

# Bound from below of the exterior energy for the wave equation and applications

Thomas Duyckaerts<sup>1</sup> (avec H. Jia<sup>2</sup>, C. Kenig<sup>3</sup> et F. Merle<sup>4</sup>)

<sup>1</sup>Institut Galilée, Université Paris 13, Sorbonne Paris Cité Université et IUF

<sup>2</sup>IAS, Princeton

<sup>3</sup>University of Chicago

<sup>4</sup>Université de Cergy-Pontoise et IHES

CIRM. June 12th, 2017

## 1 Focusing critical wave equation

# Outline

- 1 Focusing critical wave equation
- 2 Energy equirepartition

# Outline

- 1 Focusing critical wave equation
- 2 Energy equirepartition
- 3 Improved energy equirepartition

# Outline

- 1 Focusing critical wave equation
- 2 Energy equirepartition
- 3 Improved energy equirepartition
- 4 Well-prepared initial data

# Outline

- 1 Focusing critical wave equation
- 2 Energy equirepartition
- 3 Improved energy equirepartition
- 4 Well-prepared initial data

# Equation

$$(NLW) \quad \begin{cases} \partial_t^2 u - \Delta u = |u|^{\frac{4}{N-2}} u \\ \vec{u}|_{t=0} = (u_0, u_1) \in \mathcal{H} = \dot{H}^1(\mathbb{R}^N) \times L^2(\mathbb{R}^N) \end{cases}$$

where  $u : [0, T[ \times \mathbb{R}^N \rightarrow \mathbb{R}$ ,  $N \geq 3$ .

# Equation

$$(NLW) \quad \begin{cases} \partial_t^2 u - \Delta u = |u|^{\frac{4}{N-2}} u \\ \vec{u}|_{t=0} = (u_0, u_1) \in \mathcal{H} = \dot{H}^1(\mathbb{R}^N) \times L^2(\mathbb{R}^N) \end{cases}$$

where  $u : [0, T[ \times \mathbb{R}^N \rightarrow \mathbb{R}$ ,  $N \geq 3$ . Well-posedness in  $\mathcal{H} = \dot{H}^1 \times L^2$   
[Ginibre-Velo].



# Equation

$$(NLW) \quad \begin{cases} \partial_t^2 u - \Delta u = |u|^{\frac{4}{N-2}} u \\ \vec{u}|_{t=0} = (u_0, u_1) \in \mathcal{H} = \dot{H}^1(\mathbb{R}^N) \times L^2(\mathbb{R}^N) \end{cases}$$

where  $u : [0, T[ \times \mathbb{R}^N \rightarrow \mathbb{R}$ ,  $N \geq 3$ . Well-posedness in  $\mathcal{H} = \dot{H}^1 \times L^2$   
[Ginibre-Velo].

Conserved energy

$$E(\vec{u}) = \frac{1}{2} \int_{\mathbb{R}^N} |\nabla_x u(t)|^2 + \frac{1}{2} \int_{\mathbb{R}^N} |\partial_t u(t)|^2 - \frac{N-2}{2N} \int_{\mathbb{R}^N} |u(t)|^{\frac{2N}{N-2}}.$$

# Equation

$$(NLW) \quad \begin{cases} \partial_t^2 u - \Delta u = |u|^{\frac{4}{N-2}} u \\ \vec{u}|_{t=0} = (u_0, u_1) \in \mathcal{H} = \dot{H}^1(\mathbb{R}^N) \times L^2(\mathbb{R}^N) \end{cases}$$

where  $u : [0, T[ \times \mathbb{R}^N \rightarrow \mathbb{R}$ ,  $N \geq 3$ . Well-posedness in  $\mathcal{H} = \dot{H}^1 \times L^2$   
[Ginibre-Velo].

Conserved energy

$$E(\vec{u}) = \frac{1}{2} \int_{\mathbb{R}^N} |\nabla_x u(t)|^2 + \frac{1}{2} \int_{\mathbb{R}^N} |\partial_t u(t)|^2 - \frac{N-2}{2N} \int_{\mathbb{R}^N} |u(t)|^{\frac{2N}{N-2}}.$$

Scaling:  $u_\lambda(t, x) = \frac{1}{\lambda^{N/2-1}} u\left(\frac{t}{\lambda}, \frac{x}{\lambda}\right).$

The energy and the  $\mathcal{H}$ -norm are scale invariant.

