

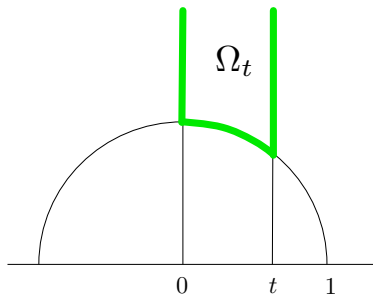
Hyperbolic triangles with no positive Neumann eigenvalues

Luc Hillairet, Université d'Orleans

Chris Judge, Indiana University

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Question



- ▶ Does there exist $u \in L^2\left(\Omega_t, \frac{dx dy}{y^2}\right)$ and $E > 0$ with
 1. $-y^2(\partial_x^2 + \partial_y^2)u = E \cdot u$, and
 2. $\frac{\partial}{\partial n}u \equiv 0$ (Neumann conditions)?
- ▶ Remark: Ω_t non-compact but finite measure.

Some answers

- ▶ Selberg proved 'yes' if $t \in \left\{ \frac{1}{2}, \frac{1}{\sqrt{2}}, \frac{\sqrt{3}}{2} \right\}$
- ▶ Phillips and Sarnak conjectured 'no' if

$$t \in \{ \cos(\pi/n) : n \neq 3, 4, 6 \}$$

Theorem (Hillairet-J.)

The answer is 'no' for all but countably many $t \in]0, 1[$.

Goal of this talk: Describe some ingredients of the proof.

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Spectral dichotomy for hyperbolic surfaces

- ▶ $\mathbb{H} = \{(x, y) \mid y > 0\}$: arclength $y^{-1} ds_{\text{Euc}}$; measure $y^{-2} dx dy$.
- ▶ $\Delta_{\mathbb{H}} = -y^2 (\partial_x^2 + \partial_y^2)$
- ▶ $\Gamma \subset \text{Isom}(\mathbb{H})$ cofinite discrete subgroup with 'cusps'.
- ▶ $\Delta : L^2(\mathbb{H}/\Gamma) \rightarrow L^2(\mathbb{H}/\Gamma)$ non-negative self-adjoint operator.
- ▶ spectral decomposition (Selberg)

$$\Delta u = \sum_{k \in \mathbb{N}} E_k \langle u, u_k \rangle u_k + \int_{\mathbb{R}} (1/4 + r^2) \langle u, G_r \rangle G_r dr$$

- ▶ Conjectured dichotomy (Sarnak):
 Γ 'arithmetic' \Leftrightarrow pure point spectrum has Weyl density 1.

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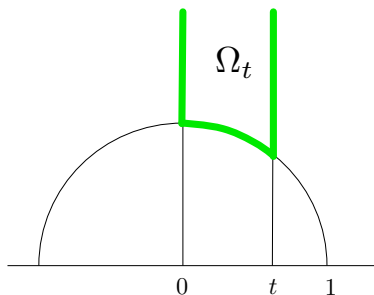
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Return to Ω_t : Fourier decomposition in x



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Return to Ω_t : Fourier decomposition in x

- ▶ pure point vs. continuous spectrum explained by Fourier decomposition:
- ▶ For u smooth Neumann function on Ω_t and $y > 1$,

$$u(x, y) = \sum_{\ell=0}^{\infty} u^{\ell}(y) \cdot \cos\left(\frac{\pi\ell}{t} \cdot x\right)$$

- ▶ If u eigenfunction with eigenvalue E , then u^{ℓ} satisfies the ODE

$$(u^{\ell})'' = \left(\left(\frac{\ell\pi}{t} \right)^2 - \frac{E}{y^2} \right) \cdot u^{\ell}$$

Solutions to the ODE

$$(u^\ell)'' = \left(\left(\frac{\ell\pi}{t} \right)^2 - \frac{E}{y^2} \right) \cdot u^\ell$$

▶ $u^0(y) = y^{\frac{1}{2}} (A \cdot \cos(r \cdot \ln(y)) + B \cdot r^{-1} \cdot \sin(r \cdot \ln(y)))$

$$E = 1/4 + r^2$$

▶ $\ell > 0$: For $y \gg$ zero of $(\ell\pi/t)^2 - E \cdot y^{-2}$

$$u^\ell \sim c \cdot \exp\left(\pm \frac{\ell\pi}{t} \cdot y\right)$$

▶ u is L^2 eigenfunction $\Leftrightarrow u^0 \equiv 0$ and u^ℓ subexp

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Cut-off Laplacian (Lax-Phillips, Colin de Verdiere)

