Symmetry breaking for a problem in optimal insulation

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Shape optimization and Isoperimetric and Functional Inequalities

Joint work with...

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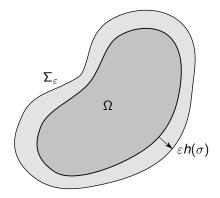
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The model

Conducting body Ω of conductivity 1.

Insulator $\Sigma_{\varepsilon} = \{\sigma + t\nu(\sigma) : \sigma \in \partial\Omega, \ 0 \le t < \varepsilon h(\sigma)\}$ of conductivity $\delta << 1$



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A problem of insulation (stationary case)

The temperature *u* of the conducting body Ω , with heat sources $f \in L^2(\Omega)$ is given by the solution of the PDE

$$\begin{cases} -\Delta u = f & \text{in } \Omega \\ -\Delta u = 0 & \text{in } \Sigma_{\varepsilon} \\ u = 0 & \text{on } \partial(\Omega \cup \Sigma_{\varepsilon}) \\ \frac{\partial u^{-}}{\partial \nu} = \delta \frac{\partial u^{+}}{\partial \nu} & \text{on } \partial \Omega. \end{cases}$$

or equivalently by the minimization of the functional

$$F_{\varepsilon\delta}(u) = \frac{1}{2} \int_{\Omega} |\nabla u|^2 \, dx + \frac{\delta}{2} \int_{\Sigma_{\varepsilon}} |\nabla u|^2 \, dx - \int_{\Omega} f u \, dx$$

in the Sobolev space $H^1_0(\Omega \cup \Sigma_{\varepsilon})$

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If $\varepsilon\approx\delta$ we get something different...

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One way to see what happens, is to look at the functional

$$F_{\varepsilon}(u) = \frac{1}{2} \int_{\Omega} |\nabla u|^2 \, dx + \frac{\varepsilon}{2} \int_{\Sigma_{\varepsilon}} |\nabla u|^2 \, dx - \int_{\Omega} f u \, dx \quad u \in H^1_0(\Omega \cup \Sigma_{\varepsilon})$$

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and notice that it **Γ**-converges to

$$E(u,h) = \frac{1}{2} \int_{\Omega} |\nabla u|^2 \, dx + \frac{1}{2} \int_{\partial \Omega} \frac{u^2}{h} \, d\mathcal{H}^{d-1} - \int_{\Omega} f u \, dx \quad u \in H^1(\Omega)$$

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Therefore this time we have the heat equation with Robin boundary conditions.

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We fix the total mass of insulator to be m > 0

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the case $f \equiv 1$ means that we are maximizing the total heat.

We want to find

$$\min_{h\in\mathcal{H}_m} \min_{u\in\mathcal{H}^1(\Omega)} \left\{ \frac{1}{2} \int_{\Omega} |\nabla u|^2 \, dx + \frac{1}{2} \int_{\partial\Omega} \frac{u^2}{h} \, d\mathcal{H}^{N-1} - \int_{\Omega} f u \, dx \right\}$$

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Therefore, the optimization problem can be rewritten as

$$\min\left\{\frac{1}{2}\int_{\Omega}|\nabla u|^{2}\,dx+\frac{1}{2m}\Big(\int_{\partial\Omega}|u|\,d\mathcal{H}^{d-1}\Big)^{2}-\int_{\Omega}fu\,dx\,:\,u\in H^{1}(\Omega)\right\}.$$

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Existence and uniqueness

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Image: A matrix

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Poincaré inequality enforces coercivity. There exists in fact a constant *C* such that

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Moreover

Proposition

Assume Ω is connected. Then the functional

$$u \mapsto F(u) = rac{1}{2} \int_{\Omega} |\nabla u|^2 \, dx + rac{1}{2m} \Big(\int_{\partial \Omega} |u| \, d\mathcal{H}^{d-1} \Big)^2$$

is strictly convex on $H^1(\Omega)$, hence for every $f \in L^2(\Omega)$ the minimization problem admits a unique solution.

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Let $\Omega = B_B$ and f = 1.

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Therefore the optimal insulation is given by a constant thickness:

$$h = rac{m}{d\omega_d R^{d-1}}$$

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There are two different cases

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• If $R_1 > R_2$ then $c_1 = 0$ and $c_2 = \frac{m}{d^2 \omega_d R_2^{d-2}}$. All the insulator must be uniformly distributed around the smaller ball, while the larger ball remains not insulated.

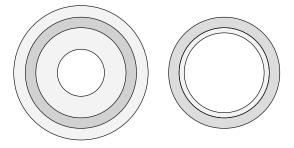
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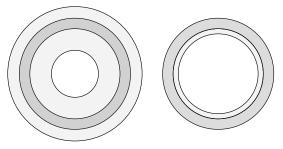
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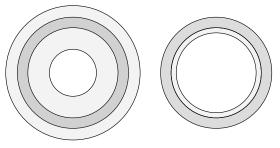


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In the non stationary case we consider the parabolic heat eqution

$$\begin{cases} u_t = \Delta u + f & \text{in } \Omega \\ u_t = \delta \Delta u & \text{in } \Sigma_{\varepsilon} \\ u = 0 & \text{on } \partial(\Omega \cup \Sigma_{\varepsilon}) \\ \frac{\partial u^-}{\partial \nu} = \delta \frac{\partial u^+}{\partial \nu} & \text{on } \partial \Omega. \end{cases}$$

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The first eigenvalue λ provides the heat loss rate.