# Ground state

*Stationary solutions of (NLW):*

$$(E) \quad -\Delta Q = |Q|^{\frac{4}{N-2}} Q, \quad Q: \mathbb{R}^N \rightarrow \mathbb{R}, \quad Q \in \dot{H}^1(\mathbb{R}^N).$$

“Unique” Minimal energy solution of (E) (ground state):

$$W(x) = \left( 1 + \frac{|x|^2}{N(N-2)} \right)^{1-\frac{N}{2}}$$

# Ground state

*Stationary solutions of (NLW):*

$$(E) \quad -\Delta Q = |Q|^{\frac{4}{N-2}} Q, \quad Q: \mathbb{R}^N \rightarrow \mathbb{R}, \quad Q \in \dot{H}^1(\mathbb{R}^N).$$

“Unique” Minimal energy solution of (E) (ground state):

$$W(x) = \left(1 + \frac{|x|^2}{N(N-2)}\right)^{1-\frac{N}{2}}$$

Threshold for the dynamics: [Kenig-Merle 2008].

$$T_+(u) < \infty \implies \limsup_{t \rightarrow T_+} \|\nabla u\|_{L^2}^2 + \frac{N-2}{2} \|\partial_t u\|_{L^2}^2 \geq \|\nabla W\|_{L^2}^2$$

Existence of solutions of (E) with arbitrary large energy: [W.Y. Ding 1986], [Del Pino, Musso, Pacard, Pistoia 2013].

## Other examples of solutions

Solitary waves (**solitons**): if  $\mathbf{p} \in \mathbb{R}^3$  et  $\rho = |\mathbf{p}| < 1$  and  $Q$  is a solution of (E):

$$Q_{\mathbf{p}}(t, x) = Q \left( \left( -\frac{t}{\sqrt{1-\rho^2}} + \frac{1}{\rho^2} \left( \frac{1}{\sqrt{1-\rho^2}} - 1 \right) \mathbf{p} \cdot x \right) \mathbf{p} + x \right)$$
$$Q_{\mathbf{p}}(t, x) = Q_{\mathbf{p}}(0, x - t\mathbf{p}).$$

## Other examples of solutions

Solitary waves (**solitons**): if  $\mathbf{p} \in \mathbb{R}^3$  et  $\rho = |\mathbf{p}| < 1$  and  $Q$  is a solution of (E):

$$Q_{\mathbf{p}}(t, x) = Q \left( \left( -\frac{t}{\sqrt{1-\rho^2}} + \frac{1}{\rho^2} \left( \frac{1}{\sqrt{1-\rho^2}} - 1 \right) \mathbf{p} \cdot x \right) \mathbf{p} + x \right)$$
$$Q_{\mathbf{p}}(t, x) = Q_{\mathbf{p}}(0, x - t\mathbf{p}).$$

Type II **Blow-up solutions**:

$$\vec{u}(t) = \left( \frac{1}{\lambda(t)^{\frac{N-2}{2}}} W \left( \frac{\cdot}{\lambda(t)} \right), 0 \right) + (v_0, v_1) + o(1), \quad t \rightarrow T_+,$$

where  $(v_0, v_1) \in \mathcal{H}$  and  $\lambda(t) \ll T_+ - t$ , see [Krieger Schlag & Tataru 09]. See also [Hillairet & Raphaël 2012, Krieger & Schlag 2014, Jendrej 2015].

## Other examples of solutions

Solitary waves (**solitons**): if  $\mathbf{p} \in \mathbb{R}^3$  et  $\rho = |\mathbf{p}| < 1$  and  $Q$  is a solution of (E):

$$Q_{\mathbf{p}}(t, x) = Q \left( \left( -\frac{t}{\sqrt{1-\rho^2}} + \frac{1}{\rho^2} \left( \frac{1}{\sqrt{1-\rho^2}} - 1 \right) \mathbf{p} \cdot x \right) \mathbf{p} + x \right)$$
$$Q_{\mathbf{p}}(t, x) = Q_{\mathbf{p}}(0, x - t\mathbf{p}).$$

Type II **Blow-up solutions**:

$$\vec{u}(t) = \left( \frac{1}{\lambda(t)^{\frac{N-2}{2}}} W \left( \frac{\cdot}{\lambda(t)} \right), 0 \right) + (v_0, v_1) + o(1), \quad t \rightarrow T_+,$$

where  $(v_0, v_1) \in \mathcal{H}$  and  $\lambda(t) \ll T_+ - t$ , see [Krieger Schlag & Tataru 09]. See also [Hillairet & Raphaël 2012, Krieger & Schlag 2014, Jendrej 2015].

