

Acoustic waves in the incompressible limits of viscous fluids

Eduard Feireisl

based on joint work with D. Donatelli (L'Aquila)

Institute of Mathematics, Academy of Sciences of the Czech Republic, Prague

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A model arising in astrophysics

Primitive system

$$\partial_t \varrho + \operatorname{div}_x(\varrho \mathbf{u}) = 0$$

$$\partial_t(\varrho \mathbf{u}) + \operatorname{div}_x(\varrho \mathbf{u} \otimes \mathbf{u}) + \frac{1}{\varepsilon^2} \nabla_x(\varrho \Theta)^\gamma = \varepsilon^\alpha \operatorname{div}_x \mathbb{S}(\nabla_x \mathbf{u}) + \frac{1}{\varepsilon^2} \varrho \nabla_x F$$

$$\partial_t(\varrho \Theta) + \operatorname{div}_x(\varrho \Theta \mathbf{u}) = 0$$

$$\mathbb{S}(\nabla_x \mathbf{u}) = \left(\nabla_x \mathbf{u} + \nabla_x^t \mathbf{u} - \frac{2}{3} \operatorname{div}_x \mathbf{u} \mathbb{I} \right) + \lambda \operatorname{div}_x \mathbf{u} \mathbb{I}, \quad \lambda \geq 0,$$

Gravitational potential

$$F \in C^\infty(\mathbb{R}^3), \quad F(x) > 0, \quad |x|^2 |\nabla_x F(x)| + |x|^3 |\nabla_x^2 F(x)| \leq c \text{ for all } x \in \mathbb{R}^3$$

$$\underline{F} \frac{1}{|x|} \leq F(x) \leq \bar{F} \frac{1}{|x|} \text{ for all } |x| > R$$

Far field conditions

$$\varrho \rightarrow \bar{\varrho} > 0, \quad \mathbf{u} \rightarrow 0, \quad \Theta \rightarrow 1 \text{ as } |x| \rightarrow \infty$$

Low Mach number approximation of (certain) supernovae

Target system

$$\operatorname{div}_x(\varrho_0 \mathbf{V}) = 0$$

$$\partial_t \mathbf{V} + \mathbf{V} \cdot \nabla_x \mathbf{V} + \nabla_x \Pi = -\frac{\varrho_0}{\mathcal{R}} \nabla_x F$$

$$\partial_t(\mathcal{R}) + \operatorname{div}_x(\mathcal{R} \mathbf{V}) = 0$$

Static density profile

$$\nabla_x \varrho_0^\gamma = \varrho_0 \nabla_x F \text{ in } R^3, \varrho_0 \rightarrow \bar{\varrho} > 0 \text{ as } |x| \rightarrow \infty$$

Initial state of the primitive system

Initial data

$$\varrho_{0,\varepsilon}(0, \cdot) \approx \varrho_0 + \varepsilon \varrho_0^1$$

$$\mathbf{u}_{0,\varepsilon}(0, \cdot) \approx \mathbf{V}_0 + \nabla_x \Phi_0$$

$$\Theta_{0,\varepsilon}(0, \cdot) = 1 + \varepsilon^2 \Theta_0^{(2)}$$

Generalized Helmholtz decomposition

$$\mathbf{v} = \mathbf{H}_{\varrho_0}[\mathbf{v}] + \nabla_x \Phi$$

$$\operatorname{div}_x (\varrho_0 \nabla_x \Phi) = \operatorname{div}_x (\varrho_0 \mathbf{v}), \quad \Phi \in D^{1,2}(R^3) \text{ for } \mathbf{v} \in L^2(R^3, R^3).$$

Well vs ill prepared data

$$\text{Well: } \varrho_0^1 = 0, \quad \nabla_x \Phi = 0, \quad \Theta_0^{(2)} \in L^1 \cap L^\infty$$

Ill: otherwise

Expected conclusion

Convergence

$$\varrho_\varepsilon \approx \varrho_0 + \boxed{\varepsilon S_\varepsilon} \rightarrow \varrho_0$$

$$\mathbf{u}_\varepsilon \approx \mathbf{V} + \boxed{\nabla_x \Phi}_\varepsilon \rightarrow \mathbf{V}$$

