

Bounds for the spectrum of the magnetic Laplacian

Bruno Colbois (joint work with Alessandro Savo)

Luminy, November 23, 2016

Let (M, g) be a complete Riemannian manifold and let Ω be a compact domain of M with smooth boundary $\partial\Omega$, if non empty (if M is closed, we can choose $\Omega = M$ and $\partial\Omega = \emptyset$).

Given a smooth real 1-form A on Ω we define a connection ∇^A on $C^\infty(\Omega, \mathbb{C})$ acting on the space of smooth complex valued functions on Ω as follows:

Given a smooth real 1-form A on Ω we define a connection ∇^A on $C^\infty(\Omega, \mathbb{C})$ acting on the space of smooth complex valued functions on Ω as follows:

$$\nabla_X^A u = \nabla_X u - iA(X)u \quad (1)$$

for all vector fields X on Ω and for all $u \in C^\infty(\Omega, \mathbb{C})$.

The operator

$$\Delta_A = (\nabla^A)^* \nabla^A \quad (2)$$

is called the *magnetic Laplacian* associated to the magnetic potential A .

The operator

$$\Delta_A = (\nabla^A)^* \nabla^A \quad (2)$$

is called the *magnetic Laplacian* associated to the magnetic potential A .

$$\Delta_A u = \Delta u + |A|^2 u + 2i \langle du, A \rangle - iu \delta A.$$

In particular:

$$\Delta_A u = \Delta u + |A|^2 u - 2i \langle du, A \rangle \quad \text{whenever} \quad \delta A = 0.$$

The operator

$$\Delta_A = (\nabla^A)^* \nabla^A \quad (2)$$

is called the *magnetic Laplacian* associated to the magnetic potential A .

$$\Delta_A u = \Delta u + |A|^2 u + 2i \langle du, A \rangle - iu \delta A.$$

In particular:

$$\Delta_A u = \Delta u + |A|^2 u - 2i \langle du, A \rangle \quad \text{whenever} \quad \delta A = 0.$$

The 2-form

$$B = dA$$

being the *associated magnetic field*.

The operator

$$\Delta_A = (\nabla^A)^* \nabla^A \quad (2)$$

is called the *magnetic Laplacian* associated to the magnetic potential A .

$$\Delta_A u = \Delta u + |A|^2 u + 2i \langle du, A \rangle - iu \delta A.$$

In particular:

$$\Delta_A u = \Delta u + |A|^2 u - 2i \langle du, A \rangle \quad \text{whenever} \quad \delta A = 0.$$

The 2-form

$$B = dA$$

being the *associated magnetic field*.

If ϕ is a function on Ω , Δ_A and $\Delta_{A+d\phi}$ are unitarily equivalent and will have the same spectrum.

The operator

$$\Delta_A = (\nabla^A)^* \nabla^A \quad (2)$$

is called the *magnetic Laplacian* associated to the magnetic potential A .

$$\Delta_A u = \Delta u + |A|^2 u + 2i \langle du, A \rangle - iu \delta A.$$

In particular:

$$\Delta_A u = \Delta u + |A|^2 u - 2i \langle du, A \rangle \quad \text{whenever} \quad \delta A = 0.$$

The 2-form

$$B = dA$$

being the *associated magnetic field*.

If ϕ is a function on Ω , Δ_A and $\Delta_{A+d\phi}$ are unitarily equivalent and will have the same spectrum. If $A = 0$, we recover the usual Laplacian.

Example

Let $\Omega \subset \mathbb{R}^2$ a domain.

Example

Let $\Omega \subset \mathbb{R}^2$ a domain.

As potential, consider for $b \in \mathbb{R}$,

$$A((x_1, x_2)) = \frac{-bx_2}{2} dx^1 + \frac{bx_1}{2} dx^2.$$

Example

Let $\Omega \subset \mathbb{R}^2$ a domain.

As potential, consider for $b \in \mathbb{R}$,

$$A((x_1, x_2)) = \frac{-bx_2}{2} dx^1 + \frac{bx_1}{2} dx^2.$$

We have $B = b dx^1 \wedge dx^2$. The magnetic field is constant.