Open questions: solutions with more than one bubbles, or other bubbles than the ground state? See [Jendrej], [Martel & Merle] and also [Côte & Zaag 2012], [Côte & Martel].

# Generalities on type II blow-up

Let  $u$  be a solution of (NLW) such that

$$T_+ = T_+(u) < \infty \text{ and } \limsup_{t \rightarrow T_+} \|\nabla u\|_{L^2}^2 + \|\partial_t u\|_{L^2}^2 < \infty.$$

Then there exist  $k \geq 1$ ,  $k$  **blow-up** points  $(x_1, \dots, x_k) \in (\mathbb{R}^N)^k$  and  $(v_0, v_1) \in \mathcal{H}$  such that

$$\vec{u}(t) \xrightarrow[t \rightarrow T_+]{} (v_0, v_1)$$

and, letting  $\mathcal{R}_t = \left\{ x \in \mathbb{R}^N : \forall j \in \{1, \dots, k\}, |x - x_j| > T_+ - t \right\}$ ,

$$\lim_{t \rightarrow T_+} \int_{\mathcal{R}_t} |\nabla u(t, x) - \nabla v_0(x)|^2 dx + \int_{\mathcal{R}_t} |\partial_t u(t, x) - v_1(x)|^2 dx = 0.$$



# Outline

1 Focusing critical wave equation

**2 Energy equirepartition**

3 Improved energy equirepartition

4 Well-prepared initial data

# Linear wave equation

The asymptotic behavior for solutions of the **linear** wave equation:

$$(LW) \quad \begin{cases} \partial_t^2 u_L - \Delta u_L = 0, & x \in \mathbb{R}^N \\ \vec{u}_L|_{t=0} = (u_0, u_1) \in \mathcal{H} = \dot{H}^1(\mathbb{R}^N) \times L^2(\mathbb{R}^N). \end{cases}$$

is well-known.

# Linear wave equation

The asymptotic behavior for solutions of the **linear** wave equation:

$$(LW) \quad \begin{cases} \partial_t^2 u_L - \Delta u_L = 0, & x \in \mathbb{R}^N \\ \vec{u}_L|_{t=0} = (u_0, u_1) \in \mathcal{H} = \dot{H}^1(\mathbb{R}^N) \times L^2(\mathbb{R}^N). \end{cases}$$

is well-known. Let  $\partial$  be the tangential derivative. Then:

$$\lim_{t \rightarrow +\infty} \int \frac{1}{|x|^2} |u_L(t, x)|^2 + |\partial u_L(t, x)|^2 + |u_L(t, x)|^{\frac{2N}{N-2}} dx = 0$$

# Linear wave equation

The asymptotic behavior for solutions of the **linear** wave equation:

$$(LW) \quad \begin{cases} \partial_t^2 u_L - \Delta u_L = 0, & x \in \mathbb{R}^N \\ \vec{u}_L|_{t=0} = (u_0, u_1) \in \mathcal{H} = \dot{H}^1(\mathbb{R}^N) \times L^2(\mathbb{R}^N). \end{cases}$$

is well-known. Let  $\partial$  be the tangential derivative. Then:

$$\lim_{t \rightarrow +\infty} \int \frac{1}{|x|^2} |u_L(t, x)|^2 + |\partial u_L(t, x)|^2 + |u_L(t, x)|^{\frac{2N}{N-2}} dx = 0$$

and (see **[Friedlander 70s]**) there exists  $G_{\pm} \in L^2(\mathbb{R} \times S^{N-1})$  such that

$$\begin{aligned} \lim_{t \rightarrow +\infty} \int_0^{+\infty} \int_{S^{N-1}} \left| r^{\frac{N-1}{2}} \partial_r u_L(t, r\omega) \mp G_{\pm}(r-t, \omega) \right|^2 \\ + \left| r^{\frac{N-1}{2}} \partial_t u_L(t, r\omega) + G_{\pm}(r-t, \omega) \right|^2 dr d\omega = 0. \end{aligned}$$

