Path-by-path regularization by noise for scalar conservation laws

Benjamin Gess Max Planck Institute for Mathematics in the Sciences, Leipzig & Universität Bielefeld

CIRM conference, Stochastic Partial Differential Equations May 2018

joint work with: Panagiotis E. Souganidis, Benoit Perthame, Paul Gassiat, Mario Maurelli, Khalil Chouk

[G., Souganidis; CMS, 2014], [G., Souganidis; CPAM, 2017],

[G., Perthame, Souganidis, SINUM, 2016], [Gassiat, G.; PTRF, 2018+], [Maurelli, G.; arXiv:1701.05393], [Chouk, G.; arXiv:1512.06056].

Outline

- Introduction
- Regularization by nonlinear noise
- Path-by-path regularization by noise
- A path-by-path scaling condition

• Aim: Improved regularity and well-posedness for PDE by noise, in particular in fluid dynamics, e.g. 3*d*-Navier-Stokes equations.



- Aim: Improved regularity and well-posedness for PDE by noise, in particular in fluid dynamics, e.g. 3*d*-Navier-Stokes equations.
- Which form of noise to consider?



- Aim: Improved regularity and well-posedness for PDE by noise, in particular in fluid dynamics, e.g. 3*d*-Navier-Stokes equations.
- Which form of noise to consider?
- Additive noise, e.g. [Flandoli, Romito; TAMS, 2002]

$$du + (u \cdot \nabla)u dt + \nabla p dt = \Delta u dt + dW_t$$

div $u = 0$.

For all t > 0, \mathbb{P} -a.s. the set of singular points of u(t) is empty.



- Aim: Improved regularity and well-posedness for PDE by noise, in particular in fluid dynamics, e.g. 3*d*-Navier-Stokes equations.
- Which form of noise to consider?
- Additive noise, e.g. [Flandoli, Romito; TAMS, 2002]

$$du + (u \cdot \nabla)u dt + \nabla p dt = \Delta u dt + dW_t$$

 $div u = 0.$

For all t > 0, \mathbb{P} -a.s. the set of singular points of u(t) is empty.

• Uniqueness for the stochastic 3d-Navier-Stokes equations remains open.

- Aim: Improved regularity and well-posedness for PDE by noise, in particular in fluid dynamics, e.g. 3d-Navier-Stokes equations.
- Which form of noise to consider?
- Additive noise, e.g. [Flandoli, Romito; TAMS, 2002]

$$du + (u \cdot \nabla)u dt + \nabla p dt = \Delta u dt + dW_t$$

 $div u = 0.$

For all t > 0, \mathbb{P} -a.s. the set of singular points of u(t) is empty.

- Uniqueness for the stochastic 3*d*-Navier-Stokes equations remains open.
- More recent: Linear multiplicative noise

$$du + b(x) \cdot \nabla u \, dt = \nabla u \circ d\beta_t$$
 [Flandoli, Gubinelli, Priola; *Invent. Math.*, 2010].



• Left open: What about the nonlinear case, e.g. Burgers?

- Left open: What about the nonlinear case, e.g. Burgers?
- Linear multiplicative noise does not help anymore:

$$du + \partial_x u^2 dt = \partial_x u \circ d\beta_t.$$

Then:
$$v(t,x) := u(t,x-\beta_t)$$
 is a solution to

$$\partial_t v + \partial_x v^2 = 0.$$



- Left open: What about the nonlinear case, e.g. Burgers?
- Linear multiplicative noise does not help anymore:

$$du + \partial_x u^2 dt = \partial_x u \circ d\beta_t.$$

Then: $v(t,x) := u(t,x-\beta_t)$ is a solution to

$$\partial_t v + \partial_x v^2 = 0.$$

• In particular, weak solutions are non-unique.