Example

Let $\Omega \subset \mathbb{R}^2$ be a domain with a hole, say around a point $a = (a_1, a_2)$. As potential, consider

$$A_{a,\gamma}(x_1, x_2) = \gamma \left(\frac{-(x_2 - a_2)}{(x_1 - a_1)^2 + (x_2 - a_2)^2} dx^1 + \frac{(x_1 - a_1)}{(x_1 - a_1)^2 + (x_2 - a_2)^2} dx^2 \right)$$

with $\gamma \in (0, 1)$.

Example

Let $\Omega \subset \mathbb{R}^2$ be a domain with a hole, say around a point $a = (a_1, a_2)$. As potential, consider

$$A_{a,\gamma}(x_1, x_2) = \gamma \left(\frac{-(x_2 - a_2)}{(x_1 - a_1)^2 + (x_2 - a_2)^2} dx^1 + \frac{(x_1 - a_1)}{(x_1 - a_1)^2 + (x_2 - a_2)^2} dx^2 \right)$$

with $\gamma \in (0, 1)$.

Then, $A_{a,\gamma}$ is a closed 1-form, and $dA_{a,\gamma} = B = 0$.

Example

Let $\Omega \subset \mathbb{R}^2$ be a domain with a hole, say around a point $a = (a_1, a_2)$. As potential, consider

$$A_{a,\gamma}(x_1, x_2) = \gamma \left(\frac{-(x_2 - a_2)}{(x_1 - a_1)^2 + (x_2 - a_2)^2} dx^1 + \frac{(x_1 - a_1)}{(x_1 - a_1)^2 + (x_2 - a_2)^2} dx^2 \right)$$

with $\gamma \in (0, 1)$.

Then, $A_{a,\gamma}$ is a closed 1-form, and $dA_{a,\gamma} = B = 0$.

The circulation of $A_{a,\gamma}$ is γ .

Example

Let $\Omega \subset \mathbb{R}^2$ be a domain with a hole, say around a point $a = (a_1, a_2)$. As potential, consider

$$A_{a,\gamma}(x_1, x_2) = \gamma \left(\frac{-(x_2 - a_2)}{(x_1 - a_1)^2 + (x_2 - a_2)^2} dx^1 + \frac{(x_1 - a_1)}{(x_1 - a_1)^2 + (x_2 - a_2)^2} dx^2 \right)$$

with $\gamma \in (0, 1)$.

Then, $A_{a,\gamma}$ is a closed 1-form, and $dA_{a,\gamma} = B = 0$.

The circulation of $A_{a,\gamma}$ is γ . For a simple closed curve c around a , the circulation of A is given by

$$\Phi_c^A = \frac{1}{2\pi} \int_c A.$$

If the boundary of Ω is non empty, we will consider Neumann magnetic conditions, that is:

$$\nabla_N^A u = 0 \quad \text{on} \quad \partial\Omega, \quad (3)$$

where N denotes the inner unit normal. Then Δ_A is self-adjoint, and admits a discrete spectrum

$$0 \leq \lambda_1(\Delta_A) \leq \lambda_2(\Delta_A) \leq \dots \rightarrow \infty.$$

Note that there exist a lot of result with *Dirichlet* boundary conditions. As example:

Note that there exist a lot of result with *Dirichlet* boundary conditions. As example:

For a domain $\Omega \subset \mathbb{R}^2$ with a constant magnetic field, there is a Faber-Krahn inequality. The disc minimize the first eigenvalue (L. Erdős, 1996).

Note that there exist a lot of result with *Dirichlet* boundary conditions. As example:

For a domain $\Omega \subset \mathbb{R}^2$ with a constant magnetic field, there is a Faber-Krahn inequality. The disc minimize the first eigenvalue (L. Erdős, 1996).

Sharp upper bounds on starlike plane domains for some functionals of the eigenvalues (Laugesen-Siudeja, 2015).

Note that there exist a lot of result with *Dirichlet* boundary conditions. As example:

For a domain $\Omega \subset \mathbb{R}^2$ with a constant magnetic field, there is a Faber-Krahn inequality. The disc minimize the first eigenvalue (L. Erdős, 1996).