# Equirepartition for the linear equation

**Theorem** [TD, Kenig, Merle 2012]. *Assume that  $N$  is odd. Let  $u_L$  be a solution of the linear wave equation. Then the following holds for all  $t \geq 0$  or for all  $t \leq 0$ :*

$$\int_{|x| \geq |t|} |\nabla_{t,x} u_L(t, x)|^2 dx \geq \frac{1}{2} \int_{\mathbb{R}^N} |\nabla_{t,x} u_L(0, x)|^2 dx.$$

# Equirepartition for the linear equation

**Theorem** [TD, Kenig, Merle 2012]. *Assume that  $N$  is odd. Let  $u_L$  be a solution of the linear wave equation. Then the following holds for all  $t \geq 0$  or for all  $t \leq 0$ :*

$$\int_{|x| \geq |t|} |\nabla_{t,x} u_L(t, x)|^2 dx \geq \frac{1}{2} \int_{\mathbb{R}^N} |\nabla_{t,x} u_L(0, x)|^2 dx.$$

Proof by a symmetry argument, using the explicit solution.

# Equirepartition for the linear equation

**Theorem** [TD, Kenig, Merle 2012]. Assume that  $N$  is odd. Let  $u_L$  be a solution of the linear wave equation. Then the following holds for all  $t \geq 0$  or for all  $t \leq 0$ :

$$\int_{|x| \geq |t|} |\nabla_{t,x} u_L(t, x)|^2 dx \geq \frac{1}{2} \int_{\mathbb{R}^N} |\nabla_{t,x} u_L(0, x)|^2 dx.$$

Proof by a symmetry argument, using the explicit solution.  
Does not hold in even dimension [Côte, Kenig, Schlag 2014].

# Equirepartition for the linear equation

**Theorem** [TD, Kenig, Merle 2012]. Assume that  $N$  is odd. Let  $u_L$  be a solution of the linear wave equation. Then the following holds for all  $t \geq 0$  or for all  $t \leq 0$ :

$$\int_{|x| \geq |t|} |\nabla_{t,x} u_L(t, x)|^2 dx \geq \frac{1}{2} \int_{\mathbb{R}^N} |\nabla_{t,x} u_L(0, x)|^2 dx.$$

Proof by a symmetry argument, using the explicit solution.  
Does not hold in even dimension [Côte, Kenig, Schlag 2014].

**Question:** for which solutions of (NLW) does there exist  $\eta > 0$  such that

$$\forall t \geq 0 \text{ ou } \forall t \leq 0, \quad \int_{|x| \geq |t|} |\nabla_{t,x} u(t, x)|^2 dx \geq \eta?$$



# Equipartition for critical focusing wave

**Théorème** [TD, Kenig, Merle 2012]. Assume  $N$  is odd. There exists  $\varepsilon_0 > 0$  such that if  $u$  is a solution of (NLW) with:

$$\|(u_0, u_1)\|_{\mathcal{H}} < \varepsilon_0,$$

then the following holds for all  $t \geq 0$  or for all  $t \leq 0$ :

$$\int_{|x| \geq |t|} |\nabla_{t,x} u(t, x)|^2 dx \geq \frac{1}{4} \int_{\mathbb{R}^N} |\nabla_{t,x} u(0, x)|^2 dx.$$

# Application

**Theorem.** There exists  $\varepsilon_0 > 0$  such that, for any solution  $u$  of (NLW) such that  $T_+(u) < \infty$  and

$$\limsup_{t \rightarrow T_+} \|\nabla u\|_{L^2}^2 + \frac{N-2}{2} \|\partial_t u\|_{L^2}^2 \leq \|\nabla W\|^2 + \varepsilon_0,$$

there exists  $x_0 \in \mathbb{R}^N$ ,  $(v_0, v_1) \in \mathcal{H}$ ,  $\mathbf{p} \in \mathbb{R}^N$ ,  $\lambda(t)$  and  $x(t)$  such that

