

On the spectral analysis of the Schrodinger operator with a periodic PT-symmetric potential

Oktay Veliev

Dogus University, Istanbul

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I am going to give a talk about the one dimensional Schrödinger operator  $L(q)$  generated in  $L_2(-\infty, \infty)$  by the differential expression

$$-y''(x) + q(x)y(x), \quad (1)$$

where  $q$  is 1-periodic  $q(x+1) = q(x)$ , complex-valued and PT symmetric  $\overline{q(-x)} = q(x)$  potential. It is well-known [ Rofe-Beketov (1963), McGarvey (1965)] that the spectrum  $\sigma(L(q))$  of the operator  $L(q)$  is the union of the spectra  $\sigma(L_t(q))$  of the operators  $L_t(q)$  for  $t \in (-\pi, \pi]$  generated in  $L_2[0, 1]$  by (1) and the boundary conditions

$$y(1) = e^{it}y(0), \quad y'(1) = e^{it}y'(0). \quad (2)$$

The spectrum of  $L_t$  consists of the eigenvalues  $\lambda_1(t), \lambda_2(t), \dots$  such that  $\lambda_n(-t) = \lambda_n(t)$  and  $\lambda_n(t) \rightarrow \infty$  as  $n \rightarrow \infty$ . Moreover they can be numbered so that  $\lambda_n$  is a continuous function on  $[0, \pi]$ . Thus

$$\Gamma_n = \{\lambda_n(t) : t \in [0, \pi]\}$$

is a continuous curve in the complex plane and the spectrum of  $L(q)$  is the union of the curves  $\Gamma_n$  for  $n = 1, 2, \dots$  :

$$\sigma(L(q)) = \bigcup_{t \in [0, \pi]} \sigma(L_t(q)) = \bigcup_{n \in \mathbb{N}} \Gamma_n.$$

In the first papers [Bender et al (1999)] about the PT-symmetric periodic potential, the disappearance of real energy bands for some complex-valued PT-symmetric periodic potentials have been reported. Shin (2004) showed that the disappearance of such real energy bands implies the existence of nonreal band spectra. He involved some condition on the Hill discriminant to show the existence of nonreal curves in the spectrum. Caliceti and Graffi (2009) found explicit condition on the Fourier coefficient of the potential providing the nonreal spectra for small potentials.

I prove that the main part of the spectrum of  $L(q)$  is real and contains the large part of  $[0, \infty)$ . However, in general, the spectrum contains also infinitely many nonreal arcs. The necessary and sufficient condition on the potential for finiteness of the number of the nonreal arcs is determined. Moreover, I find necessary and sufficient conditions for the equality of the spectrum of  $L(q)$  to the half line  $[c, \infty)$ . Then, I consider the connections between spectrality of  $L(q)$  and the reality of its spectrum. Finally, I consider in detail the optical potential

$$4 \cos^2 x + 4iV \sin 2x \quad (3)$$

The steps of my investigations are the followings:

- 1. General Properties of the spectrum of  $L(q)$  with PT-symmetric potential  $q$ .**
- 2. General necessary and sufficient conditions for finiteness of the number of the nonreal arcs in the spectrum  $\sigma(L(q))$  and for  $\sigma(L(q)) = [c, \infty)$**
- 3. Necessary and sufficient condition on the potential.**
- 4. The connections between reality of  $\sigma(L(q))$  and spectrality of  $L(q)$ .**
- 5. Detail investigations of the optical potential**

## General Properties of the spectrum of $L(q)$ .

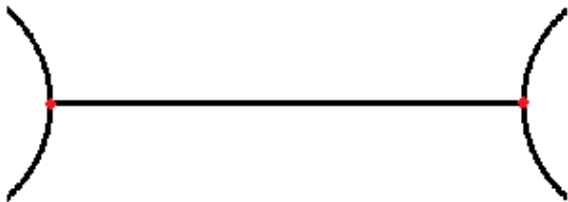
### Theorem

- (a) If  $\lambda_n(t_1)$  and  $\lambda_n(t_2)$  are real numbers, where  $0 \leq t_1 < t_2 \leq \pi$  then  $\gamma := \{\lambda_n(t) : t \in [t_1, t_2]\}$  is an interval of the real line.
- (b) If the eigenvalues  $\lambda_n(0)$  and  $\lambda_n(\pi)$  are real numbers then  $\lambda_n(t)$  are real eigenvalues of  $L_t(q)$  for all  $t \in (0, \pi)$ , that is,  $\Gamma_n$  is either  $[\lambda_n(0), \lambda_n(\pi)]$  or  $[\lambda_n(\pi), \lambda_n(0)]$ .
- (c) The spectrum of  $L(q)$  is completely real if and only if all eigenvalues of  $L_0(q)$  and  $L_\pi(q)$  are real.