- Left open: What about the nonlinear case, e.g. Burgers?
- Linear multiplicative noise does not help anymore:

$$du + \partial_x u^2 dt = \partial_x u \circ d\beta_t.$$

Then: $v(t,x) := u(t,x-eta_t)$ is a solution to

$$\partial_t v + \partial_x v^2 = 0.$$

- In particular, weak solutions are non-unique.
- Conclusion in [Flandoli, Gubinelli, Priola; Invent. Math., 2010]: "It is very easy to produce examples [...] for a stochastic version of Euler equation which show that the particular noise we use does not have any regularizing effect in this case.

- Left open: What about the nonlinear case, e.g. Burgers?
- Linear multiplicative noise does not help anymore:

$$du + \partial_x u^2 dt = \partial_x u \circ d\beta_t.$$

Then: $v(t,x) := u(t,x-eta_t)$ is a solution to

$$\partial_t v + \partial_x v^2 = 0.$$

- In particular, weak solutions are non-unique.
- Conclusion in [Flandoli, Gubinelli, Priola; Invent. Math., 2010]: "It is very easy to produce examples [...] for a stochastic version of Euler equation which show that the particular noise we use does not have any regularizing effect in this case. [...] The generalization to nonlinear transport equations, where b depends on u itself, would be a major next step for applications to fluid dynamics but it turns out to be a difficult problem."

- Left open: What about the nonlinear case, e.g. Burgers?
- Linear multiplicative noise does not help anymore:

$$du + \partial_x u^2 dt = \partial_x u \circ d\beta_t.$$

Then: $v(t,x) := u(t,x-\beta_t)$ is a solution to

$$\partial_t v + \partial_x v^2 = 0.$$

- In particular, weak solutions are non-unique.
- Conclusion in [Flandoli, Gubinelli, Priola; Invent. Math., 2010]: "It is very easy to produce examples [...] for a stochastic version of Euler equation which show that the particular noise we use does not have any regularizing effect in this case. [...] The generalization to nonlinear transport equations, where b depends on u itself, would be a major next step for applications to fluid dynamics but it turns out to be a difficult problem."
- Side-remark: In certain cases [Delarue, Flandoli, Vincenzi; CPAM, 2014] spacetime linear multiplicative noise

$$\sum_{k=1}^{\infty} e_k(x) \cdot \nabla u \circ d\beta_t^k$$

has proven useful also in nonlinear situations.



- Left open: What about the nonlinear case, e.g. Burgers?
- Linear multiplicative noise does not help anymore:

$$du + \partial_x u^2 dt = \partial_x u \circ d\beta_t.$$

Then: $v(t,x) := u(t,x-\beta_t)$ is a solution to

$$\partial_t v + \partial_x v^2 = 0.$$

- In particular, weak solutions are non-unique.
- Conclusion in [Flandoli, Gubinelli, Priola; Invent. Math., 2010]: "It is very easy to produce examples [...] for a stochastic version of Euler equation which show that the particular noise we use does not have any regularizing effect in this case. [...] The generalization to nonlinear transport equations, where b depends on u itself, would be a major next step for applications to fluid dynamics but it turns out to be a difficult problem."
- Side-remark: In certain cases [Delarue, Flandoli, Vincenzi; CPAM, 2014] spacetime linear multiplicative noise

$$\sum_{k=1}^{\infty} e_k(x) \cdot \nabla u \circ d\beta_t^k$$

has proven useful also in nonlinear situations.

Different forms of noise?

Beniamin Gess



4 / 22

Regularization by nonlinear noise

Regularization by nonlinear noise

Regularity of solutions for stochastic SCL

Consider mean field equations

$$dX_t^i = \sigma^L\left(X_t^i, \frac{1}{L}\sum_{j=1}^L \delta_{X_t^j}\right) \circ d\beta_t$$
 in \mathbb{R}^d

Taking $L o \infty$ and $\sigma^L o \sigma$ leads to stochastic scalar conservation laws

$$du + \operatorname{div}(\underbrace{\sigma(x,u)u \circ d\beta}) = 0 \quad \text{on } (0,T) \times \mathbb{R}^d.$$

Regularity of solutions for stochastic SCL

Consider mean field equations

$$dX_t^i = \sigma^L\left(X_t^i, \frac{1}{L}\sum_{j=1}^L \delta_{X_t^j}\right) \circ d\beta_t$$
 in \mathbb{R}^d

Taking $L o \infty$ and $\sigma^L o \sigma$ leads to stochastic scalar conservation laws

$$du + \operatorname{div}(\underbrace{\sigma(x,u)u}_{=:A(x,u)} \circ d\beta) = 0$$
 on $(0,T) \times \mathbb{R}^d$.