Sharp upper bounds on starlike plane domains for some functionals of the eigenvalues (Laugesen-Siudeja, 2015).

For a domain $\Omega \subset \mathbb{R}^2$ with a hole around a point $a = (a_1, a_2)$ and circulation $\frac{1}{2}$, what does occur if a approach the boundary? (Noris, Terracini, Bonnaillie, Felli, ...)

With Neumann boundary conditions: Sharp upper bounds on starlike plane domains for some functionals of the eigenvalues (Laugesen-Siudeja, 2015).

With Neumann boundary conditions: Sharp upper bounds on starlike plane domains for some functionals of the eigenvalues (Laugesen-Siudeja, 2015).

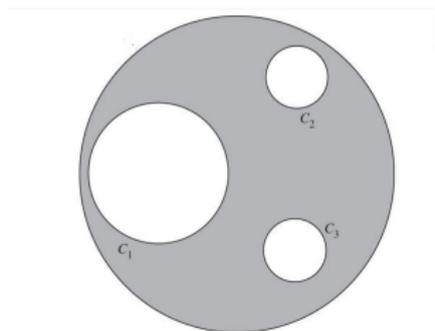
(B. Helffer, M. Hoffmann-Ostenhof, T. Hoffmann-Ostenhof, M. P. Owen 1990).

Let $\Omega \subset \mathbb{R}^2$ be a region with smooth boundary, which is homeomorphic to a disk with k holes. They look at a potential A with $dA = 0$.

With Neumann boundary conditions: Sharp upper bounds on starlike plane domains for some functionals of the eigenvalues (Laugesen-Siudeja, 2015).

(B. Helffer, M. Hoffmann-Ostenhof, T. Hoffmann-Ostenhof, M. P. Owen 1990).

Let $\Omega \subset \mathbb{R}^2$ be a region with smooth boundary, which is homeomorphic to a disk with k holes. They look at a potential A with $dA = 0$. Let c_i a closed path which parametrise the boundary of the i^{th} hole, and $\Phi_{c_i}^A = \frac{1}{2\pi} \int_{c_i} A$.



The first eigenvalue of the magnetic operator Δ_A depends only on the circulations $(\Phi_{c_1}^A, \dots, \Phi_{c_k}^A)$ of A .

The first eigenvalue of the magnetic operator Δ_A depends only on the circulations $(\Phi_{c_1}^A, \dots, \Phi_{c_k}^A)$ of A .

$\lambda_1(\Delta_A) = 0$ if and only if $\Phi_{c_i}^A \in \mathbb{Z}$ for all i . (See also Shikagawa for compact manifolds).

The first eigenvalue of the magnetic operator Δ_A depends only on the circulations $(\Phi_{c_1}^A, \dots, \Phi_{c_k}^A)$ of A .

$\lambda_1(\Delta_A) = 0$ if and only if $\Phi_{c_i}^A \in \mathbb{Z}$ for all i . (See also Shikagawa for compact manifolds).

Moreover, when $k = 1$, $\lambda_1(\Delta_A, \Omega)$ is maximal when $\Phi_c^A = \frac{1}{2}$.

There are two recent very interesting results regarding lower bound for λ_1 :

There are two recent very interesting results regarding lower bound for λ_1 :

A Cheeger type inequality (by Carsten Lange, Shiping Liu, Norbert Peyerimhoff, Olaf Post);

There are two recent very interesting results regarding lower bound for λ_1 :

A Cheeger type inequality (by Carsten Lange, Shiping Liu, Norbert Peyerimhoff, Olaf Post);

Lower bound thanks to Bochner methods (by Michela Egidi, Shiping Liu, Florentin Münch, Norbert Peyerimhoff).

Goal for today:

Goal for today:

- upper bounds for all the spectrum using geometric methods;

Goal for today:

- upper bounds for all the spectrum using geometric methods;
- lower bounds for λ_1 in a specific situation;

Goal for today:

- upper bounds for all the spectrum using geometric methods;
- lower bounds for λ_1 in a specific situation;
- highlighting the role of the circulation $\Phi_{C_i}^A = \frac{1}{2\pi} \int_{C_i} A$.