- ▶ Restrict Δ_t to $u \in C_0^\infty$ with $u^0(y) = 0$ for $y \geq b > 1$
- ▶ Friedrichs extension called the *cut-off Laplacian* Δ_t^b .
- ▶ $(\Delta_t^b + Id)^{-1}$ compact, and eigendata real-analytic in t .
- ▶ $\Delta_t u = Eu \implies \Delta_t^b u = Eu$
- ▶ $\Delta_t^b u = Eu$ and $u^0(y) = 0$ for $1 < y < b \implies \Delta_t u = Eu$
- ▶ If $t \mapsto u_t$ analytic, then $t \mapsto u_t^0(y)$ is analytic.
- ▶ Suffices to show $\nexists \Delta_t^b$ eigenbranches u_t with $u_t^0 \equiv 0$

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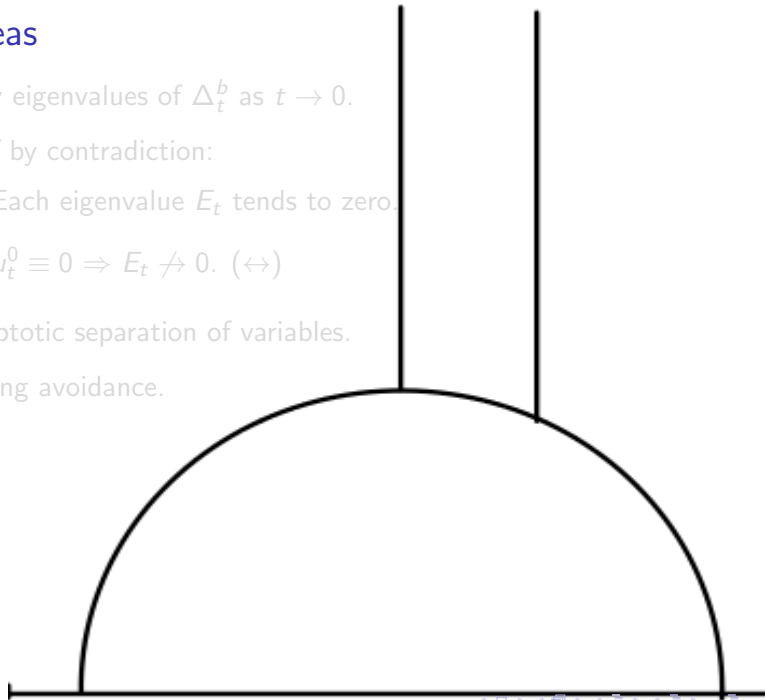
Study eigenvalues of Δ_t^b as $t \rightarrow 0$.

Proof by contradiction:

- ▶ Each eigenvalue E_t tends to zero.
- ▶ $u_t^0 \equiv 0 \Rightarrow E_t \not\rightarrow 0$. (\leftrightarrow)

asymptotic separation of variables.

crossing avoidance.