 Methods apply to general spatially homogeneous and truly nonlinear flux A, general dimension.

Regularity of solutions for stochastic SCL

Consider mean field equations

$$dX_t^i = \sigma^L \left(X_t^i, rac{1}{L} \sum_{j=1}^L \delta_{X_t^j}
ight) \circ deta_t \quad ext{in } \mathbb{R}^d$$

Taking $L o \infty$ and $\sigma^L o \sigma$ leads to stochastic scalar conservation laws

$$du + \operatorname{div}(\underbrace{\sigma(x,u)u}_{=:A(x,u)} \circ d\beta) = 0 \quad \text{on } (0,T) \times \mathbb{R}^d.$$

- Methods apply to general spatially homogeneous and truly nonlinear flux A, general dimension.
- For simplicity, in this talk restrict to

$$du + \frac{1}{2}\partial_x u^2 \circ d\beta_t = 0.$$



Consider

$$\partial_t u + rac{1}{2} \partial_{\times} u^2 = 0, \quad \text{on } (0, T) \times \mathbb{R}^d$$
 $u(0) = u_0 \in L^{\infty}(\mathbb{R}^d).$

Consider

$$\partial_t u + \frac{1}{2} \partial_x u^2 = 0$$
, on $(0, T) \times \mathbb{R}^d$
 $u(0) = u_0 \in L^{\infty}(\mathbb{R}^d)$.

For

$$\chi(t,x,v) = \chi(u(t,x),v) = 1_{v < u(t,x)} - 1_{v < 0}$$

we get the kinetic form

$$\partial_t \chi + v \partial_x \chi = \partial_v m$$
 on $(0, T) \times \mathbb{R}^d \times \mathbb{R}$.

Consider

$$\partial_t u + rac{1}{2} \partial_{\times} u^2 = 0, \quad \text{on } (0, T) \times \mathbb{R}^d$$
 $u(0) = u_0 \in L^{\infty}(\mathbb{R}^d).$

For

$$\chi(t,x,v) = \chi(u(t,x),v) = 1_{v < u(t,x)} - 1_{v < 0}$$

we get the kinetic form

$$\partial_t \chi + v \partial_{\times} \chi = \partial_v m$$
 on $(0, T) \times \mathbb{R}^d \times \mathbb{R}$.

Dissipation-dispersion approximations lead to

Definition (De Lellis, Otto, Westdickenberg, 2003)

A function $u \in L^{\infty}([0,T] \times \mathbb{R}^d)$ is said to be a quasi-solution if $\chi(t,x,v) = \chi(u(t,x),v)$ satisfies

$$\partial_t \chi + v \partial_x \chi = \partial_v m$$
 on $(0, T) \times \mathbb{R}^d \times \mathbb{R}$

for some finite (signed) measure m.

Theorem (De Lellis, Westdickenberg, 2003; Jabin, Perthame 2002)

Consider

$$\partial_t u + \frac{1}{2} \partial_x u^2 = 0$$
, on $(0, T) \times \mathbb{R}$.

Then

Theorem (De Lellis, Westdickenberg, 2003; Jabin, Perthame 2002)

Consider

$$\partial_t u + \frac{1}{2} \partial_x u^2 = 0$$
, on $(0, T) \times \mathbb{R}$.

Then

• Each quasi-solution satisfies, for all $\lambda \in (0, \frac{1}{3})$,

$$u \in L^1([0,T]; W^{\lambda,1}(\mathbb{R})).$$

Theorem (De Lellis, Westdickenberg, 2003; Jabin, Perthame 2002)

Consider

$$\partial_t u + \frac{1}{2} \partial_{\times} u^2 = 0$$
, on $(0, T) \times \mathbb{R}$.