Now to consider, the reality of the spectrum in detail, we investigate the points in which the spectrum ceases to be real. These points are crucial and can be defined as follows.

## Definition

A real number  $\lambda \in \sigma(L)$  is said to be a left (right) complexation point of the spectrum if there exists  $\varepsilon > 0$  such that  $[\lambda, \lambda + \varepsilon] \subset \sigma(L)$  ( $[\lambda - \varepsilon, \lambda] \subset \sigma(L)$ ) and  $\sigma(L)$  contains a nonreal number in any neighborhood of  $\lambda$ . Both left and right complexation points are called complexation points.



Pic. 1. Complexation points

## Theorem

*If  $\lambda_n(t)$  is a complexation point, then it is a multiple eigenvalue of  $L_t(q)$  and spectral singularities of  $L(q)$ . Moreover, the multiplicity of  $\lambda_n(t)$  is greater than 2 if  $t = 0, \pi$ . In any case if  $\lambda_n(t)$  is a complexation point then the eigenspace corresponding to  $\lambda_n(t)$  contains an associated function.*

The spectral singularities was defined as follows.

## Definition

We say that  $\lambda \in \sigma(L(q))$  is a spectral singularity of  $L(q)$  if for all  $\varepsilon > 0$  there exists a sequence  $\{\gamma_n\}$  of the arcs  $\gamma_n \subset \{z \in \mathbb{C} : |z - \lambda| < \varepsilon\}$  such that  $\gamma_n \subset \sigma(L)$  and

$$\lim_{n \rightarrow \infty} \|P(\gamma_n)\| = \infty. \quad (4)$$



## Theorem

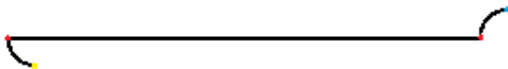
Suppose that  $n$  is a large number.

(a) There may exist at most one number  $\varepsilon_n$  and  $\pi - \delta_n$  in the neighborhoods of 0 and  $\pi$  respectively, such that  $\lambda_n(\varepsilon_n)$  and  $\lambda_n(\pi - \delta_n)$  are double eigenvalues.

(b) The eigenvalues  $\lambda_n(t)$  for  $t \in [0, \pi] \setminus \{\varepsilon_n, \pi - \delta_n\}$  are simple and are not the complexation points. Moreover  $\varepsilon_n \rightarrow 0$  and  $\delta_n \rightarrow 0$  as  $n \rightarrow \infty$ .

(c) The double eigenvalues  $\lambda_n(\varepsilon_n)$  and  $\lambda_n(\pi - \delta_n)$  are the complexation points if and only if  $\varepsilon_n \neq 0$  and  $\delta_n \neq 0$  respectively.

(d) The main part  $\{\lambda_n(t) : t \in [\varepsilon_n, \pi - \delta_n]\}$  of  $\Gamma_n$  is real. The other parts  $\{\lambda_n(t) : t \in [0, \varepsilon_n)\}$  and  $\{\lambda_n(t) : t \in (\pi - \delta_n, \pi]\}$  (if exist) are the pure nonreal parts of  $\Gamma_n$  and are called the tails of  $\Gamma_n$ .



Pic. 2,  $\Gamma_n$  for large  $n$ .