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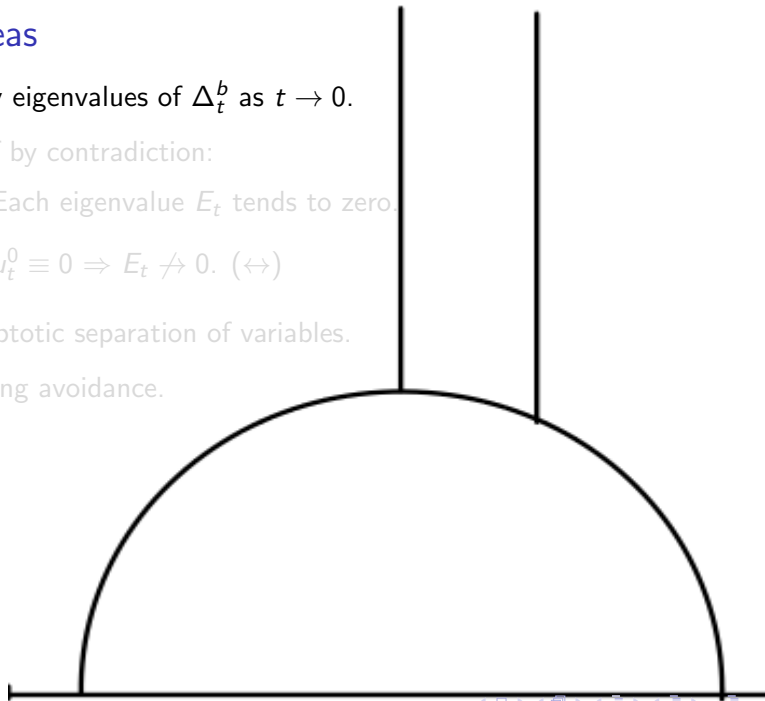
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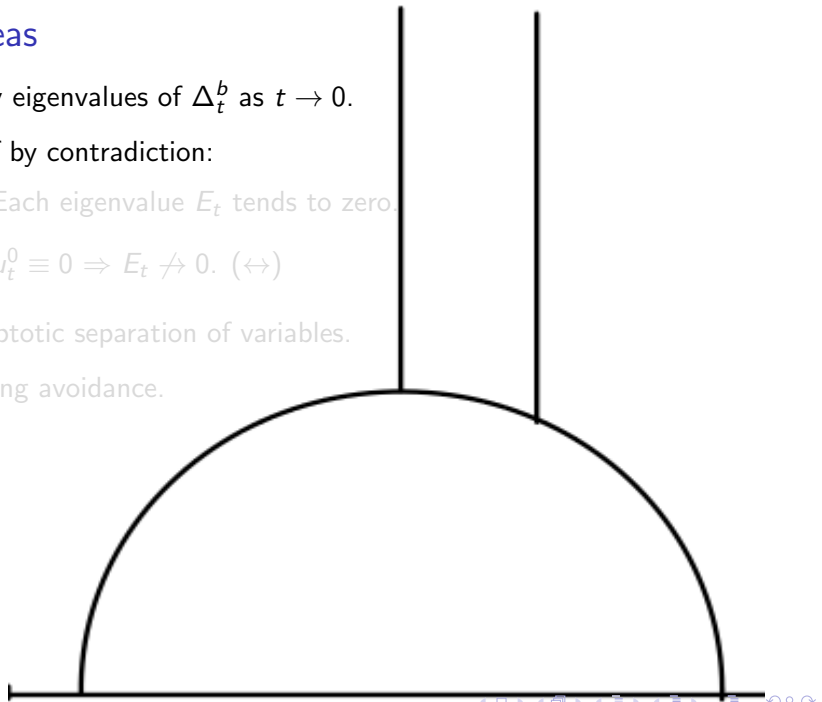
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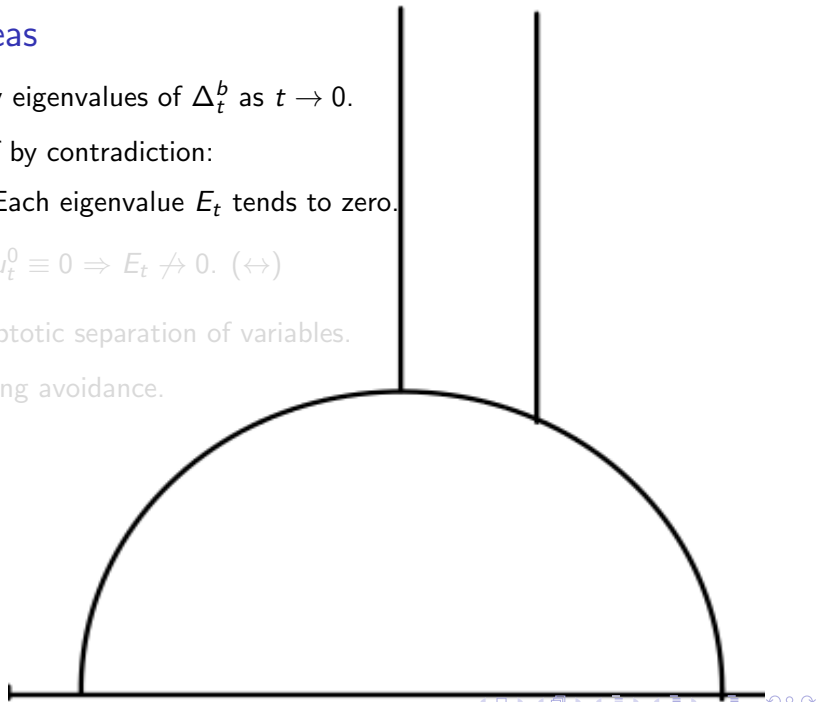
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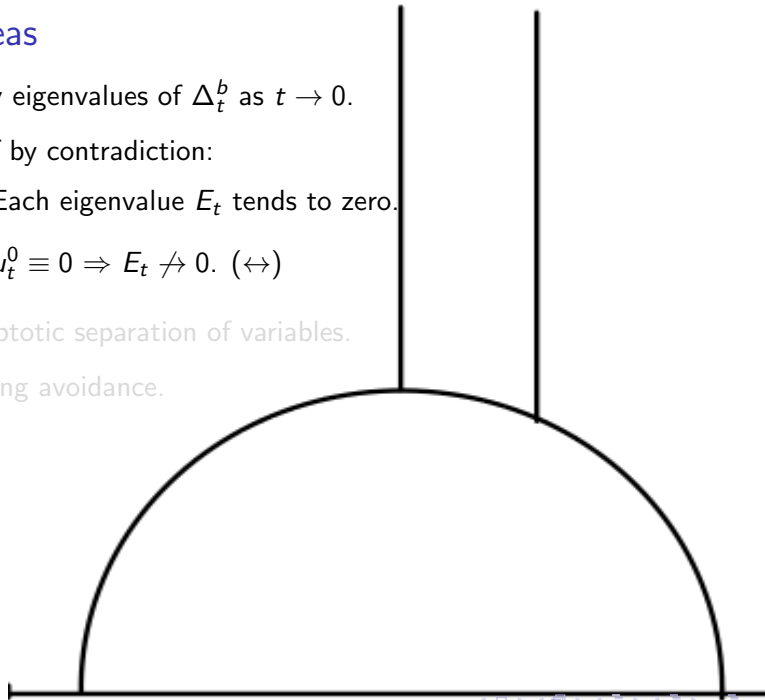