Then

• Each quasi-solution satisfies, for all $\lambda \in (0, \frac{1}{3})$,

$$u \in L^1([0,T]; W^{\lambda,1}(\mathbb{R})).$$

② For each $\lambda > \frac{1}{3}$ there exists a quasi-solution u such that

$$u \notin L^1([0,T]; W^{\lambda,1}(\mathbb{R})).$$

Theorem (G., Souganidis; CPAM, 2016)

Let u be a quasi-solution to

$$du + \frac{1}{2}\partial_x u^2 \circ d\beta_t = 0$$
 on $(0, T) \times \mathbb{T}$.

Then,

$$u \in L^1([0,T];W^{\lambda,1}(\mathbb{T}))$$
 for all $\lambda \in (0,\frac{1}{2}), \mathbb{P}$ -a.s..

If u is an entropy solution, then

$$u(t) \in W^{\lambda,1}(\mathbb{T})$$
 for all $t > 0, \lambda \in (0,\frac{1}{2}), \mathbb{P}$ -a.s.. (\star)

Theorem (G., Souganidis; CPAM, 2016)

Let u be a quasi-solution to

$$du + rac{1}{2}\partial_{x}u^{2} \circ deta_{t} = 0 \quad \text{on } (0,T) imes \mathbb{T}.$$

Then,

$$u \in L^1([0,T];W^{\lambda,1}(\mathbb{T}))$$
 for all $\lambda \in (0,\frac{1}{2}), \mathbb{P}$ -a.s..

If u is an entropy solution, then

$$u(t) \in W^{\lambda,1}(\mathbb{T})$$
 for all $t > 0, \lambda \in (0,\frac{1}{2}), \mathbb{P}$ -a.s.. (*)

Two resulting questions:

• Can the zero set in (\star) be chosen uniformly in t?

Theorem (G., Souganidis; CPAM, 2016)

Let u be a quasi-solution to

$$du + rac{1}{2}\partial_{x}u^{2} \circ deta_{t} = 0 \quad \text{on } (0,T) imes \mathbb{T}.$$

Then,

$$u \in L^1([0,T];W^{\lambda,1}(\mathbb{T}))$$
 for all $\lambda \in (0,\frac{1}{2}), \mathbb{P}$ -a.s..

If u is an entropy solution, then

$$u(t) \in W^{\lambda,1}(\mathbb{T})$$
 for all $t > 0, \lambda \in (0,\frac{1}{2}), \mathbb{P}$ -a.s.. (\star)

Two resulting questions:

- Can the zero set in (\star) be chosen uniformly in t?
- Characterize the properties of Brownian paths leading to (*).



Path-by-path regularization by noise

Path-by-path regularization by noise

Framework

• Model example:

$$\partial_t u + \frac{1}{2} \partial_{\times} u^2 \circ dw_t = 0 \quad \text{on } (0, T) \times \mathbb{T},$$

with $w \in C([0,T];\mathbb{R})$.

Framework

• Model example:

$$\partial_t u + \frac{1}{2} \partial_X u^2 \circ dw_t = 0 \quad \text{on } (0, T) \times \mathbb{T},$$

with $w \in C([0,T];\mathbb{R})$.

• Again: Results are given for general truely nonlinear flux A.

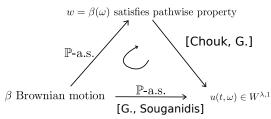
Framework

Model example:

$$\partial_t u + \frac{1}{2} \partial_x u^2 \circ dw_t = 0 \quad \text{on } (0, T) \times \mathbb{T},$$

with $w \in C([0,T];\mathbb{R})$.

- Again: Results are given for general truely nonlinear flux A.
- How to classify pathwise properties of w leading to improved regularity?



