30th Oct.~3rd Nov. 2017

"Collisionless Boltzmann (Vlasov) Equation and Modeling of Self-Gravitating Systems and Plasmas"

@ CIRM, Marseille



Perturbative description of Vlasov-Poisson system in cosmology : Approaching and going beyond shell-crossing

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Plan of talk

Perturbative description of 'shell-crossing' and beyond in nonlinear regime of cosmic structure formation

Refs. AT & Colombi, MNRAS 470, 4858 ('17), Saga, Colombi & AT (in prep.)

Introduction & motivation

- Beyond shell-crossing in ID
- Approaching shell-crossing in 3D

• Summary

Large-scale structure of the Universe

Matter inhomogeneity over Giga parsec scales 1000 Mpc = 3*10^9 light years

is dominated by hypothetical invisible matter (i.e., cold dark matter)

It has evolved from tiny fluctuations (most likely seeded by inflation) under influence of cosmic expansion and gravity

- Provide a wealth of cosmological information
- Key observable in the era of precision cosmology

Macroscopically,

self-gravitating collisionless system in an expanding universe

Cosmological Vlasov-Poisson system

Vlasov-Poisson system in a cosmological background:

$$\begin{bmatrix} \frac{\partial}{\partial t} + \frac{\mathbf{p}}{ma^2} \frac{\partial}{\partial \mathbf{x}} - m \frac{\partial \Phi}{\partial \mathbf{x}} \frac{\partial}{\partial \mathbf{p}} \end{bmatrix} f(\mathbf{x}, \mathbf{p}) = 0,$$

$$= \frac{2}{m} \int d\mathbf{x} d\mathbf{x$$

$$\nabla^2 \Phi(\boldsymbol{x}) = 4\pi \, G \, a^2 \begin{bmatrix} \frac{m}{a^3} \int d^3 \boldsymbol{p} \, f(\boldsymbol{x}, \boldsymbol{p}) - \rho_{\rm m} \end{bmatrix} \begin{array}{l} {\rm a(t): scale \ factor \ of \ the \ Universe} \end{array}$$

Cold initial flow (or single-stream flow): Dirac's delta function

$$f(\boldsymbol{x}, \boldsymbol{p}) = \overline{n} a^3 \{1 + \delta_{m}(\boldsymbol{x})\} \delta_{D} [\boldsymbol{p} - m a \boldsymbol{v}(\boldsymbol{x})]$$

Mass density field Velocity field

System at an early phase is reduced to pressureless fluid system

Cosmic fluid and perturbation theory

Assuming single-stream flow,

cosmological Vlasov-Poisson system is reduced to fluid system



Shell-crossing: inevitable feature

Fluid description of matter flow, however, breaks down at late time, followed by

Shell crossing

e.g., Shandarin & Zel'dovich ('89)

- Singular density field
- Development of multi-stream flow after shell-crossing



http://www.vlasix.org/index.php?n=Main.ColDICE

Analytic description of shell-crossing

How well one can analytically describe shell-crossing and beyond ?

Self-similar solution under a certain symmetry (spherical, planar, ...)



Filmore & Goldreich ('84); Bertschinger ('85); Ryden ('93); Zukin & Bertschinger ('10); Lithwick & Dalal ('11);Alard ('13)

In this talk,

Perturbative description of formation of shell-crossing and beyond

Key Lagrangian description of matter flow:

 $\boldsymbol{x}(\boldsymbol{q},t) = \boldsymbol{q} + \boldsymbol{\Psi}(\boldsymbol{q},t), \qquad \boldsymbol{v}(\boldsymbol{q},t) = \frac{d\boldsymbol{\Psi}(\boldsymbol{q},t)}{dt}$

q : Lagrangian coordinate Ψ : displacement field

Zel'dovich solution

Zel'dovich ('70), Novikov ('69)

Exact single-stream solution in ID:

Shandarin & Zel'dovich ('89)

$$x(q;\tau) = q + \psi(q) D_{+}(\tau) \qquad \mathbf{v}(q;\tau) = \psi(q) \frac{dD_{+}(\tau)}{d\tau}$$

 $D_+(au)$: linear growth factor $\psi(q)$: arbitrary function of q



Beyond shell-crossing, solution is no longer exact

→ exploiting new treatment with post-collapse PT in ID
 In 2D & 3D, it is no more exact even at single-stream regime
 → improved description by higher-order Lagrangian PT in 3D