Remark

On the sphere S^2 (with its canonical metric), there exists a family A_k of potentials, such that $\lambda_1(\Delta_{A_k}) \rightarrow \infty$ as $k \rightarrow \infty$. If $B_k = dA_k$, we have $\|B_k\|_2 \rightarrow \infty$ as $k \rightarrow \infty$ (Besson-C-Courtois).

Remark

On the sphere S^2 (with its canonical metric), there exists a family A_k of potentials, such that $\lambda_1(\Delta_{A_k}) \rightarrow \infty$ as $k \rightarrow \infty$. If $B_k = dA_k$, we have $\|B_k\|_2 \rightarrow \infty$ as $k \rightarrow \infty$ (Besson-C-Courtois). So, we need to take account of $\|B\|_2$ if we want to find upper bounds.

A remark about the potential A .

A remark about the potential A .

Because of the Gauge invariance and the Hodge decomposition, we can suppose that A has the following expression

$$A = h + \delta\psi$$

A remark about the potential A .

Because of the Gauge invariance and the Hodge decomposition, we can suppose that A has the following expression

$$A = h + \delta\psi$$

Here, ψ is a 2-form and h is a 1-form satisfying $dh = \delta h = 0$.

Moreover, if $\partial\Omega \neq \emptyset$ and N is the normal derivative to the boundary, we can choose ψ and h tangential (i.e. $\psi(N, \cdot) = 0$ and $h(N) = 0$).

A remark about the potential A .

Because of the Gauge invariance and the Hodge decomposition, we can suppose that A has the following expression

$$A = h + \delta\psi$$

Here, ψ is a 2-form and h is a 1-form satisfying $dh = \delta h = 0$.

Moreover, if $\partial\Omega \neq \emptyset$ and N is the normal derivative to the boundary, we can choose ψ and h tangential (i.e. $\psi(N, \cdot) = 0$ and $h(N) = 0$).

We denote by $Har_1(\Omega)$ the 1-forms satisfying $dh = \delta h = 0$ and $h(N) = 0$ on $\partial\Omega$.

A remark about the potential A .

Because of the Gauge invariance and the Hodge decomposition, we can suppose that A has the following expression

$$A = h + \delta\psi$$

Here, ψ is a 2-form and h is a 1-form satisfying $dh = \delta h = 0$.

Moreover, if $\partial\Omega \neq \emptyset$ and N is the normal derivative to the boundary, we can choose ψ and h tangential (i.e. $\psi(N, \cdot) = 0$ and $h(N) = 0$).

We denote by $Har_1(\Omega)$ the 1-forms satisfying $dh = \delta h = 0$ and $h(N) = 0$ on $\partial\Omega$.

In particular, we have $B = d\delta\psi$, and $B = 0$ if and only if $\psi = 0$.

A remark about the potential A .

Because of the Gauge invariance and the Hodge decomposition, we can suppose that A has the following expression

$$A = h + \delta\psi$$

Here, ψ is a 2-form and h is a 1-form satisfying $dh = \delta h = 0$. Moreover, if $\partial\Omega \neq \emptyset$ and N is the normal derivative to the boundary, we can choose ψ and h tangential (i.e. $\psi(N, \cdot) = 0$ and $h(N) = 0$).

We denote by $Har_1(\Omega)$ the 1-forms satisfying $dh = \delta h = 0$ and $h(N) = 0$ on $\partial\Omega$.

In particular, we have $B = d\delta\psi$, and $B = 0$ if and only if $\psi = 0$.