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In the non stationary case we consider the parabolic heat eqution

$$\begin{cases} u_t = \Delta u + f & \text{in } \Omega \\ u_t = \delta \Delta u & \text{in } \Sigma_{\varepsilon} \\ u = 0 & \text{on } \partial(\Omega \cup \Sigma_{\varepsilon}) \\ \frac{\partial u^-}{\partial \nu} = \delta \frac{\partial u^+}{\partial \nu} & \text{on } \partial \Omega. \end{cases}$$

for which it is interesting to consider the eigenvalue problem

$$\begin{cases} -\Delta u = \lambda u & \text{in } \Omega \\ -\delta \Delta u = \lambda u & \text{in } \Sigma_{\varepsilon} \\ u = 0 & \text{on } \partial(\Omega \cup \Sigma_{\varepsilon}) \\ \frac{\partial u^{-}}{\partial \nu} = \delta \frac{\partial u^{+}}{\partial \nu} & \text{on } \partial \Omega. \end{cases}$$

The first eigenvalue λ provides the heat loss rate.

Again we consider the limit $\varepsilon, \delta \rightarrow 0$.

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If $\varepsilon = o(\delta)$ then we get Dirichlet Laplacian eigenvalue problem in Ω

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If $\delta = o(\varepsilon)$ then we get Neumann Laplacian eigenvalue problem in Ω

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$$\begin{cases} -\Delta u = \lambda u & \text{in } \Omega \\ \frac{\partial u}{\partial \nu} = 0 & \text{on } \partial \Omega \end{cases}$$

If $\varepsilon = \delta$ we get again Robin Laplacian eigenvalue problem

$$\begin{cases} -\Delta u = \lambda u & \text{in } \Omega \\ h \frac{\partial u}{\partial \nu} + u = 0 & \text{on } \partial \Omega \end{cases}$$

The first Robin eigenvalue $\lambda(h)$ is also characterized by

$$\lambda(h) = \inf \left\{ \frac{\int_{\Omega} |\nabla u|^2 \, dx + \int_{\partial \Omega} h^{-1} u^2 \, d\mathcal{H}^{d-1}}{\int_{\Omega} u^2 \, dx} : u \in H^1(\Omega), \ u \neq 0 \right\}.$$

Any *u* achieving such a minimum is an eigenfunction.

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The optimization consists in minimizing the heat loss rate $\min \{\lambda(h) : h \in \mathcal{H}_m\}.$

Therefore

$$\min_{h\in\mathcal{H}_m}\inf\left\{\frac{\int_{\Omega}|\nabla u|^2\,dx+\int_{\partial\Omega}h^{-1}u^2\,d\mathcal{H}^{d-1}}{\int_{\Omega}u^2\,dx}:\ u\in H^1(\Omega),\ u\neq 0\right\}.$$

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And again interchanging the minimization we get...

 $\min\left\{\lambda(h) : h \in \mathcal{H}_m\right\} =$

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$$\min\left\{\chi(H) : H \in \mathcal{H}_m\right\} =$$

$$\min\left\{\frac{\int_{\Omega} |\nabla u|^2 \, dx + \frac{1}{m} \left(\int_{\partial \Omega} |u| \, d\mathcal{H}^{d-1}\right)^2}{\int_{\Omega} u^2 \, dx} : u \in H^1(\Omega)\right\}$$

 $\min\left(\left(h\right), h \in \mathcal{U}\right)$

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$$\min\left\{\frac{\int_{\Omega}|\nabla u|^2\,dx+\frac{1}{m}\left(\int_{\partial\Omega}|u|\,d\mathcal{H}^{d-1}\right)^2}{\int_{\Omega}u^2\,dx}:\ u\in H^1(\Omega)\right\}.$$

 $\min \left\{ \lambda(h) : h \in \mathcal{H}_m \right\} =$

Let *u* be a minimizer,

$$\min\left\{\lambda(h) : h \in \mathcal{H}_m\right\} = \\ \min\left\{\frac{\int_{\Omega} |\nabla u|^2 \, dx + \frac{1}{m} \left(\int_{\partial \Omega} |u| \, d\mathcal{H}^{d-1}\right)^2}{\int_{\Omega} u^2 \, dx} : u \in H^1(\Omega)\right\}$$

Let u be a minimizer, again the optimal density h_{opt} is proportional to u

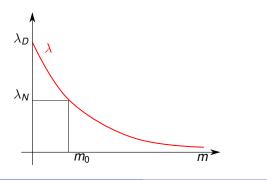
$$h_{opt} = m rac{u}{\int_{\partial\Omega} u \, d\mathcal{H}^{d-1}}$$
 .

Theorem

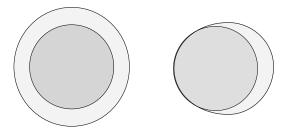
Let Ω be a ball. Then there exists $m_0 > 0$ such that the solution of the variational problem is radial if $m > m_0$, while the solution is not radial for $0 < m < m_0$. As a consequence, the optimal density h_{opt} is not constant if $m < m_0$.

Theorem

Let Ω be a ball. Then there exists $m_0 > 0$ such that the solution of the variational problem is radial if $m > m_0$, while the solution is not radial for $0 < m < m_0$. As a consequence, the optimal density h_{opt} is not constant if $m < m_0$.



The picture is therefore the following



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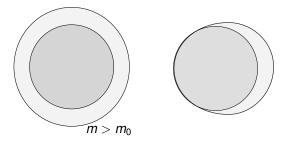


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The picture is therefore the following

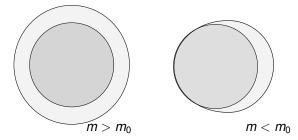


Image: A matrix

Thank you!

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