$$(u(t), \partial_t u(t)) - (v_0, v_1) - \left( \frac{\iota_0}{\lambda(t)^{\frac{N}{2}-1}} W_{\mathbf{p}} \left( 0, \frac{\cdot - x(t)}{\lambda(t)} \right), \frac{\iota_0}{\lambda(t)^{\frac{N}{2}}} (\partial_t W_{\mathbf{p}}) \left( 0, \frac{\cdot - x(t)}{\lambda(t)} \right) \right) \xrightarrow{t \rightarrow T_+} 0$$

in  $\mathcal{H}$  and

$$\lim_{t \rightarrow T_+} \frac{\lambda(t)}{T_+ - t} = 0, \quad \lim_{t \rightarrow T_+} \frac{x(t) - x_0}{T_+ - t} = \mathbf{p}, \quad |\mathbf{p}| \leq C\varepsilon_0^{1/4}.$$

# Outline

- 1 Focusing critical wave equation
- 2 Energy equirepartition
- 3 Improved energy equirepartition**
- 4 Well-prepared initial data

# Exterior energy for radial data in 3D

**Proposition.** Let  $u_L$  be a radial solution of the linear wave equation in *space dimension 3*. Let  $A > 0$ . Assume  $u_0 \perp \frac{1}{r}$  in  $\dot{H}^1(\{r > A\})$ , i.e.  $u_0(A) = 0$ . Then:

$$\forall t \geq 0 \text{ or } \forall t \leq 0, \quad \int_{A+|t|}^{+\infty} (\partial_{t,r} u_L(t, r))^2 r^2 dr \geq \frac{1}{2} \int_A^{+\infty} (\partial_{t,r} u_L(0, r))^2 r^2 dr.$$

# Exterior energy for radial data in 3D

**Proposition.** Let  $u_L$  be a radial solution of the linear wave equation in space dimension 3. Let  $A > 0$ . Assume  $u_0 \perp \frac{1}{r}$  in  $\dot{H}^1(\{r > A\})$ , i.e.  $u_0(A) = 0$ . Then:

$$\forall t \geq 0 \text{ or } \forall t \leq 0, \quad \int_{A+|t|}^{+\infty} (\partial_{t,r} u_L(t, r))^2 r^2 dr \geq \frac{1}{2} \int_A^{+\infty} (\partial_{t,r} u_L(0, r))^2 r^2 dr.$$

Generalization to other odd dimensions: [Kenig, Lawrie, Baoping Liu, Schlag 2015]

# Rigidity theorem

**Theorem** Assume  $N = 3$ . Let  $u$  be a global and radial solution of (NLW). Assume

$$\forall A > 0, \quad \liminf_{t \rightarrow \pm\infty} \int_{|x| \geq A+|t|} |\nabla_{t,x} u|^2 dx = 0.$$

Then  $u = 0$  or there exist  $\lambda > 0$ ,  $\iota \in \{\pm 1\}$  such that  $u(t, x) = \frac{\iota}{\lambda^{1/2}} W\left(\frac{x}{\lambda}\right)$ .

# Rigidity theorem

**Theorem** Assume  $N = 3$ . Let  $u$  be a global and radial solution of (NLW). Assume

$$\forall A > 0, \quad \liminf_{t \rightarrow \pm\infty} \int_{|x| \geq A+|t|} |\nabla_{t,x} u|^2 dx = 0.$$

Then  $u = 0$  or there exist  $\lambda > 0$ ,  $\iota \in \{\pm 1\}$  such that  $u(t, x) = \frac{\iota}{\lambda^{1/2}} W\left(\frac{x}{\lambda}\right)$ .

Recall that  $W(x) = \frac{1}{\left(1 + \frac{|x|^2}{3}\right)^{\frac{1}{2}}}$ , so that

$$\frac{\iota}{\lambda^{1/2}} W\left(\frac{x}{\lambda}\right) \approx \frac{\sqrt{3}\lambda^{1/2}}{|x|}, \quad |x| \rightarrow \infty.$$