## General necessary and sufficient condition

### Theorem

*If  $n$  is a large number, then the component  $\Gamma_n$  has left (right) nonreal tail, that is, it contains a left (right) complexation points point if and only if  $\lambda_n(0)$  ( $\lambda_n(\pi)$ ) is a nonreal number.*

### Theorem

*If  $n$  is a large number, then the followings are equivalent*

- (a) Component  $\Gamma_n$  is the real interval.*
- (b) The eigenvalues  $\lambda_n(0)$  and  $\lambda_n(\pi)$  are the real numbers.*
- (c)  $\lambda_n(t)$  for  $t \in (0, \pi)$  are simple eigenvalues of  $L_t(q)$ .*
- (d) The eigenvalues  $\lambda_n(t)$  for  $t \in (0, \pi)$  are not the complexation point.*

### Theorem

*The spectrum of  $L(q)$  is a half line if and only if the followings hold: (i) one eigenvalue of  $L_0(q)$  is simple and the all others are double, (ii) all eigenvalues of  $L_\pi(q)$  are double.*

## Necessary and sufficient condition on the potential for the reality of $\Gamma_n$ for large $n$ .

Let  $S_p$  be the set of 1 periodic PT-symmetric functions  $q \in W_1^p[0, 1]$  such that

$$q(1) = q(0), q'(1) = q'(0), \dots, q^{(s-1)}(1) = q^{(s-1)}(0) \quad (5)$$

for some  $s \leq p$  and there exist positive constants  $c_1, c_2, c_3$  and  $N$  satisfying

$$|q_n| > c_1 n^{-s-1} \quad \& \quad c_2 |q_n| \leq |q_{-n}| \leq c_3 |q_n|, \quad \forall n > N,$$

where  $q_n = \int_0^1 q(x) e^{-i2\pi x} dx$ . In particular,  $S_0$  is the set of 1 periodic PT-symmetric functions  $q \in L_1[0, 1]$  satisfying

$$|q_n| > c_1 n^{-1} \quad \& \quad c_2 |q_n| \leq |q_{-n}| \leq c_3 |q_n|, \quad \forall n > N.$$

### Theorem

Suppose that  $q \in S_p$  and  $n$  is a large number. Then  $\Gamma_n \subset \mathbb{R}$  if and only if

$$q_n q_{-n} > 0. \quad (6)$$

Let  $Q_n$  and  $S_n$  be the Fourier coefficients of the function  $Q$  and  $S$  defined by

$$Q(x) = \int_0^x q(t) dt, \quad S(x) = Q^2(x).$$

Suppose that equality (5) and the inequality

$$|P_n| > cn^{-2s-2}, \quad (7)$$

holds, where  $P_n = q_n q_{-n} - q_n (S_{-n} - 2Q_0 Q_{-n}) - q_{-n} (S_n - 2Q_0 Q_n)$ .

### Theorem

*If  $n$  is a large number and  $P_n < 0$  then  $\Gamma_n$  has the nonreal tails.*

## The connections between reality of $\sigma(L(q))$ and spectrality of $L(q)$ .

### Theorem

(a) Let  $n$  be a large number and  $t \in (0, \pi)$ . Then the followings are equivalent.

- 1)  $\lambda_n(t)$  is a spectral singularity of  $L(q)$ .
- 2)  $\lambda_n(t)$  is a complexation points point of  $\sigma(L(q))$ .
- 3)  $\lambda_n(t)$  is a multiple eigenvalue of  $L_t(q)$ .
- 4) If  $t \in (0, h]$ , where  $h$  is a small number, then  $\{\lambda_n(s) : s \in [0, t)\}$ ; if  $t \in [\pi - h, \pi)$  then  $\{\lambda_n(s) : s \in (t, \pi]\}$  is a nonreal tail of  $\Gamma_n$ .

(b) Let  $q \in S_p$  for some  $p = 0, 1, \dots$ . Then the followings are equivalent.

- 1)  $L(q)$  is an asymptotically spectral operator.
- 2) There exists a large number  $m$  such that (6) holds for all  $|n| > m$ .
- 3) There exist a large number  $m$  such that  $\Gamma_n$  for  $n > m$  are real pairwise disjoint intervals.

## Detail investigations of the optical potential

$$4 \cos^2 x + 4iV \sin 2x = 2 + (1 + 2V)e^{i2x} + (1 - 2V)e^{-i2x}, \quad V \geq 0. \quad (3)$$

For the first time, the explanation of the nonreality of the spectrum of  $L(q)$ , redenoted by  $L(V)$ , for  $V > 0.5$  was done by Makris et al (2008). For  $V = 0,85$  they sketch the real and imaginary parts of the first two bands by using the numerical methods. Midya et al (2010) reduce the operator  $L(q)$  with potential (3) to the Mathier operator and using the tabular values establish that there is second critical point  $V_2 \sim 0.888437$  after which no part of the first and second bands remains real. Note that the corresponding Mathier operator in the case  $0 \leq V < 1/2$  is self-adjoint and hence the spectrum consists of the real intervals. The case  $V = 1/2$  for the first time is considered by Gasimov (1980) and it was proved that the spectrum is  $[0, \infty)$ .

