### DIMERS AND RELATED MODELS IN

#### STATISTICAL MECHANICS

Béatrice de Tilière Université Paris-Est Créteil

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## STATISTICAL MECHANICS

Study the macroscopic properties of a physics system whose interactions are described on the microscopic level

The general framework is the following.

The structure of the physics system is represented by a finite graph G = (V, E).



## STATISTICAL MECHANICS

• Set of configurations on the graph G:  $\mathcal{C}(G)$ ,

- vertex configurations,
- edge configurations,
- vertex/edge configurations.
- Parameters representing:
  - ▶ the intensity of interactions between microscopic components,

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- the external temperature.
- $\Rightarrow$  Positive weight function *w* on edges/vertices.

## STATISTICAL MECHANICS

- To a configuration C, one assigns an energy  $\mathcal{E}_w(C)$ .
- Boltzmann probability on configurations:

$$\forall \mathbf{C} \in \mathcal{C}(\mathbf{G}), \quad \mathbb{P}(\mathbf{C}) = \frac{e^{-\mathcal{E}_w(\mathbf{C})}}{Z(\mathbf{G}, w)},$$

where  $Z(\mathbf{G}, w) = \sum_{\mathbf{C} \in \mathcal{C}(\mathbf{G})} e^{-\mathcal{E}_w(\mathbf{C})}$  is the partition function.

Understand the model when the graph is large (infinite).

Adsorption of di-atomic molecules on the surface of a crystal





Sir Ralph H. Fowler (1889-1944) Congrès Solvay 1927.

George S. Rushbrooke (1915-1995)

- Graph G = (V, E).
- ► A dimer configuration or perfect matching: subset of edges such that every vertex is incident to exactly one edge.

 $\Rightarrow \mathcal{M}(G) = \text{set of dimer configurations.}$ 

► A dimer configuration



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- ▶ Positive weight function on the edges:  $v = (v_e)_{e \in E}$ .
- Energy of a configuration M:  $\mathcal{E}_{\nu}(M) = -\sum_{e \in M} \log \nu_e$ .
- Dimer Boltzmann measure:

$$\forall M \in \mathcal{M}(G), \quad \mathbb{P}_{dimer}(M) = \frac{\prod_{e \in M} \nu_e}{Z_{dimer}(G, \nu)}$$

• The highest the weight  $v_e$ , the more likely is the edge e.

#### Model of ferromagnetism / mixture of two materials





Wilhelm Lenz (1888-1957)

Ernst Ising (1900-1998)

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- Graph G = (V, E).
- A spin configuration  $\sigma$  assigns to every vertex x of the graph G a spin  $\sigma_x \in \{-1, 1\}$ .

 $\Rightarrow C(G) = \{-1, 1\}^{V} = \text{set of spin configurations.}$ 

► A spin configuration

• A spin configuration / two interpretations.

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Magnetic moments:

+1/ $\rightarrow$ , -1/ $\leftarrow$ 



• A spin configuration / two interpretations.



## THE ISING MODEL

- Positive weight function: coupling constants  $J = (J_e)_{e \in E}$ .
- Energy of a spin configuration:  $\mathcal{E}_{J}(\sigma) = -\sum_{e=xy\in E} J_{xy}\sigma_{x}\sigma_{y}$ .
- Ising Boltzmann measure:

$$\forall \, \sigma \in \{-1,1\}^{\mathsf{V}}, \quad \mathbb{P}_{\mathrm{Ising}}(\sigma) = \frac{e^{-\mathcal{E}_{\mathsf{J}}(\sigma)}}{Z_{\mathrm{Ising}}(\mathsf{G},\mathsf{J})}.$$

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- Two neighboring spins  $\sigma_x, \sigma_y$  tend to align.
- The higher the coupling  $J_{xy}$ , the strongest is this tendency.



- **Percolation**: flow of a liquid through a porous material.
- SPANNING TREES/FORESTS: related to electrical networks and random walks.

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- RANDOM CLUSTER MODEL: includes percolation, Ising, Potts, spanning trees.
- ► VERTEX MODELS (6-8-···): 6-vertex is a model for ice.
- O(n) loop models.
- ▶ ...

# Macroscopic behavior

 ▶ Rhombus tilings ↔ dimers on the honeycomb lattice (illustration by R. Kenyon)



# Macroscopic behavior

### ▶ Ising model on the square lattice (illustrations by R. Cerf)



J small



#### J critical



J large

## Macroscopic behavior

- Identification of the phase transition.
- ► Understanding of the sub/super critical regimes.
- Understanding the critical model (at the phase transition):
  - Universality and conformal invariance.
  - Conjectures: Cardy, Duplantier, Nienhuis ...
    Proofs: Lawler, Schramm, Werner, D. Chelkak, S. Smirnov ...

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### EXACTLY SOLVABLE MODELS

• One of the tools to study the macroscopic behavior is the partition function:

$$Z(\mathsf{G},w) = \sum_{\mathsf{C}\in\mathcal{C}(\mathsf{G})} e^{-\mathcal{E}_w(\mathsf{C})},$$

the normalizing constant of the Boltzmann measure

$$\forall \mathbf{C} \in \mathcal{C}(\mathbf{G}), \quad \mathbb{P}(\mathbf{C}) = \frac{e^{-\mathcal{E}_w(\mathbf{C})}}{Z(\mathbf{G}, w)}.$$

- The model is exactly solvable if there exists an exact, explicit formula, for the partition function
- ▶ The two models considered are exactly solvable in 2d:
  - Ising: Onsager (1944) Kaufman Kac & Ward (1952) Fisher (1966).
  - Dimers: Kasteleyn Temperley & Fisher (1961).

# Partition function for the dimer model Kasteleyn, Temperley & Fisher

- Dimer model on a finite, planar graph G with weights v.
- Computation of the partition function,

$$Z_{\text{dimer}}(\mathsf{G}, \nu) = \sum_{\mathsf{M} \in \mathcal{M}(\mathsf{G})} \prod_{\mathsf{e} \in \mathsf{E}} \nu_{\mathsf{e}}.$$

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### Superimposition of two dimer configurations

• Let  $M_1$ ,  $M_2$  be two dimer configurations of G and let  $M_1 \cup M_2$  be their superimposition.



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### SUPERIMPOSITION OF TWO DIMER CONFIGURATIONS

• Let  $M_1$ ,  $M_2$  be two dimer configurations of G and let  $M_1 \cup M_2$  be their superimposition.



•  $M_1 \cup M_2$  is a disjoint union of alternating cycles, where an alternating cycle has edges alternating between  $M_1$  and  $M_2$ . An alternating cycle of length 2 is a doubled edge.

► A graph is isoradial if it is planar and can be embedded in the plane in such a way that all inner faces are inscribable in a circle of radius 1 and all circumcenters are in the interior of the faces (Duffin-Mercat-Kenyon).



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• Take the center of the circumcircles.



▶ Join each center to the "neighboring" vertices of G.
 ⇒ Associated rhombus graph G<sup>◦</sup>.



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► To every edge e, one assigns the half-angle \(\bar{\theta}\)<sub>e</sub> of the corresponding rhombus.



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