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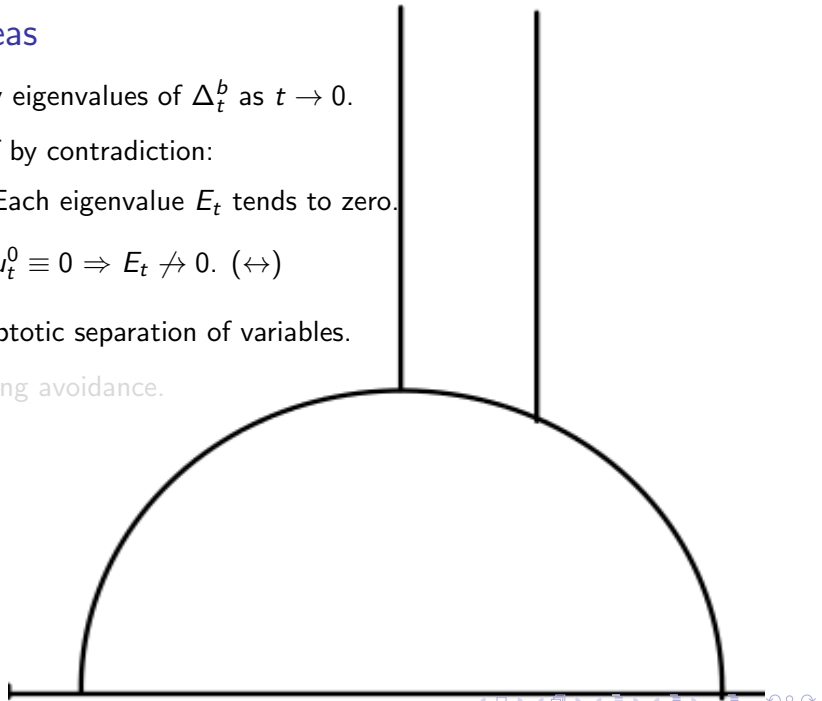
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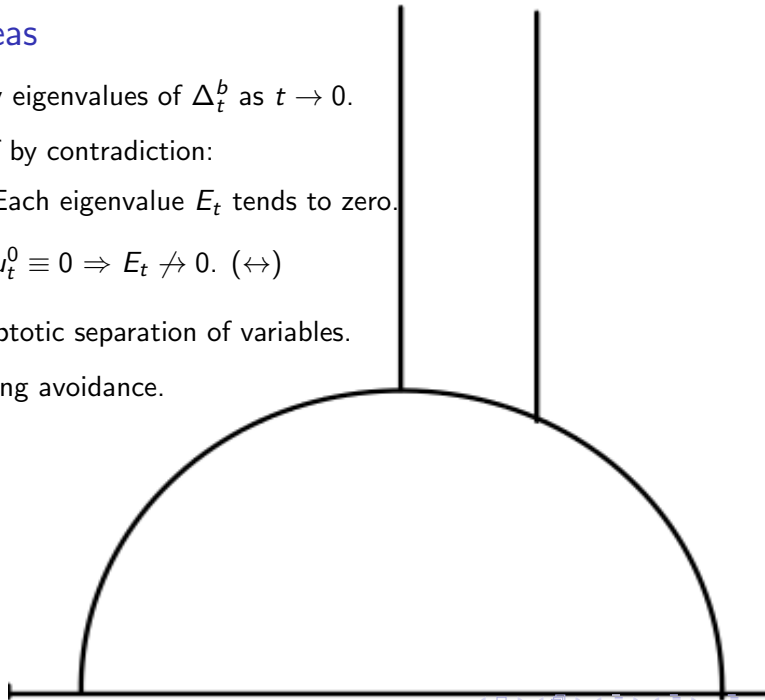
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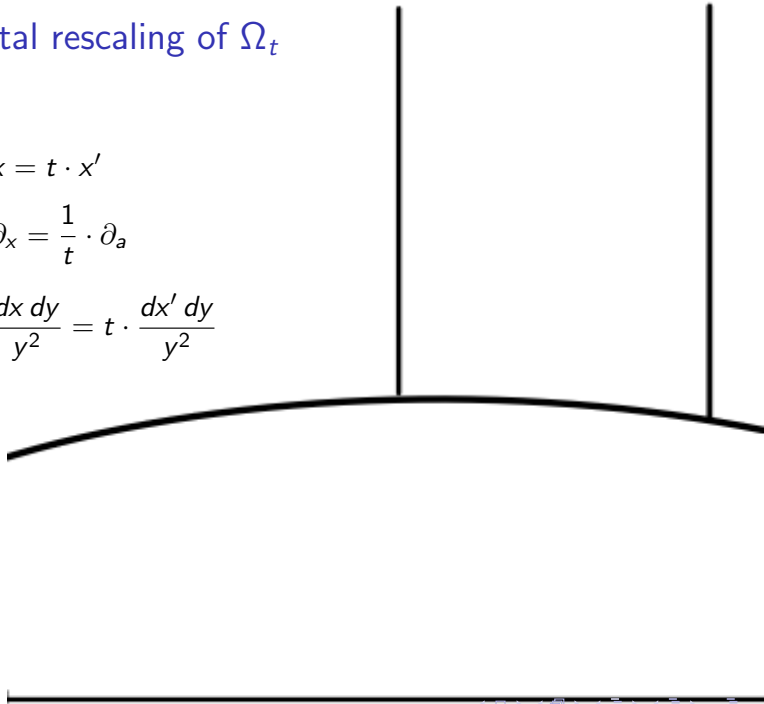