• Ideas of the proof of regularity for

$$\partial_t u + rac{1}{2} \partial_{\scriptscriptstyle X} u^2 \circ deta_t = 0 \quad ext{on } (0, \mathcal{T}) imes \mathbb{T}.$$

• Ideas of the proof of regularity for

$$\partial_t u + rac{1}{2} \partial_{\times} u^2 \circ deta_t = 0 \quad ext{on } (0,T) imes \mathbb{T}.$$

Kinetic formulation:

$$d\chi + v\partial_{x}\chi \circ d\beta_{t} = \partial_{v}m,$$

for some finite Radon measure m.

• Ideas of the proof of regularity for

$$\partial_t u + rac{1}{2} \partial_{\times} u^2 \circ d\beta_t = 0 \quad \text{on } (0,T) imes \mathbb{T}.$$

Kinetic formulation:

$$d\chi + v\partial_{x}\chi \circ d\beta_{t} = \partial_{v}m,$$

for some finite Radon measure m.

Change of variables gives

$$\chi(t,x,v) = \chi_0(x+v\beta_t,v) + \int_0^t \partial_v m(s,x+v(\beta_t-\beta_s),v)ds.$$

Averaging over velocity

$$u(t,x) = \int_{V} \chi = \int_{V} \chi_0(x + v\beta_t, v) dv + \int_0^t \int_{V} \partial_V m(s, x + v(\beta_t - \beta_s), v) dv ds.$$

Averaging over velocity

$$u(t,x) = \int_{V} \chi = \int_{V} \chi_0(x + v\beta_t, v) dv + \int_{0}^{t} \int_{V} \partial_V m(s, x + v(\beta_t - \beta_s), v) dv ds.$$

• The averaging effect appears since the velocity average in *v* contains averaging of the *x*-variable.

Averaging over velocity

$$u(t,x) = \int_{V} \chi = \int_{V} \chi_{0}(x + v\beta_{t}, v)dv + \int_{0}^{t} \int_{V} \partial_{V} m(s, x + v(\beta_{t} - \beta_{s}), v)dvds.$$

- The averaging effect appears since the velocity average in v contains averaging of the x-variable.
- Rigorously, this can be seen by Fourier transform, that is,

$$\hat{u}(t,n) = \int_{V} e^{-iv\beta_t n} \hat{\chi}_0(n,v) dv + \int_0^t \int_{V} e^{-iv(\beta_t - \beta_s)n} \partial_V \hat{m}(s,n,v) dv ds.$$

Averaging over velocity

$$u(t,x) = \int_{V} \chi = \int_{V} \chi_0(x + v\beta_t, v) dv + \int_{0}^{t} \int_{V} \partial_V m(s, x + v(\beta_t - \beta_s), v) dv ds.$$

- The averaging effect appears since the velocity average in v contains averaging of the x-variable.
- Rigorously, this can be seen by Fourier transform, that is,

$$\hat{u}(t,n) = \int_{V} e^{-iv\beta_{t}n} \hat{\chi}_{0}(n,v) dv + \int_{0}^{t} \int_{V} e^{-iv(\beta_{t}-\beta_{s})n} \partial_{V} \hat{m}(s,n,v) dv ds.$$

ullet The oscillatory integrals have a regularizing effect, both in u and in $eta_t - eta_s$.

Framework

• For SDE this has been considered by [Catellier, Gubinelli; *SPA*, 2016]: A path $w \in C(\mathbb{R}_+; \mathbb{R}^d)$ is said to be (ρ, γ) -irregular if

$$|\int_{s}^{t} e^{iw_{r} \cdot n} dr| \lesssim (1+|n|)^{-\rho} |t-s|^{\gamma} \quad \forall n \in \mathbb{R}^{d}, \, s < t.$$

Framework

• For SDE this has been considered by [Catellier, Gubinelli; *SPA*, 2016]: A path $w \in C(\mathbb{R}_+; \mathbb{R}^d)$ is said to be (ρ, γ) -irregular if

$$|\int_{s}^{t} e^{iw_{r} \cdot n} dr| \lesssim (1+|n|)^{-\rho} |t-s|^{\gamma} \quad \forall n \in \mathbb{R}^{d}, \, s < t.$$

Note:

$$\int_{s}^{t} e^{iw_{r} \cdot n} dr = \int_{\mathbb{R}} e^{ix \cdot n} dL_{w}^{s,t}(x) = L_{w}^{\hat{s},t}(n)$$

the Fourier transform of the local time.