Potential temperature

$$\Theta_\varepsilon \approx 1 + \varepsilon^2 \Theta_\varepsilon^{(2)}, \quad \Theta_\varepsilon^{(2)} \rightarrow \frac{\varrho_0}{\mathcal{R}}$$

Acoustic waves

Acoustic equation

$$\varepsilon \partial_t s_\varepsilon + \operatorname{div}_x [\varrho_0 \nabla_x \Phi_\varepsilon] = 0$$

$$\varepsilon \varrho_0 \partial_t \nabla_x \Phi_\varepsilon + \varrho_0 \nabla_x \left[\frac{p'(\varrho_0)}{\varrho_0} s_\varepsilon \right] = 0$$

or, in a concise form

$$\varepsilon^2 \partial_{t,t}^2 \Phi_\varepsilon - \frac{p'(\varrho_0)}{\varrho_0} \operatorname{div}_x (\varrho_0 \nabla_x \Phi_\varepsilon) = 0$$

Initial data

$$s_\varepsilon(0, \cdot) = \varrho_0^{(1)}, \quad \nabla_x \Phi_\varepsilon(0, \cdot) = \nabla_x \Phi_0$$

Acoustic propagator

$$\mathcal{A}_{\varrho_0} : v \mapsto -\frac{p'(\varrho_0)}{\varrho_0} \operatorname{div}_x (\varrho_0 \nabla_x v)$$

Self-adjoint operator

$$\mathcal{A}_{\varrho_0} : v \mapsto -\frac{p'(\varrho_0)}{\varrho_0} \operatorname{div}_x (\varrho_0 \nabla_x v)$$

$$v \in L^2(\mathbb{R}^3), \langle u; v \rangle = \int_{\mathbb{R}^3} uv \frac{\varrho_0}{p'(\varrho_0)} dx$$

Spectrum

- the point spectrum of \mathcal{A}_{ϱ_0} is empty
- the operator \mathcal{A}_{ϱ_0} satisfies the *Limiting absorption principle*

Operators

$$\mathcal{V} \circ (\mathcal{A}_{\varrho_0} - \lambda \pm i\eta)^{-1} \circ \mathcal{V} : L^2(\Omega) \rightarrow L^2(\Omega),$$

$$\mathcal{V}[v] = (1 + |x|^2)^{-s/2}, \quad s > 1$$

are bounded uniformly for $\lambda \in [a, b]$, $0 < a < b$, $\eta > 0$,

Local decay properties

Theorem (Kato)

Let C be a closed densely defined linear operator and H a self-adjoint densely defined linear operator in a Hilbert space X . For $\lambda \notin R$, let $R_H[\lambda] = (H - \lambda \text{Id})^{-1}$ denote the resolvent of H . Suppose that

$$\Gamma = \sup_{\lambda \notin R, v \in \mathcal{D}(C^*), \|v\|_X=1} \|C \circ R_H[\lambda] \circ C^*[v]\|_X < \infty.$$

Then

$$\sup_{w \in X, \|w\|_X=1} \frac{\pi}{2} \int_{-\infty}^{\infty} \|C \exp(-itH)[w]\|_X^2 dt \leq \Gamma^2.$$

Local decay estimates

$$\int_{-\infty}^{\infty} \left\| \varphi G(\mathcal{A}_{\ell_0}) \exp\left(\pm i\sqrt{\mathcal{A}_{\ell_0}} t\right) [h] \right\|_{L^2(\mathbb{R}^3)}^2 dt \leq c(\varphi, G) \|h\|_{L^2(\mathbb{R}^3)}^2$$

$$\varphi \in C_c^\infty(\mathbb{R}^3), G \in C_c^\infty(0, \infty)$$

Improved decay estimates

Strichartz estimates (frequency localized)

$$\int_{-\infty}^{\infty} \left\| G(\mathcal{A}_{\varrho_0}) \exp\left(\pm i\sqrt{\mathcal{A}_{\varrho_0}} t\right) [h] \right\|_{L^q(\mathbb{R}^3)}^p dt \leq c(G) \|h\|_{L^2(\mathbb{R}^3)}^p$$