Horizontal rescaling of Ω_t

▶ $x = t \cdot x'$

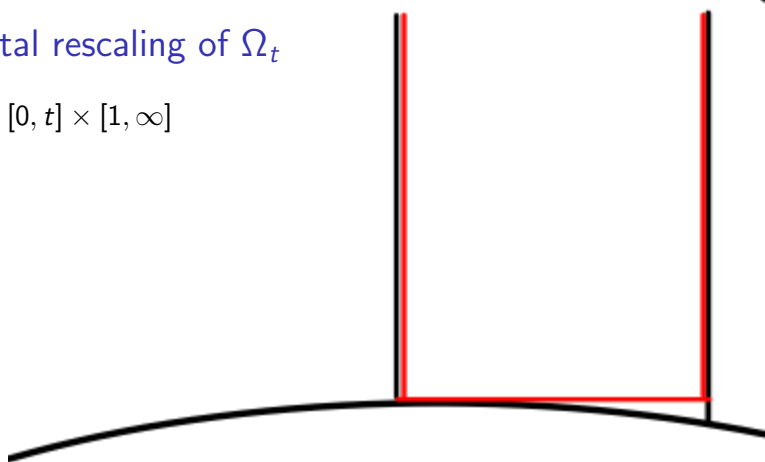
▶ $\partial_x = \frac{1}{t} \cdot \partial_{x'}$

▶ $\frac{dx dy}{y^2} = t \cdot \frac{dx' dy}{y^2}$



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$$\Omega_t \sim [0, t] \times [1, \infty]$$