Main result

Theorem

Let $w \in C^{\eta}([0,T],\mathbb{R}^d)$ for some $\eta>0$ be (ρ,γ) -irregular, u a quasi-solution solution to

$$\partial_t u + \frac{1}{2} \partial_x u^2 \circ dw_t = 0$$
 on $(0, T) \times \mathbb{T}$.

Then, for all

$$\lambda < \frac{\rho(\eta+1) - (1-\gamma)}{(\rho \vee 1)(\eta+1) + (1-\gamma)},$$

we have

$$u \in L^1([0,T]; W^{\lambda,1}(\mathbb{T})).$$

Corollary

Let β^H be a fractional Brownian motion with Hurst parameter $H \in (0, \frac{1}{2}]$ and u be a quasi-solution to

$$\partial_t u + \frac{1}{2} \partial_x u^2 \circ d\beta_t^H = 0$$
 on \mathbb{T} . (1)

Then, \mathbb{P} -a.s. for all $\lambda < \frac{1}{1+2H}$,

$$u \in L^1([0,T]; W^{\lambda,1}(\mathbb{T})).$$

Corollary

Let β^H be a fractional Brownian motion with Hurst parameter $H \in (0, \frac{1}{2}]$ and u be a quasi-solution to

$$\partial_t u + \frac{1}{2} \partial_x u^2 \circ d\beta_t^H = 0 \quad \text{on } \mathbb{T}.$$
 (1)

Then, \mathbb{P} -a.s. for all $\lambda < \frac{1}{1+2H}$,

$$u \in L^1([0,T]; W^{\lambda,1}(\mathbb{T})).$$

• Note: Fully recover the probabilistic result from [G., Souganidis; *CPAM*, 2016]: For $H = \frac{1}{2}$ get $\lambda < \frac{1}{2}$.

A path-by-path scaling condition

A path-by-path scaling condition

• The proof given in [G., Souganidis; *CPAM*, 2016] uses the *scaling property* of Brownian motion and independence of increments.

- The proof given in [G., Souganidis; *CPAM*, 2016] uses the *scaling property* of Brownian motion and independence of increments.
- However: (ρ, γ) -irregularity depends on two parameters, also encoding a time regularity. Hence, does not seem to be optimal.

- The proof given in [G., Souganidis; *CPAM*, 2016] uses the *scaling property* of Brownian motion and independence of increments.
- However: (ρ, γ) -irregularity depends on two parameters, also encoding a time regularity. Hence, does not seem to be optimal.
- Moreover: (ρ, γ) -irregularity not easy to check.

- The proof given in [G., Souganidis; CPAM, 2016] uses the scaling property of Brownian motion and independence of increments.
- However: (ρ, γ) -irregularity depends on two parameters, also encoding a time regularity. Hence, does not seem to be optimal.
- Moreover: (ρ, γ) -irregularity not easy to check.
- To avoid the use of oscillatory integrals: Completely avoid Fourier methods in the proof.

Consider

$$\partial_t u + \frac{1}{2} \partial_x u^2 \circ dw_t = 0.$$

Consider

$$\partial_t u + \frac{1}{2} \partial_x u^2 \circ dw_t = 0.$$

Kinetic form

$$\partial_t \chi + v \partial_x \chi \circ dw_t = \partial_v m.$$

Consider

$$\partial_t u + \frac{1}{2} \partial_x u^2 \circ dw_t = 0.$$

Kinetic form

$$\partial_t \chi + v \partial_x \chi \circ dw_t = \partial_v m.$$

• Rewrite as, for $\lambda > 0$,

$$\partial_t \chi + v \partial_x \chi \circ dw_t + \lambda \chi = \partial_v m + \lambda \chi.$$