The 1-forms $\delta\psi$ and h are L^2 -orthogonal on Ω :

$$\int_{\Omega} \langle \delta\psi, h \rangle dvol_g = 0.$$

Let Ω be a domain in (M, g) . We choose a family of closed curves (c_1, \dots, c_m) , basis of the homology of degree 1 of Ω and we consider the dual basis of harmonic 1-forms $A_1, \dots, A_m \in \text{Har}_1(\Omega)$: we have

$$\Phi_{c_i}^{A_j} = \frac{1}{2\pi} \int_{c_i} A_j = \delta_{ij}.$$

Let Ω be a domain in (M, g) . We choose a family of closed curves (c_1, \dots, c_m) , basis of the homology of degree 1 of Ω and we consider the dual basis of harmonic 1-forms $A_1, \dots, A_m \in \text{Har}_1(\Omega)$: we have

$$\Phi_{c_i}^{A_j} = \frac{1}{2\pi} \int_{c_i} A_j = \delta_{ij}.$$

Given $A \in \text{Har}_1(\Omega)$, we write

$$\Phi^A = (\Phi_{c_1}^A, \dots, \Phi_{c_m}^A)$$

Let Ω be a domain in (M, g) . We choose a family of closed curves (c_1, \dots, c_m) , basis of the homology of degree 1 of Ω and we consider the dual basis of harmonic 1-forms $A_1, \dots, A_m \in \text{Har}_1(\Omega)$: we have

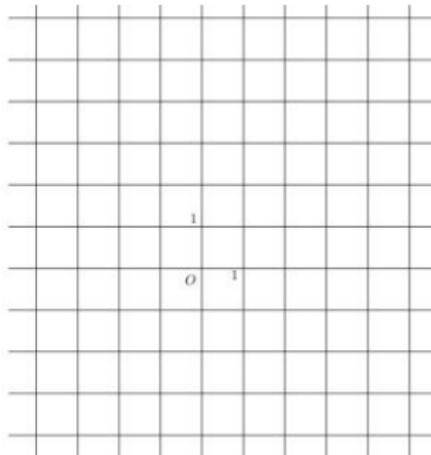
$$\Phi_{c_i}^{A_j} = \frac{1}{2\pi} \int_{c_i} A_j = \delta_{ij}.$$

Given $A \in \text{Har}_1(\Omega)$, we write

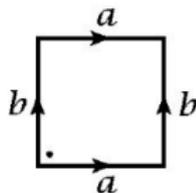
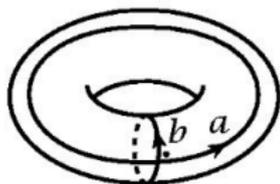
$$\Phi^A = (\Phi_{c_1}^A, \dots, \Phi_{c_m}^A)$$

and we denote by $d(\Phi^A, \mathbb{Z}^m)$ the Euclidean distance between Φ^A and the Euclidean lattice:

$$d(\Phi^A, \mathbb{Z}^m)^2 = \min \left\{ \sum_{j=1}^m (\Phi_{c_j}^A - k_j)^2 : (k_1, \dots, k_m) \in \mathbb{Z}^m \right\}.$$



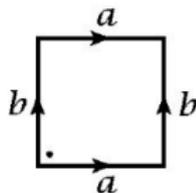
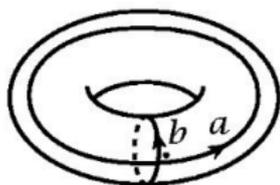
Example of a torus:



The curve c_1 and c_2 correspond to a and b .

In coordinates, if the length of c_i is α_i , $A_1 = \frac{2\pi}{\alpha_1} dx^1$, $A_2 = \frac{2\pi}{\alpha_2} dx^2$.

Example of a torus:



The curve c_1 and c_2 correspond to a and b .

In coordinates, if the length of c_i is α_i , $A_1 = \frac{2\pi}{\alpha_1} dx^1$, $A_2 = \frac{2\pi}{\alpha_2} dx^2$.