Thus the shapes of the spectrum of  $L(V)$  for  $0 \leq V \leq 1/2$  are well-known



Pic. 3.  $0 \leq V < 1/2$ . As in the self-adjoint Mathier operator



Pic.4.  $V = 1/2$ . All periodic (except first) and antiperiodic eigenvalues are double and spectral singularities of  $L(q)$ .

Note that the reality of the spectrum for  $0 \leq V \leq 1/2$  also follows from the following. I proved that if  $ab = cd$ , then

$$\sigma(L(ae^{-i2x} + be^{i2x})) = \sigma(L(ce^{-i2x} + de^{i2x}))$$

Therefore

$$\sigma(L((1 + 2V)e^{i2x} + (1 - 2V)e^{-i2x})) = \sigma(L(2c \cos 2x)), \quad (8)$$

where  $c = \sqrt{1 - 4V^2}$

### Theorem

*The PT-symmetric operator  $L(ae^{-i2x} + be^{i2x})$  is a spectral operator if and only if  $a = b$ , that is,  $q(x)$  is the real potential  $2a \cos 2x$ . The optical operator with potential (3) is spectral if and only if  $V = 0$ .*

I give a complete description, provided with a mathematical proof, of the shape of the spectrum of the Hill operator  $L(V)$  with potential (3), when  $V$  changes from  $1/2$  to  $\sqrt{5}/2$ , that is,  $c = ir$  and  $r \in (0, 2)$ .



For this first we prove the followings

### Theorem

*If  $1/2 < V < \sqrt{5}/2$ , then all antiperiodic eigenvalues are nonreal and simple.*

### Theorem

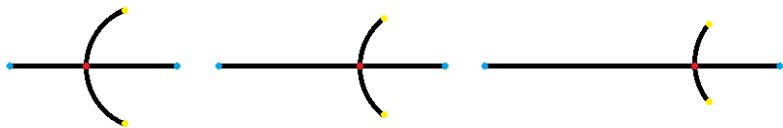
*If  $1/2 < V < \sqrt{5}/2$ , then all periodic eigenvalues (except first and second) are real and simple.*

### Theorem

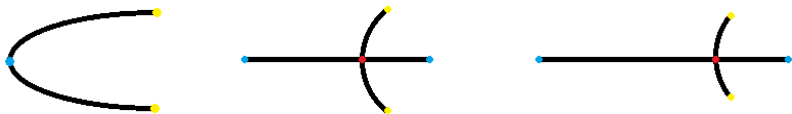
*There exists unique number  $V_2$  from  $(1/2, \sqrt{5}/2)$  such that*

- (a) If  $1/2 < V < V_2$  then the first and second eigenvalues are real and simple.*
- (b) If  $V = V_2$ , then  $\lambda_1(0) = \lambda_2(0) \in \mathbb{R}$*
- (c) If  $V_2 < V < \sqrt{5}/2$  then the first and second periodic eigenvalues  $\lambda_1(0)$  and  $\lambda_2(0)$  are simple and nonreal and  $\lambda_2(0) = \overline{\lambda_1(0)}$ .*

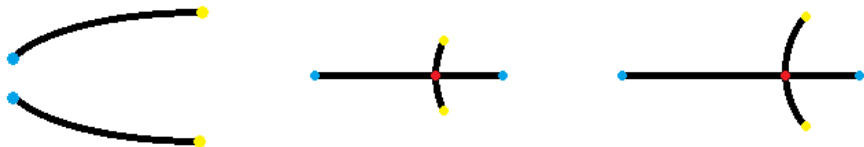
I prove that  $V_2$  is the second critical point (after which no part of the first and second bands remains real) and is a number between 0.8884370025 and 0.8884370117. Moreover, the last theorem shows that  $V_2$  is the unique degeneration point for the first periodic eigenvalue, in the sense that the first periodic eigenvalue of the potential (3) is simple for all  $V \in [0, \sqrt{5}/2) \setminus \{V_2\}$  and is double if  $V = V_2$ . My approach give the possibility to find the arbitrary close values of the  $k$ -th critical point  $V_k$  and prove that no part of the  $(2k - 3)$ -th and  $(2k - 2)$ -th bands remains real for  $V_k < V < V_k + \varepsilon$  for some positive  $\varepsilon$ , where  $k = 2, 3, \dots$  Using the last theorems we prove that the spectrum of  $L(V)$  has the following shape