First step of the proof:  $\exists \ell \in \mathbb{R}$  such that

$$\lim_{r \rightarrow \infty} ru_0(r) = \ell.$$

# Classification of type II blow-up solution

**Theorem.** Assume  $N = 3$ . Let  $u$  be a radial solution of (NLW) such that  $T_+(u) < +\infty$ . Then

$$\lim_{t \rightarrow T_+(u)} \|\vec{u}(t)\|_{\mathcal{H}} = +\infty$$

or there exist  $J \geq 1$ ,

- $(v_0, v_1) \in \mathcal{H}$ ,
- signs  $\iota_j \in \{\pm 1\}$ ,  $j = 1 \dots J$ ,
- parameters  $\lambda_j(t)$ ,  $0 < \lambda_1(t) \ll \lambda_2(t) \ll \dots \ll \lambda_J(t) \ll T_+ - t$ ,

such that



# Classification of type II blow-up solution

**Theorem.** Assume  $N = 3$ . Let  $u$  be a radial solution of (NLW) such that  $T_+(u) < +\infty$ . Then

$$\lim_{t \rightarrow T_+(u)} \|\vec{u}(t)\|_{\mathcal{H}} = +\infty$$

or there exist  $J \geq 1$ ,

- $(v_0, v_1) \in \mathcal{H}$ ,
- signs  $\iota_j \in \{\pm 1\}$ ,  $j = 1 \dots J$ ,
- parameters  $\lambda_j(t)$ ,  $0 < \lambda_1(t) \ll \lambda_2(t) \ll \dots \ll \lambda_J(t) \ll T_+ - t$ ,

such that

$$\vec{u}(t) = (v_0, v_1) + \sum_{j=1}^J \left( \frac{\iota_j}{\lambda_j^{\frac{1}{2}}(t)} W \left( \frac{x}{\lambda_j(t)} \right), 0 \right) + \vec{\varepsilon}(t),$$

where  $\lim_{t \rightarrow T_+} \|\vec{\varepsilon}(t)\|_{\mathcal{H}} = 0$ .

# Outline

1 Focusing critical wave equation

2 Energy equirepartition

3 Improved energy equirepartition

4 Well-prepared initial data

# Lower bound for the exterior energy

**Lemma.** Let  $\gamma \in (0, 1)$ . There exists  $\varepsilon = \varepsilon(\gamma) > 0$  with the following property. Let  $u_L$  be a solution of (LW) with initial data  $(u_0, u_1)$  such that

$$\begin{cases} (u_0, u_1) \in \dot{H}^1 \times L^2 & \text{if } N \geq 3 \\ |\nabla u_0| \in L^2, u_1 \in L^2 \text{ and } u_0 \equiv u_\infty & \text{si } N = 2 \end{cases}$$

(where  $u_\infty \in \mathbb{R}$ ) and

$$\|(\nabla u_0, u_1)\|_{L^2(B_{1+\varepsilon}^c \cup B_{1-\varepsilon})} + \|\partial u_0\|_{L^2} + \|\partial_r u_0 + u_1\|_{L^2} \leq \varepsilon \|(\nabla u_0, u_1)\|_{L^2}.$$

# Lower bound for the exterior energy

**Lemma.** Let  $\gamma \in (0, 1)$ . There exists  $\varepsilon = \varepsilon(\gamma) > 0$  with the following property. Let  $u_L$  be a solution of (LW) with initial data  $(u_0, u_1)$  such that

$$\begin{cases} (u_0, u_1) \in \dot{H}^1 \times L^2 & \text{if } N \geq 3 \\ |\nabla u_0| \in L^2, u_1 \in L^2 \text{ and } u_0 \equiv u_\infty & \text{si } N = 2 \end{cases}$$

(where  $u_\infty \in \mathbb{R}$ ) and

$$\|(\nabla u_0, u_1)\|_{L^2(B_{1+\varepsilon}^c \cup B_{1-\varepsilon})} + \|\partial_t u_0\|_{L^2} + \|\partial_r u_0 + u_1\|_{L^2} \leq \varepsilon \|(\nabla u_0, u_1)\|_{L^2}.$$

Then, for all  $t \geq 0$ ,

$$\int_{|x| \geq \gamma+t} |\nabla_{x,t} u_L|^2(x, t) dx \geq \gamma \|(\nabla u_0, u_1)\|_{L^2}^2.$$

# Lower bound for the exterior energy

**Lemma.** Let  $\gamma \in (0, 1)$ . There exists  $\varepsilon = \varepsilon(\gamma) > 0$  with the following property. Let  $u_L$  be a solution of (LW) with initial data  $(u_0, u_1)$  such that

$$\begin{cases} (u_0, u_1) \in \dot{H}^1 \times L^2 & \text{if } N \geq 3 \\ |\nabla u_0| \in L^2, u_1 \in L^2 \text{ and } u_0 \equiv u_\infty & \text{si } N = 2 \end{cases}$$