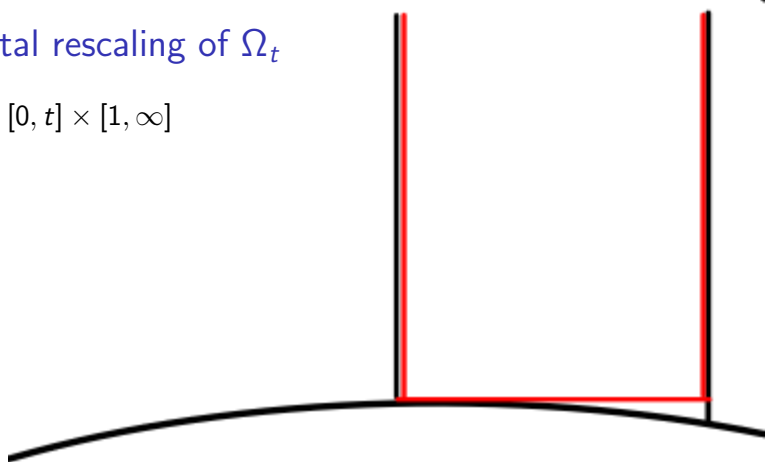


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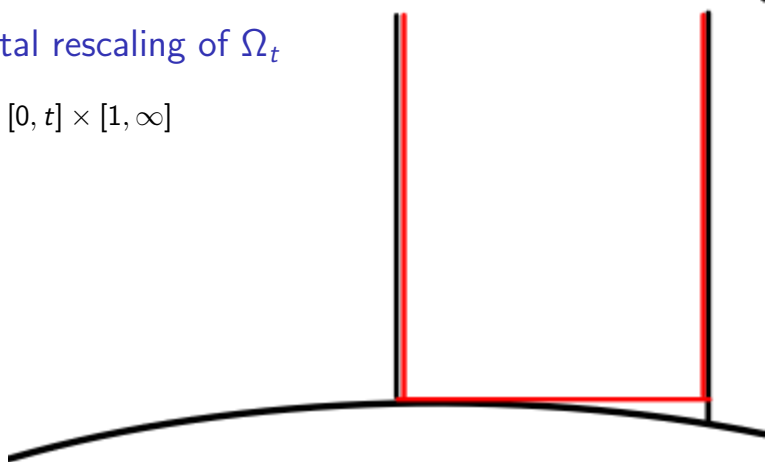


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A tale of two quadratic forms

- ▶ Horizontally stretched domain Ω_t
- ▶ Rectangle $R = [0, 1] \times [1, \infty]$
- ▶ Map R onto Ω_t via diffeo 'supported' on $y < b$.
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- ▶ Expansion at $t = 0$

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- ▶ Expansion at $t = 0$

$$q_t(u) = \int_R \left(\frac{(u_{x'})^2}{t^2} + (u_y)^2 \right) t \, dx \, dy + O(1)$$

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$$q_t(u) = \underbrace{\int_R ((u_x)^2 + t^2 \cdot (u_y)^2) dx dy}_{a_t(u)} + O(t)$$

A tale of two quadratic forms

- ▶ Horizontally stretched domain Ω_t
- ▶ Rectangle $R = [0, 1] \times [1, \infty]$
- ▶ Map R onto Ω_t via diffeo 'supported' on $y < b$.
- ▶ q_t pull-back of t times $\langle \Delta_t \cdot, \cdot \rangle$ to $C_0^\infty(R)$
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$$q_t(u) = \underbrace{\int_R ((u_x)^2 + t^2 \cdot (u_y)^2) dx dy}_{a_t(u)} + O(t \cdot a_t(u))$$

Some generalities: 'Asymptotic at first order'

- ▶ a_t and q_t nonnegative quadratic forms.
- ▶ $|q_t(u, v) - a_t(u, v)| \leq C \cdot t \cdot a_t(u)^{\frac{1}{2}} \cdot a_t(v)^{\frac{1}{2}}$
- ▶ $|\dot{q}_t(u) - \dot{a}_t(u)| \leq C \cdot a_t(u)$ (dot signifies t -derivative)
- ▶ $\dot{a}(u) = 2t \cdot \int u_y^2 dx dy$
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A first consequence: convergence of eigenvalue branches

Proposition

If a_t and q_t asymptotic at first order and $\dot{a}_t \geq 0$, then each real-analytic eigenvalue branch $t \mapsto E_t$ converges as $t \rightarrow 0$.

Proof.

If u_t is the associated eigenfunction branch, then

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for small t because asymptotic at first order.