If $A = \beta_1 A_1 + \beta_2 A_2$,

$$d(\Phi^A, \mathbb{Z}^2)^2 = \min\{(\beta_1 - k_1)^2 + (\beta_2 - k_2)^2 : (k_1, k_2) \in \mathbb{Z}^2\}$$

We also introduce the lattice generated by the dual basis (A_1, \dots, A_m) :

$$\mathcal{L}_{\mathbf{Z}} = \{k_1 A_1 + \dots + k_m A_m : k_j \in \mathbf{Z}\}$$

which is an abelian subgroup of $Har_1(\Omega)$. Given $A \in Har_1(\Omega)$, we define its minimum distance to the lattice $\mathcal{L}_{\mathbf{Z}}$ by the formula:

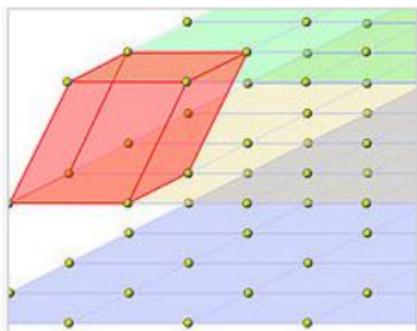
$$d(A, \mathcal{L}_{\mathbf{Z}})^2 = \min \left\{ \|\omega - A\|^2, \omega \in \mathcal{L}_{\mathbf{Z}} \right\}.$$

We also introduce the lattice generated by the dual basis (A_1, \dots, A_m) :

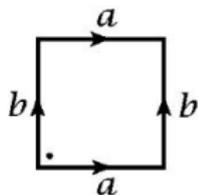
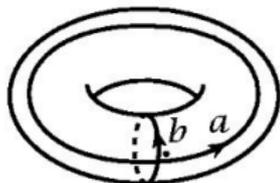
$$\mathcal{L}_{\mathbf{Z}} = \{k_1 A_1 + \dots + k_m A_m : k_j \in \mathbf{Z}\}$$

which is an abelian subgroup of $Har_1(\Omega)$. Given $A \in Har_1(\Omega)$, we define its minimum distance to the lattice $\mathcal{L}_{\mathbf{Z}}$ by the formula:

$$d(A, \mathcal{L}_{\mathbf{Z}})^2 = \min \left\{ \|\omega - A\|^2, \omega \in \mathcal{L}_{\mathbf{Z}} \right\}.$$



For the example of a torus:



If A is an harmonic form, $A = \beta_1 dx^1 + \beta_2 dx^2$,

$$d(A, \mathcal{L}_Z)^2 = \min \left\{ \left\| \left(k_1 \frac{2\pi}{\alpha_1} - \beta_1 \right)^2 + \left(k_2 \frac{2\pi}{\alpha_2} - \beta_2 \right)^2 \right\|^2, (k_1, k_2) \in \mathbb{Z}^2 \right\}$$

Our results: Ω domain of (M, g) .

Our results: Ω domain of (M, g) .

Recall that we write

$$A = \delta\psi + h$$

with $dh = \delta h = 0$ and $d(\delta\psi) = B$.

Upper bound for λ_1 :

Upper bound for λ_1 :

We have

$$\lambda_1(\Delta_A) \leq \frac{d(A, \mathcal{L}_Z)^2}{|\Omega|} + \frac{\|B\|_2^2}{\lambda_{1,1}(\Omega)|\Omega|},$$

where $\lambda_{1,1}$ denotes the first nonzero eigenvalue of the Laplacian on co-exact 1-forms and $|\Omega|$ denotes the volume of Ω .

Upper bound for λ_1 :

We have

$$\lambda_1(\Delta_A) \leq \frac{d(A, \mathcal{L}_Z)^2}{|\Omega|} + \frac{\|B\|_2^2}{\lambda_{1,1}(\Omega)|\Omega|},$$

where $\lambda_{1,1}$ denotes the first nonzero eigenvalue of the Laplacian on co-exact 1-forms and $|\Omega|$ denotes the volume of Ω .

In particular, if $B = 0$, we get

$$\lambda_1(\Delta_A) \leq \frac{d(A, \mathcal{L}_Z)^2}{|\Omega|},$$

and this inequality is sharp (equality in the case of a flat rectangular torus).

Some comments

Some comments

In the first inequality, we need to take account of B , but the presence of $\|B\|_2^2$ is probably not optimal. L. Erdős obtains an estimate with $\|B\|_1$ in the case of surfaces. The proof is much more difficult.

Some comments

In the first inequality, we need to take account of B , but the presence of $\|B\|_2^2$ is probably not optimal. L. Erdős obtains an estimate with $\|B\|_1$ in the case of surfaces. The proof is much more difficult.