(where  $u_\infty \in \mathbb{R}$ ) and

$$\|(\nabla u_0, u_1)\|_{L^2(B_{1+\varepsilon}^c \cup B_{1-\varepsilon})} + \|\partial_t u_0\|_{L^2} + \|\partial_r u_0 + u_1\|_{L^2} \leq \varepsilon \|(\nabla u_0, u_1)\|_{L^2}.$$

Then, for all  $t \geq 0$ ,

$$\int_{|x| \geq \gamma+t} |\nabla_{x,t} u_L|^2(x, t) dx \geq \gamma \|(\nabla u_0, u_1)\|_{L^2}^2.$$

**Application for critical semilinear wave equation:** lower bound of the exterior energy for well-prepared initial data, and soliton resolution along a sequence of times.

# Wave maps

$$(WM) \quad \begin{cases} \partial_t^2 u - \Delta u = (|\nabla u|^2 - |\partial_t u|^2) u, & x \in \mathbb{R}^2 \\ \vec{u}|_{t=0} = (u_0, u_1), & u_0 \cdot u_1 = 0. \end{cases}$$

$$u : [0, T[ \times \mathbb{R}^2 \rightarrow \mathbb{S}^2.$$

“Well-posedness” in  $\mathcal{H} = \dot{H}^1(\mathbb{R}^2) \times L^2(\mathbb{R}^2)$  [Tao, Tataru].

To fix ideas, consider classical solutions:  $(u_0, u_1) \in C^\infty$ ,  $u_0$  constant at infinity,  $u_1 \equiv 0$  at infinity.

# Wave maps

$$(WM) \quad \begin{cases} \partial_t^2 u - \Delta u = (|\nabla u|^2 - |\partial_t u|^2) u, & x \in \mathbb{R}^2 \\ \vec{u}|_{t=0} = (u_0, u_1), & u_0 \cdot u_1 = 0. \end{cases}$$

$$u : [0, T[ \times \mathbb{R}^2 \rightarrow \mathbb{S}^2.$$

“Well-posedness” in  $\mathcal{H} = \dot{H}^1(\mathbb{R}^2) \times L^2(\mathbb{R}^2)$  [Tao, Tataru].

To fix ideas, consider classical solutions:  $(u_0, u_1) \in C^\infty$ ,  $u_0$  constant at infinity,  $u_1 \equiv 0$  at infinity.

The **energy**

$$E_M(\vec{u}) = \frac{1}{2} \int_{\mathbb{R}^2} |\nabla_x u(t)|^2 + \frac{1}{2} \int_{\mathbb{R}^2} |\partial_t u(t)|^2$$

is conserved.

# Wave maps

$$(WM) \quad \begin{cases} \partial_t^2 u - \Delta u = (|\nabla u|^2 - |\partial_t u|^2) u, & x \in \mathbb{R}^2 \\ \vec{u}|_{t=0} = (u_0, u_1), & u_0 \cdot u_1 = 0. \end{cases}$$

$$u : [0, T[ \times \mathbb{R}^2 \rightarrow \mathbb{S}^2.$$

“Well-posedness” in  $\mathcal{H} = \dot{H}^1(\mathbb{R}^2) \times L^2(\mathbb{R}^2)$  [Tao, Tataru].

To fix ideas, consider classical solutions:  $(u_0, u_1) \in C^\infty$ ,  $u_0$  constant at infinity,  $u_1 \equiv 0$  at infinity.

The **energy**

$$E_M(\vec{u}) = \frac{1}{2} \int_{\mathbb{R}^2} |\nabla_x u(t)|^2 + \frac{1}{2} \int_{\mathbb{R}^2} |\partial_t u(t)|^2$$

is conserved.

**Scaling:**  $u_\lambda(t, x) = u\left(\frac{t}{\lambda}, \frac{x}{\lambda}\right)$ .

The energy is invariant by the scaling.



# References for Wave Maps.

Global and local Cauchy theory in the critical space: [Tao 2001], [Tataru 2001 & 2005], [Sterbenz & Tataru 2010].