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Thus,

$$-2C \leq \frac{\dot{E}_t}{E_t}$$

Quasimodes from spectral projections

- ▶ u eigenfunction of q_t with respect to $L^2(dx dy/y^2)$ norm:

$$q_t(u, v) = E \cdot \langle u, v \rangle \text{ for all } v$$

- ▶ $I =$ interval in \mathbb{R} .
- ▶ $w :=$ projection of u onto a_t eigenspaces with $\lambda \in I$

Proposition

w is an a_t quasimode at order t and energy E : For each v

$$|a_t(w, v) - E \cdot \langle w, v \rangle| \leq C \cdot t \cdot \|w_t\| \cdot \|v\|$$

for all v . Moreover,

$$a_t(u - w) + \|w - u\|^2 \leq C \cdot t \cdot \|u\|^2.$$

Consequences for real-analytic eigenbranches of q_t

- ▶ $t \mapsto u_t$ eigenfunction branch of q_t with eigenvalue E_t .
- ▶ w_t corresponding spectral projection associated to $I \subset \mathbb{R}$

Proposition

The function $t \mapsto \frac{\dot{a}(w_t)}{\|w_t\|^2}$ is integrable over $]0, t_0[$

and $\left| \dot{a}(w_t) - \dot{E}_t \cdot \|u_t\|^2 \right| \leq C \cdot \|u_t\|$

Variational formula in Ω_t context

To obtain finer information we use our specific context:

$$a_t(u) = \int (u_x)^2 + t^2 \cdot (u_y)^2$$

$$a_t(u) - E \cdot \|u\|^2 = \int (u_x)^2 + t^2 \cdot (u_y)^2 - E \int \frac{u^2}{y^2}$$

$$t^2 \cdot (u_y)^2 = a_t(u) - E \cdot \|u\|^2 + \int \frac{E}{y^2} \cdot u^2 - (u_x)^2$$

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Spectrum of a_t via separation of variables

▶ $u(x, y) = \sum_{\ell} u^{\ell}(y) \cdot \cos(\pi \cdot \ell \cdot x)$

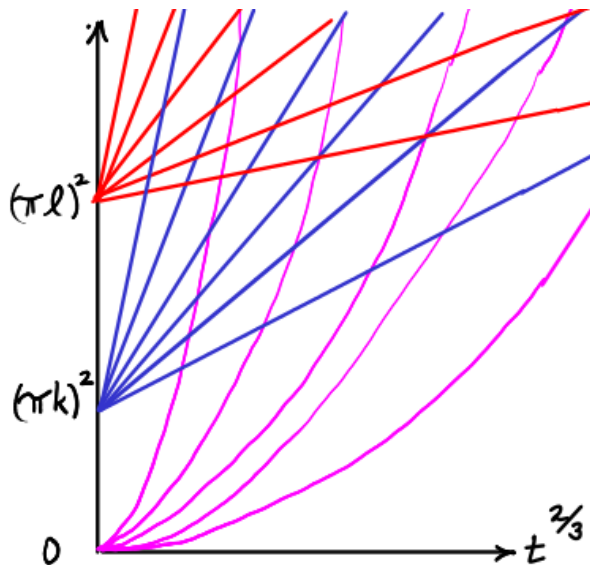
▶ $a_t(u) = \sum a_t^{\ell}(u^{\ell})$ where

$$a_t^{\ell}(v) = \int_1^{\infty} (t^2 \cdot (v')^2 + (\pi\ell)^2 \cdot v^2) dy$$

▶ spectrum of $a_t = \bigcup_{\ell}$ spectra of a_t^{ℓ}

▶ $\ell > 0 \Rightarrow$ eigenvalue $\lambda_t = (\pi\ell)^2 + c \cdot t^{\frac{2}{3}} + O(t)$ (Airy)

Eigenvalue branches of a_t



Variation and non-concentration

- ▶ \dot{a} is 'diagonal' with respect to a_t , and so $\dot{a}_t = \sum_{\ell} \dot{a}_t^{\ell}$ where

$$\dot{a}_t^{\ell}(v) = \frac{2}{t} \left(a_t^{\ell}(v) - E_t \cdot \|v\|^2 \right) + \frac{2}{t} \int (E_t \cdot y^{-2} - (\pi\ell)^2) \cdot v^2$$

Proposition (Nonconcentration at the turning point)

If w is an a_t^{ℓ} quasimode of order t at energy $E \gg (\ell\pi)^2$, then there exists $\kappa > 0$ so that for all small t

$$\int \left(\frac{E}{y^2} - (\ell\pi)^2 \right) w^2 \geq \kappa \cdot \|w\|^2.$$