The term $\lambda_{1,1}$ reflects the presence of the geometry. In his estimates for surfaces, L. Erdős has a term depending on the curvature and injectivity radius of the surface. However, in dimension 2, $\lambda_{1,1}$ is equal to the first eigenvalue of the Laplacian on functions, and this is no longer the case in higher dimensions.

The other eigenvalues: we get our estimates with respect to the lower bound of the Ricci curvature.

The other eigenvalues: we get our estimates with respect to the lower bound of the Ricci curvature.

There exist $c_1(n), c_2(n), c_3(n)$ depending only on the dimension n of M , such that for a domain $\Omega \subset (M, g)$ with $Ric(M, g) \geq -a^2(n-1)$, we have

The other eigenvalues: we get our estimates with respect to the lower bound of the Ricci curvature.

There exist $c_1(n), c_2(n), c_3(n)$ depending only on the dimension n of M , such that for a domain $\Omega \subset (M, g)$ with $Ric(M, g) \geq -a^2(n-1)$, we have

$$\lambda_k(\Delta_A, \Omega) \leq C_1(\Omega, A) + c_2(n)a^2 + c_3(n) \left(\frac{k}{|\Omega|} \right)^{2/n},$$

The other eigenvalues: we get our estimates with respect to the lower bound of the Ricci curvature.

There exist $c_1(n), c_2(n), c_3(n)$ depending only on the dimension n of M , such that for a domain $\Omega \subset (M, g)$ with $Ric(M, g) \geq -a^2(n-1)$, we have

$$\lambda_k(\Delta_A, \Omega) \leq C_1(\Omega, A) + c_2(n)a^2 + c_3(n) \left(\frac{k}{|\Omega|} \right)^{2/n},$$

with

$$C_1(\Omega, A) \leq \frac{c_1(n)}{|\Omega|} \left(d(A, \mathcal{L}_Z)^2 + \frac{\|B\|^2}{\lambda_{1,1}(\Omega)} \right).$$

1. If $A = 0$, we get

$$\lambda_k(\Delta_A, \Omega) \leq c_2(n)a^2 + c_3(n) \left(\frac{k}{|\Omega|} \right)^{2/n},$$

which is the same estimate as for functions (P. Buser for closed Ω with $c_2 = \frac{(n-1)^2}{4}$ and C-Maerten for general domains).

1. If $A = 0$, we get

$$\lambda_k(\Delta_A, \Omega) \leq c_2(n)a^2 + c_3(n) \left(\frac{k}{|\Omega|} \right)^{2/n},$$

which is the same estimate as for functions (P. Buser for closed Ω with $c_2 = \frac{(n-1)^2}{4}$ and C-Maerten for general domains).

2. The term $C_1(\Omega, A)$ contains all the contribution associated to the presence of A .

1. If $A = 0$, we get

$$\lambda_k(\Delta_A, \Omega) \leq c_2(n)a^2 + c_3(n) \left(\frac{k}{|\Omega|} \right)^{2/n},$$

which is the same estimate as for functions (P. Buser for closed Ω with $c_2 = \frac{(n-1)^2}{4}$ and C-Maerten for general domains).

2. The term $C_1(\Omega, A)$ contains all the contribution associated to the presence of A .
3. The estimate is compatible with the Weyl law.

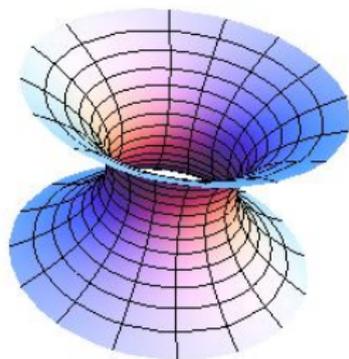
Lower bounds for λ_1 .

We will give lower bounds for λ_1 in the very specific situation where A is a closed form (that is $B = 0$) and the manifold is a *cylinder*

Lower bounds for λ_1 .