Global existence below the energy of the ground state: [Christodoulou, Tahvildar-Zadeh 1993], [Struwe 2003] (equivariant case), [Sterbenz & Tataru 2010]. See also [Tao], [Krieger & Schlag].

Explicit blow-up solutions: [Krieger, Schlag & Tataru 2008], [Raphaël & Rodnianski 2012], [Jendrej 2016].

Soliton resolution for **equivariant** solutions below a natural threshold: [Côte, Kenig, Lawrie & Schlag 2015].

Soliton resolution along a sequence of times for equivariant solutions: [Côte 2015].

Soliton resolution strictly inside the wave cone for equivariant solutions: [Grinis 2016].

# Well-prepared initial data for wave maps

**Theorem.** *Let  $\gamma \in (0, 1)$ . There exists  $\varepsilon = \varepsilon(\gamma) > 0$  with the following property. Let  $u$  be a classical solution of (WM) with initial data  $(u_0, u_1)$  such that*

$$E_M(u_0, u_1) \leq \varepsilon$$

*and*

$$\|(\nabla u_0, u_1)\|_{L^2(B_{1+\varepsilon}^c \cup B_{1-\varepsilon})} + \|\not\partial u_0\|_{L^2} + \|\partial_r u_0 + u_1\|_{L^2} \leq \varepsilon \|(\nabla u_0, u_1)\|_{L^2}.$$

# Well-prepared initial data for wave maps

**Theorem.** Let  $\gamma \in (0, 1)$ . There exists  $\varepsilon = \varepsilon(\gamma) > 0$  with the following property. Let  $u$  be a classical solution of (WM) with initial data  $(u_0, u_1)$  such that

$$E_M(u_0, u_1) \leq \varepsilon$$

and

$$\|(\nabla u_0, u_1)\|_{L^2(B_{1+\varepsilon}^c \cup B_{1-\varepsilon})} + \|\not\partial u_0\|_{L^2} + \|\partial_r u_0 + u_1\|_{L^2} \leq \varepsilon \|(\nabla u_0, u_1)\|_{L^2}.$$

Then, for all  $t \geq 0$ ,

$$\int_{|x| \geq \gamma+t} |\nabla_{x,t} u|^2(t, x) dx \geq \gamma \|(\nabla u_0, u_1)\|_{L^2}^2.$$

# Small blow-up solutions for wave maps

**Theorem** Let  $u$  be a classical of (WM) such that  $E_M(\vec{u}(0)) < E_M(W, 0) + \epsilon_0^2$ , blowing-up in finite time  $T_+$  at  $x = 0$ . Then  $\exists \mathbf{p} \in \mathbb{R}^2$  such that  $|\mathbf{p}| \ll 1$ ,  $x(t) \in \mathbb{R}^2$ ,  $\lambda(t) > 0$  with

$$\lim_{t \rightarrow T_+} \frac{x(t)}{T_+ - t} = \mathbf{p}, \quad \lim_{t \rightarrow T_+} \frac{\lambda(t)}{T_+ - t} = 0,$$

and  $(v_0, v_1) \in \mathcal{H} \cap C^\infty(\mathbb{R}^2 \setminus \{0\})$  with  $(v_0 - u_\infty, v_1)$  compactly supported, such that

$$(i) \quad \inf \left\{ \left\| \vec{u}(t) - (v_0, v_1) - (Q_{\mathbf{p}}, \partial_t Q_{\mathbf{p}}) \right\|_{\mathcal{H}} : Q_{\mathbf{p}} \in \mathcal{M}_{\mathbf{p}} \right\} \xrightarrow{t \rightarrow T_+} 0,$$

$$(ii) \quad \left\| (\nabla u(t), \partial_t u(t)) - (\nabla v_0, v_1) \right\|_{L^2(\mathbb{R}^2 \setminus B_{\lambda(t)}(x(t)))} \xrightarrow{t \rightarrow T_+} 0,$$

where  $B_{\lambda(t)}(x(t)) = \{x \in \mathbb{R}^2 : |x - x(t)| < \lambda(t)\}$ ,  $\mathcal{M}_{\mathbf{p}}$  is the set of all geometrical transforms of  $W_{\mathbf{p}}$  (space translation, scaling, and  $\mathbb{S}^2$  isometries), and  $W$  is the ground state.