for any any $G \in C_c^\infty(0, \infty)$

$$\frac{1}{p} + \frac{3}{q} = \frac{1}{2}.$$

Decay for the scaled equation

$$\begin{aligned} & \int_0^T \left\| G(\mathcal{A}_{\varrho_0}) \exp\left(\pm i\sqrt{\mathcal{A}_{\varrho_0}} \frac{t}{\varepsilon}\right) [h] \right\|_{L^q(\mathbb{R}^3)}^p dt \\ &= \varepsilon \int_0^T \frac{1}{\varepsilon} \left\| G(\mathcal{A}_{\varrho_0}) \exp\left(\pm i\sqrt{\mathcal{A}_{\varrho_0}} \frac{t}{\varepsilon}\right) [h] \right\|_{L^q(\mathbb{R}^3)}^p dt \\ &= \varepsilon \int_0^{T/\varepsilon} \left\| G(\mathcal{A}_{\varrho_0}) \exp\left(\pm i\sqrt{\mathcal{A}_{\varrho_0}} t\right) [h] \right\|_{L^q(\mathbb{R}^3)}^p dt \leq \boxed{\varepsilon} c(G) \|h\|_{L^2(\mathbb{R}^3)}^p \end{aligned}$$

Strichartz estimates by Metcalfe and Tataru

A wave equation

$$\partial_{t,t}^2 V - \operatorname{div}_x \left(\tilde{\mathbf{A}}(x) \nabla_x V \right) + \tilde{\mathbf{B}}(x) \cdot \nabla_x V = Z, \quad V(0, \cdot) = V_0, \quad \partial_t V(0, \cdot) = V_1$$

Strichartz estimates

$$\|V\|_{L^p(-\infty, \infty; L^q(\mathbb{R}^3))} \leq (\|\nabla_x V_0\|_{L^2(\mathbb{R}^3; \mathbb{R}^3)} + \|V_1\|_{L^2(\mathbb{R}^3)} + \|Z\|_{L^r(-\infty, \infty; L^s(\mathbb{R}^3))})$$
$$\frac{1}{p} + \frac{3}{q} = \frac{1}{2} = \frac{1}{r} + \frac{3}{s} - 2$$

Flatness of coefficients

$$\sum_{j \in \mathbb{Z}} \sup_{x \in A_j} \left(|x|^2 |\nabla_x^2 \tilde{\mathbf{A}}(x)| + |x| |\tilde{\mathbf{A}}(x)| + |\tilde{\mathbf{A}}(x) - \bar{\mathbf{A}}| \right) \leq \delta, \quad \bar{\mathbf{A}} > 0$$

$$\sum_{j \in \mathbb{Z}} \sup_{x \in A_j} \left(|x|^2 |\nabla_x \tilde{\mathbf{B}}(x)| + |x| |\tilde{\mathbf{B}}(x)| \right) \leq \delta$$

$$A_j = \{2^j \leq |x| \leq 2^{j+1}\}$$

Decomposition method

A new density profile

$\tilde{\varrho}, \tilde{\varrho}$ flat on a large ball $\{|x| < R\}$, $\tilde{\varrho} = \varrho_0$ for $\{|x| > 2R\}$

Modified operator

$$\begin{aligned}\mathcal{A}_{\tilde{\varrho}}[v] &= \operatorname{div}_x (p'(\tilde{\varrho})\nabla_x v) + \tilde{\varrho}Q''(\tilde{\varrho})\nabla_x \tilde{\varrho} \cdot \nabla_x v, \\ \tilde{\mathbf{A}}(x) &= p'(\tilde{\varrho})(x), \quad \tilde{\mathbf{B}}(x) = \tilde{\varrho}Q''(\tilde{\varrho})\nabla_x \tilde{\varrho}.\end{aligned}$$

Strichartz estimates

The operator \mathcal{A}_{ϱ_0} satisfies the flatness condition of Metcalfe and Tataru

A method by Smith and Sogge

Desired estimates

$$\int_{-\infty}^{\infty} \left\| G(\mathcal{A}_{\ell_0}) \exp(\pm i\sqrt{\mathcal{A}_{\ell_0}} t) [h] \right\|_{L^q(\mathbb{R}^3)}^p dt \leq c(G) \|h\|_{L^2(\mathbb{R}^3)}^p$$