Beyond shell-crossing in I

AT & Colombi, MNRAS 470, 4858('17) see also, Colombi, ibid. 446, 2902 ('15)

Post-collapse perturbation theory (PT) formalism

Computing back-reaction to the Zel'dovich flow:

New

I. Expand the displacement field around shell-crossing (Lagrangian) point, **q**o:

 $x(q;\tau) \simeq A(q_0;\tau) - B(q_0;\tau)(q-q_0) + C(q_0;\tau)(q-q_0)^3$



2. Compute force $F(x(q;\tau)) = -\nabla_x \Phi(x(q;\tau))$ at multi-valued region Newton potential

 $\Delta \mathbf{v}(Q;\tau,\,\tau_{\mathbf{q}}) = \int_{\tau}^{\tau} d\tau' F(x(Q,\tau')), \quad \Delta x(Q;\tau,\,\tau_{\mathbf{q}}) = \int_{\tau}^{\tau} d\tau' \,\Delta \mathbf{v}(Q;\tau',\tau_{\mathbf{q}})$ **Backreaction** ••••• polynomial function of $Q=q-q_0$ up to 7th order

Post-collapse PT: single cluster

AT & Colombi ('17)

Simulation

Post-collapse PT basically can capture the postcollapse dynamics until 2nd shell-crossing happens



Even after 2nd crossing,

it still gives a reasonable description for density profiles

Post-collapse PT: random initial condition

Initial condition

Planck

AT & Colombi ('17)



Approaching shell-crossing in 3D

Improved description of shell-crossing by Poster by Shohei Saga

higher-order Lagrangian perturbation theory (LPT)

Lagrangian $\Psi({\bm q},t) = \Psi^{(1)}({\bm q},t) + \Psi^{(2)}({\bm q},t) + \Psi^{(3)}({\bm q},t) + \cdots$ displacement field

Eq. of motion for displacement field assuming single-stream flow

Longitudinal: $(\hat{\mathcal{T}} - 4\pi G \bar{\rho}_{m}) \Psi_{k,k} = -\epsilon_{ijk} \epsilon_{ipq} \Psi_{j,p} (\hat{\mathcal{T}} - 2\pi G \bar{\rho}_{m}) \psi_{k,q}$ $-\frac{1}{2} \epsilon_{ijk} \epsilon_{pqr} \Psi_{i,p} \Psi_{j,q} (\hat{\mathcal{T}} - \frac{4\pi G}{3} \bar{\rho}_{m}) \Psi_{k,r}$ Transverse: $\epsilon_{ijk} \hat{\mathcal{T}} \Psi_{j,k} = -\epsilon_{ijk} \Psi_{p,j} \hat{\mathcal{T}} \Psi_{p,k}$ $\hat{\mathcal{T}} f(t) \equiv \ddot{f}(t) + 2H\dot{f}(t)$

Moutarde et al. ('91); Bouchet et al. ('92); Buchert ('92); Buchert & Ehlers ('93); Bouchet et al. ('95), ..., Matsubara ('15), Rampf & Frisch ('17)

Approaching shell-crossing in 3D

Improved description of shell-crossing by

higher-order Lagrangian perturbation theory (LPT)



Simulation: ColDICE (Sousbie & Colombi '16)

sine-wave initial condition (Moutarde et al. '91)

Poster by Shohei Saga

Approaching shell-crossing in 3D

Improved description of shell-crossing by

higher-order Lagrangian perturbation theory (LPT)



Simulation: ColDICE (Sousbie & Colombi '16)

(Moutarde et al. '91)

Poster by Shohei Saga

Beyond shell-crossing in 3D (in progress)

Matching higher-order LPT with post-collapse PT (in ID)

ignoring interactions along y-/z-directions

Vlasov simulation LPT+post-collapse PT



Multi-dimensional effect significantly alters multi-stream flow. Naive prediction with ID post-collapse PT fails \rightarrow need further progress

Summary

Analytical description of shell-crossing and beyond with perturbation theory (PT) calculations in cosmic structure formation

	Single-stream	Multi-stream
ID	Zel'dovich solution	Post-collapse PT
2D & 3D	Higher-order Lagrangian PT	Work in progress

- Combining with adaptive smoothing, Post-collapse PT gives an accurate statistical prediction of random density field
- Higher-order Lagrangian PT

precisely matches results of the state-of-the-art Vlasov simulation

Extension to 3D case is still challenging, but worth for investigation