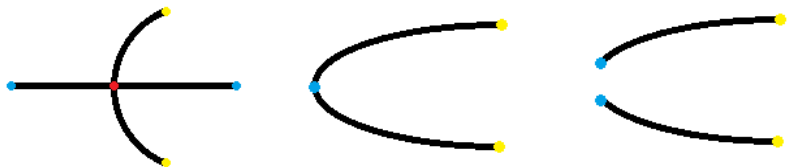
Pic. 5.  $1/2 < V < V_2$ . All antiperiodic eigenvalues are nonreal and periodic eigenvalues are real



Pic. 6.  $V = V_2$ .



Pic. 7.  $V_2 < V < \sqrt{5}/2$ . All periodic eigenvalues (except first and second) are real, the first and second periodic eigenvalues  $\lambda_1(0)$  and  $\lambda_2(0)$  are simple and nonreal and  $\lambda_2(0) = \overline{\lambda_1(0)}$ . Thus the first and second bands have different shapes in the following 3 cases:



Pic. 8. Case 1:  $1/2 < V < V_2$ , Case 2:  $V = V_2$ , Case 3:  $V_2 < V < \sqrt{5}/2$ .

In Case 1, the real part of the first component  $\Omega_1 = \Gamma_1 \cup \Gamma_2$  is the closed interval  $I_1 =: [\lambda_1(0), \lambda_2(0)]$ . We prove that if  $V$  approaches  $V_2$  from the left, then the eigenvalues  $\lambda_1(0)$  and  $\lambda_2(0)$  get close to each other and the length of the interval  $I_1$  approaches zero.

As a result if  $V = V_2$ , then we get the equality  $\lambda_1(0) = \lambda_2(0)$  which means that the first and second bands  $\Gamma_1$  and  $\Gamma_2$  have only one real point which is their common point

$$I_1 = \text{Re } \Omega_1 = \text{Re } \Gamma_1 = \text{Re } \Gamma_2 = \{\lambda_1(0)\}.$$

The other parts of the bands  $\Gamma_1$  and  $\Gamma_2$  are nonreal and symmetric with respect to the real line. Then we prove that if  $V > V_2$ , then the eigenvalues  $\lambda_1(0)$  and  $\lambda_2(0)$  get off the real line and hence  $I_1$  becomes the empty set. As a results, the first and second bands  $\Gamma_1$  and  $\Gamma_2$  became the nonreal curves symmetric with respect to the real line.

## Definition

A real number  $V_2 \in [0, \sqrt{5}/2)$  is called the second critical point or the degeneration point for the first periodic eigenvalue if the first real eigenvalue of  $L_0(V_2)$  is a double eigenvalue or equivalently if the first real component of  $\sigma(L(V_2))$  is a point.

Now let us describe briefly the shapes of all bands. I prove that the spectrum of  $L(V)$  in Case 1 has the following properties:

**Pr. 1.** *The real part  $\sigma(H(a)) \cap \mathbb{R}$  of the spectrum of  $H(a)$  consist of the intervals*

$$I_1 = [\lambda_1(0), \lambda_2(0)], I_2 = [\lambda_3(0), \lambda_4(0)], \dots, I_n = [\lambda_{2n-1}(0), \lambda_{2n}(0)], \dots \quad (9)$$

**Pr. 2.** *For each  $n = 1, 2, \dots$ , the interval  $I_n$  is the real part of  $\Omega_n =: \Gamma_{2n-1} \cup \Gamma_{2n}$ .*

**Pr. 3.** *The bands  $\Gamma_{2n-1}$  and  $\Gamma_{2n}$  have only one common point  $\Lambda_n$  which is interior point of  $I_n$ . Moreover,  $\Lambda_n$  is a double eigenvalue of  $L_{t_n}(V)$  for some  $t_n \in (0, \pi)$  and a spectral singularity of  $L(V)$  and hence*

$$\Gamma_{2n-1} \cap \Gamma_{2n} = \Lambda_n = \lambda_{2n-1}(t_n) = \lambda_{2n}(t_n) \in \mathbb{R}.$$