Change of variables, take velocity integral and drop i.c.:

$$u(t,x) = \int_{V} \chi(t,x,v) dv = \int_{0}^{t} \int_{V} e^{-\lambda(t-s)} (\partial_{V} m)(s,x-vw_{s,t},v) dvds$$
$$+ \lambda \int_{0}^{t} \int_{V} e^{-\lambda(t-s)} \chi(s,x-vw_{s,t},v) dvds.$$

Change of variables, take velocity integral and drop i.c.:

$$u(t,x) = \int_{V} \chi(t,x,v) dv = \int_{0}^{t} \int_{V} e^{-\lambda(t-s)} (\partial_{V} m)(s,x-vw_{s,t},v) dvds$$
$$+ \lambda \int_{0}^{t} \int_{V} e^{-\lambda(t-s)} \chi(s,x-vw_{s,t},v) dvds.$$

Introduce the random X-ray transform

$$(Tg)(t,x) := \int_0^t \int_V g(s,x-vw_{s,t},v)e^{-\lambda(t-s)} dvds.$$

Change of variables, take velocity integral and drop i.c.:

$$u(t,x) = \int_{V} \chi(t,x,v) dv = \int_{0}^{t} \int_{V} e^{-\lambda(t-s)} (\partial_{V} m)(s,x-vw_{s,t},v) dvds$$
$$+ \lambda \int_{0}^{t} \int_{V} e^{-\lambda(t-s)} \chi(s,x-vw_{s,t},v) dvds.$$

Introduce the random X-ray transform

$$(Tg)(t,x) := \int_0^t \int_V g(s,x-vw_{s,t},v)e^{-\lambda(t-s)} dvds.$$

Hence,

$$u := \int_{\mathcal{X}} \chi = T(\partial_{\nu} m) + \lambda T \chi.$$

where m is a finite measure and $\chi(t,x,v) := 1_{[0,u(t,x)]}(v)$.



Change of variables, take velocity integral and drop i.c.:

$$u(t,x) = \int_{V} \chi(t,x,v) dv = \int_{0}^{t} \int_{V} e^{-\lambda(t-s)} (\partial_{V} m)(s,x-vw_{s,t},v) dvds$$
$$+ \lambda \int_{0}^{t} \int_{V} e^{-\lambda(t-s)} \chi(s,x-vw_{s,t},v) dvds.$$

Introduce the random X-ray transform

$$(Tg)(t,x) := \int_0^t \int_V g(s,x-vw_{s,t},v)e^{-\lambda(t-s)} dvds.$$

Hence,

$$u := \int_{\mathcal{X}} \chi = T(\partial_{\mathcal{X}} m) + \lambda T \chi.$$

where m is a finite measure and $\chi(t,x,v):=1_{[0,u(t,x)]}(v)$.

• Strategy: Estimate the regularity of $T(\partial_{\nu} m)$, $T\chi$ then use real interpolation.

Path-by-path scaling condition

• This leads to: Path-by-path scaling condition: Assume that there is a $\iota \in (0,1]$ such that for every $\sigma \in [0,1)$, $\lambda \geq 1$ we have

$$\int_0^T \int_0^{T-r} e^{-\lambda t} |\underbrace{w_{t+r} - w_r}|^{-\sigma} dt dr \lesssim \lambda^{-1+\iota\sigma}.$$

Path-by-path scaling condition

• This leads to: Path-by-path scaling condition: Assume that there is a $\iota \in (0,1]$ such that for every $\sigma \in [0,1)$, $\lambda \geq 1$ we have

$$\int_0^T \int_0^{T-r} e^{-\lambda t} |\underbrace{w_{t+r} - w_r}|^{-\sigma} dt dr \lesssim \lambda^{-1+\iota\sigma}.$$

• Easy to see: (ρ, γ) -irregularity implies path-by-path scaling.

Theorem

Let u be a quasi-solution to

$$\partial_t u + \frac{1}{2} \partial_x u^2 \circ dw_t = 0$$
 on \mathbb{R}

and suppose that $w \in C^\eta$ satisfies path-by-path scaling. Then, for all $\lambda < \frac{1+\eta-\iota}{1+\eta+\iota}$,

$$u \in L^1([0,T]; W^{\lambda,1}(\mathbb{T})).$$