Ansatz

$$U = G(\mathcal{A}_{\ell_0}) \exp(\pm i\sqrt{\mathcal{A}_{\ell_0}} t) [h]$$

$$\chi \in C_c^\infty(\mathbb{R}), \chi(z) = \begin{cases} 1 & \text{for } |z| < R \\ 0 & \text{for } |z| > 2R, \end{cases}$$

$$U = \chi(|x|)U + (1 - \chi(|x|))U,$$

Local + Strichartz

$\chi(|x|)U$ satisfies the localized frequency cut-off estimates as U does

Decay estimates for the far field component

Far field vector field

$$V = (1 - \chi(|x|))U$$

Wave equation

$$\partial_{t,t}^2 V - \operatorname{div}_x (\tilde{\mathbf{A}}(x) \nabla_x V) + \tilde{\mathbf{B}}(x) \cdot \nabla_x V = Z,$$

$$Z = \tilde{\mathbf{A}} \nabla_x \chi \cdot \nabla_x U + \operatorname{div}_x (\tilde{\mathbf{A}} \nabla_x \chi U) - \tilde{\mathbf{B}} \cdot \nabla_x \chi U$$

$$V(0, \cdot) = (1 - \chi)U_0, \quad \partial_t V(0, \cdot) = (1 - \chi)U_1.$$

Strichartz estimates

The result of Metcalfe and Tataru applies to V

Convergence proof

Relative energy

$$\begin{aligned} \mathcal{E}(\varrho, \Theta, \mathbf{u} \mid r, \mathbf{U}) \\ = \int_{R^3} \left[\frac{1}{2} \varrho |\mathbf{u} - \mathbf{U}|^2 + \frac{H(\varrho\Theta) - H'(r)(\varrho\Theta - r) - H(r)}{\varepsilon^2} \right] dx \end{aligned}$$

$$H(Z) = \frac{1}{\gamma - 1} Z^\gamma$$

Test functions

$$(r - \bar{\varrho}) \in C_c^\infty([0, T] \times R^3), \quad r > 0, \quad \mathbf{U} \in C_c^\infty([0, T] \times R^3; R^3)$$

Relative energy inequality

Inequality

$$\begin{aligned} & \left[\mathcal{E}(\varrho, \Theta, \mathbf{u} \mid r, \mathbf{U}) \right]_{t=0}^{t=\tau} + \varepsilon^\alpha \int_0^\tau \int_{R^3} \mathbb{S}(\nabla_x(\mathbf{u} - \mathbf{U})) : \nabla_x(\mathbf{u} - \mathbf{U}) \, dx dt \\ & \leq \int_0^\tau \int_{R^3} \varrho (\partial_t \mathbf{U} + \mathbf{u} \cdot \nabla_x \mathbf{U}) \cdot (\mathbf{U} - \mathbf{u}) + \varepsilon^\alpha \mathbb{S}(\nabla_x \mathbf{U}) : \nabla_x(\mathbf{U} - \mathbf{u}) \, dx dt \\ & \quad + \frac{1}{\varepsilon^2} \int_0^\tau \int_{R^3} \left[(r - \varrho \Theta) \partial_t H'(r) + \nabla_x H'(r) \cdot (r \mathbf{U} - \varrho \Theta \mathbf{u}) \right] \, dx dt \\ & \quad - \frac{1}{\varepsilon^2} \int_0^\tau \int_{R^3} \left[\operatorname{div}_x \mathbf{U} \left((\varrho \Theta)^\gamma - r^\gamma \right) + \varrho \nabla_x F \cdot (\mathbf{U} - \mathbf{u}) \right] \, dx dt \end{aligned}$$

Ansatz

$$r = r_\varepsilon = \varrho_0 + \underbrace{\varepsilon \mathcal{S}_\varepsilon}_{\text{acoustic component}}, \quad \mathbf{U} = \mathbf{U}_\varepsilon = \mathbf{V} + \underbrace{\nabla_x \Phi_\varepsilon}_{\text{acoustic component}}$$