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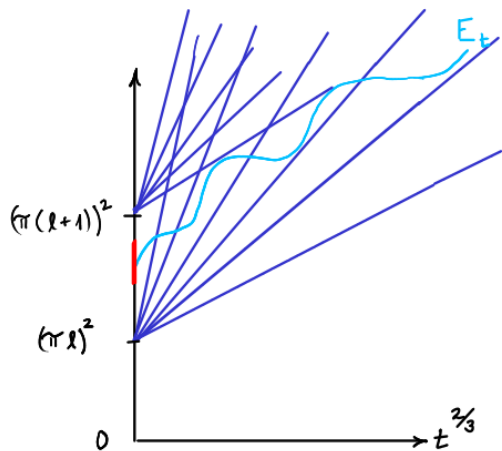
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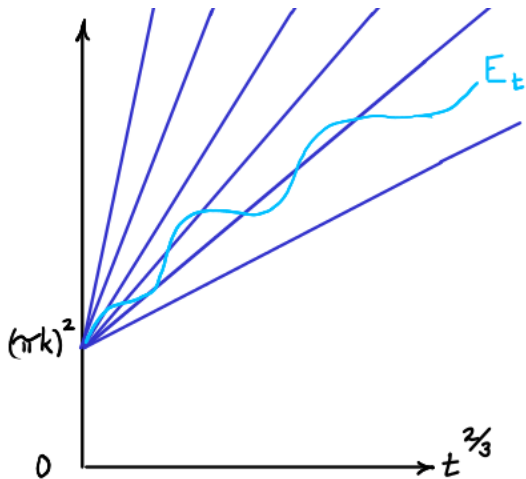
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$$E_t \rightarrow (k\pi)^2$$



► Integrality condition

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Relatively small k -mode implies large variation

w_t^ℓ = projection of w_t onto a_t eigenspaces with $\lambda_t \rightarrow (\pi\ell)^2$.

Lemma

Let $\rho < 1$. There exists $\delta > 0$ such that for sufficiently small t

$$\frac{\|w_t^k\|}{\|w_t\|} < \rho \implies \dot{E}_t \geq \frac{\delta}{t}.$$

Crossings imply relatively small k -mode

Lemma

Let $\rho < 1$. If for sufficiently small η and t

$$\text{dist}(E_t, \text{spec}(a_t^0)) < \eta \cdot t^{\frac{5}{3}}$$

then

$$\|w_t^k\| / \|w_t\| \leq \rho.$$

Idea of Proof:

$$(E_t - \lambda_t^0) \cdot \langle u_t, \psi_t \rangle = (q_t - a_t)(u_t, \psi_t^0) = t \cdot b_t(u_t, \psi_t^0) + O(t^2) \cdot \|u_t\| \cdot \|\psi_t^0\|.$$

ODE (Airy) approximation \Rightarrow there exists $\delta > 0$ so that

$$|b_t(u_t, \psi_t^0)| \geq \delta \cdot \left(\|w_t^k\| \cdot t^{\frac{2}{3}} - O(t^\delta) \|w_t\| \right) \cdot \|\psi_t^0\|$$

Here

$$b_t(u) = 2t \int_{y \leq b} p(y) \cdot u_x \cdot u_y \, dx \, dy$$

with p is smooth and $p(1) = 1$

Zeroth order approximation of proof of main theorem

- ▶ ($\rightarrow \leftarrow$) Suppose there exists u_t with $u_t \equiv 0$ and $E_t \not\rightarrow 0$.
- ▶ $E_0 > 0 \implies$ infinitely many eigenvalue branches of a_t^0 cross E_t .
- ▶ previous two lemmas imply that at each crossing $\dot{E}_t > \frac{\delta}{t}$
- ▶ Summing over crossings leads to $E_t \rightarrow 0$.

Additional considerations for summing

- ▶ Frequency of crossings: $\exists t_n \rightarrow 0$ with $E_{t_n} \in \text{spec}(a_{t_n})$ and

$$\lim_{n \rightarrow \infty} n \cdot t_n = k \cdot \ln(b).$$

- ▶ Width of crossings: $\lambda_t \in \text{spec}(a_t^0) \Rightarrow$

$$|s - t| < t^{\frac{8}{3}} \implies \text{dist}(\lambda_s, \text{spec}(a_s^0))$$

- ▶ 'Tracking': There exists unique eigenvalue branch λ_t^* of a_t^k so that $|E_t - \lambda_t^*| < Ct$ for t small.

- ▶ In truth, we use crossings to show that

$$\int_0^t (\dot{E}_s - \dot{\lambda}_s^*) ds > O(t^{\frac{2}{3}}).$$

a contradiction.