormal system $\{e_n\}_{n=1}^m \subset \dot{H}^s(D)$ of m elements, there holds

$$\sum_{n=1}^m \|Te_n\|_{L^2(D)}^2 \geq \sum_{n=1}^m \mu_n^2.$$

Together with the previous result, we arrive at

$$\begin{aligned} \sum_{n=1}^{\infty} \lambda_n \|T\phi_n\|_{L^2(D)}^2 &= (\lambda_1 - \lambda_2) \|T\phi_1\|_{L^2(D)}^2 + (\lambda_2 - \lambda_3) \left(\|T\phi_1\|_{L^2(D)}^2 + \|T\phi_2\|_{L^2(D)}^2 \right) \\ &\quad + (\lambda_3 - \lambda_4) \left(\|T\phi_1\|_{L^2(D)}^2 + \|T\phi_2\|_{L^2(D)}^2 + \|T\phi_3\|_{L^2(D)}^2 \right) + \cdots \\ &\geq (\lambda_1 - \lambda_2)\mu_1^2 + (\lambda_2 - \lambda_3)(\mu_1^2 + \mu_2^2) + (\lambda_3 - \lambda_4)(\mu_1^2 + \mu_2^2 + \mu_3^2) + \cdots \\ &= \sum_{n=1}^{\infty} \lambda_n \mu_n^2 \geq \sum_{n=1}^{\infty} \lambda_n C_{\text{weyl}}(d)^s \text{diam}(D)^{-2s} n^{\frac{2s}{d}}. \end{aligned}$$

$$\{\lambda_n\}_{n=1}^{\infty} \in \ell_{\frac{d}{d+2s}, 1} \quad \text{and} \quad \|\{\lambda_n\}_{n=1}^{\infty}\|_{\ell_{\frac{d}{d+2s}, 1}} \leq C_{\text{weyl}}(d)^{-s} \text{diam}(D)^{2s} \|\kappa\|_{L^2(\Omega, \dot{H}^s(D))}^2.$$

Numerical simulations¹

Examples

1. $\kappa(x, y) = \exp(-|x - y|) \in H^{3/2-\epsilon}([0, 1]^2)$ for $(x, y) \in [0, 1]^2$;
2. $\kappa(x, y) = (1 + |x - y|) \exp(-|x - y|) \in H^{7/2-\epsilon}([0, 1]^2)$ for $(x, y) \in [0, 1]^2$.

Then by our result, $\sqrt{\lambda_n} \approx n^{-2}$ and n^{-4} , respectively.

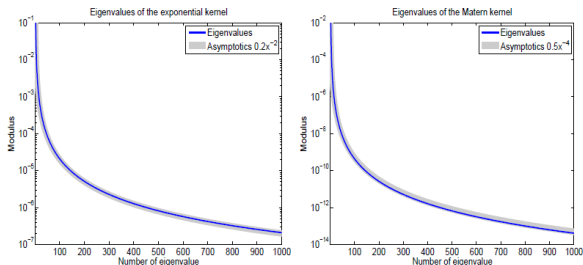


FIG. 5. The eigenvalues of the exponential kernel (left) and the Matérn kernel (right).

¹Griebel and Harbrecht, *IMA J. Numer. Analysis*, 2012.

Related publications

- ▶ Griebel and **GL**, *SIAM J. Numer. Anal.*, 56(2), 974-993, 2018.

Acknowledgement to



Thanks for your attention!