**Pr. 4.** *The real parts of the bands  $\Gamma_{2n-1}$  and  $\Gamma_{2n}$  are respectively the intervals  $[\lambda_{2n-1}(0), \Lambda_n] = \{\lambda_{2n-1}(t) : t \in [0, t_n]\}$  and  $[\Lambda_n, \lambda_{2n}(0)] = \{\lambda_{2n}(t) : t \in [0, t_n]\}$*

**Pr. 5.** *The nonreal parts of  $\Gamma_{2n-1}$  and  $\Gamma_{2n}$  are respectively the analytic curves*

$$\gamma_{2n-1}(a) = \{\lambda_{2n-1}(t) : t \in (t_n, \pi]\} , \quad \gamma_{2n}(a) = \{\lambda_{2n}(t) : t \in (t_n, \pi]\}$$

and  $\gamma_{2n}(a) = \{\bar{\lambda} : \lambda \in \gamma_{2n-1}(a)\}$

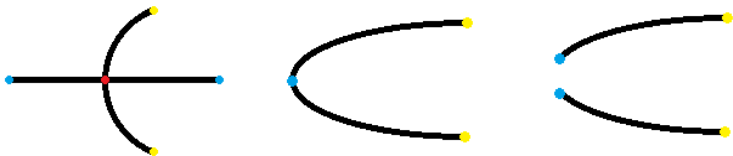
Thus the bands  $\Gamma_{2n-1}$  and  $\Gamma_{2n}$  are joined by  $\Lambda_n$  and hence they form together the connected subset of the spectrum. The spectrum  $\sigma(L(a))$  consist of the connected sets  $\Omega_1 =: \Gamma_1 \cup \Gamma_2$ ,  $\Omega_2 =: \Gamma_3 \cup \Gamma_4, \dots$  Moreover in Case 1 ( $1/2 < V < V_2$ ) we prove that

**Pr. 6** *The sets  $\Omega_1, \Omega_2, \dots$  are connected separated subset of the spectrum.*

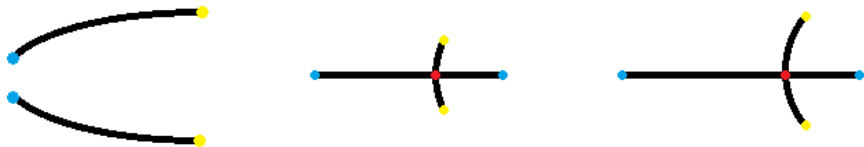
Moreover, in all cases the intervals (9) are pairwise disjoint sets. Therefore they are called the real components of  $\sigma(L(V))$  if  $V \in (1/2, \sqrt{5}/2)$ .



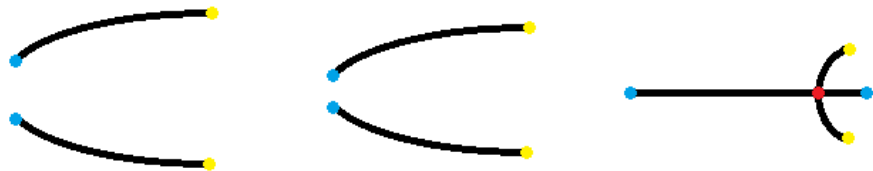
Note that in Case 2 ( $V = V_2$ ) and Case 3 ( $V_2 < V < \sqrt{5}/2$ ) the shapes of the components  $\Omega_2, \Omega_3, \dots$  are as in Case 1. In this way one can prove that there exists  $k$ -th critical point, denoted by  $V_k$ , such that for  $\frac{1}{2} < V < V_k$ ,  $V = V_k$  and  $V_k < V < V_k + \varepsilon$  the set  $\Omega_{k-1} =: \Gamma_{2k-3} \cup \Gamma_{2k-2}$  have the shape as  $\Omega_1$  in Case 1, Case 2 and Case 3 respectively.



Pic.9 The shape of  $\Omega_{k-1}$  in the cases  $\frac{1}{2} < V < V_k$ ,  $V = V_k$  and  $V_k < V < V_{k+1}$



Pic.10. The shape of the spectrum for  $V_2 < V < V_3$



Pic.11. The shape of the spectrum for  $V_3 < V < V_4$

# THANK YOU