We will give lower bounds for λ_1 in the very specific situation where A is a closed form (that is $B = 0$) and the manifold is a *cylinder*

A *Riemannian cylinder* is a domain (Ω, g) diffeomorphic to $[0, 1] \times \mathbb{S}^1$, endowed with a Riemannian metric g . We denote by Σ_1 and Σ_2 the boundaries of the cylinder.



We foliate the cylinder by the (regular) level curves of a smooth function ψ .

Let \mathcal{F}_Ω , the family of smooth real-valued functions on Ω which have no critical points in Ω and which are constant on each component of the boundary of Ω .

We foliate the cylinder by the (regular) level curves of a smooth function ψ .

Let \mathcal{F}_Ω , the family of smooth real-valued functions on Ω which have no critical points in Ω and which are constant on each component of the boundary of Ω .

If $\psi \in \mathcal{F}_\Omega$, we set:

$$K = K_{\Omega, \psi} = \frac{\sup_{\Omega} |\nabla \psi|}{\inf_{\Omega} |\nabla \psi|}.$$

We foliate the cylinder by the (regular) level curves of a smooth function ψ .

Let \mathcal{F}_Ω , the family of smooth real-valued functions on Ω which have no critical points in Ω and which are constant on each component of the boundary of Ω .

If $\psi \in \mathcal{F}_\Omega$, we set:

$$K = K_{\Omega, \psi} = \frac{\sup_{\Omega} |\nabla \psi|}{\inf_{\Omega} |\nabla \psi|}.$$

It is clear that, in the definition of the constant K , we can assume that the range of ψ is the interval $[0, 1]$, and that $\psi = 0$ on Σ_1 and $\psi = 1$ on Σ_2 .

Theorem

Let (Ω, g) be a Riemannian cylinder, and let A be a closed 1-form on Ω . Assume that Ω is K -foliated by the level curves of the smooth function $\psi \in \mathcal{F}_\Omega$. Then:

$$\lambda_1(\Omega, A) \geq \frac{4\pi^2}{KL^2} \cdot d(\Phi^A, \mathbf{Z})^2,$$

where L is the maximum length of a level curve of ψ and Φ^A is the flux of A across any of the boundary components of Ω .

Theorem

Let (Ω, g) be a Riemannian cylinder, and let A be a closed 1-form on Ω . Assume that Ω is K -foliated by the level curves of the smooth function $\psi \in \mathcal{F}_\Omega$. Then:

$$\lambda_1(\Omega, A) \geq \frac{4\pi^2}{KL^2} \cdot d(\Phi^A, \mathbf{Z})^2,$$

where L is the maximum length of a level curve of ψ and Φ^A is the flux of A across any of the boundary components of Ω .

Equality holds if and only if the cylinder Ω is a Riemannian product.

Theorem

Let (Ω, g) be a Riemannian cylinder, and let A be a closed 1-form on Ω . Assume that Ω is K -foliated by the level curves of the smooth function $\psi \in \mathcal{F}_\Omega$. Then:

$$\lambda_1(\Omega, A) \geq \frac{4\pi^2}{KL^2} \cdot d(\Phi^A, \mathbf{Z})^2,$$

where L is the maximum length of a level curve of ψ and Φ^A is the flux of A across any of the boundary components of Ω .

Equality holds if and only if the cylinder Ω is a Riemannian product.

Note that $K \geq 1$; we will see that in many interesting situations (for example, for revolution cylinders) one has in fact $K = 1$. However, in full generality, it is difficult to estimate K .

A case where we can get a good estimate of K :

A case where we can get a good estimate of K :

Let Ω be a topological annulus in \mathbb{R}^2 bounded by the inner curve Σ_1 and the outer curve Σ_2 , both convex. Let $\Phi^A = \frac{1}{2\pi} \int_c A$, where c is the closed curve around the hole. Then, we have

A case where we can get a good estimate of K :

Let Ω be a topological annulus in \mathbb{R}^2 bounded by the inner curve Σ_1 and the outer curve Σ_2 , both convex. Let $\Phi^A = \frac{1}{2\pi} \int_c A$, where c is the closed curve around the hole. Then, we have

$$\lambda_1(\Omega, A) \geq \frac{4\pi^2\beta^2}{B^2L^2} d(\Phi^A, \mathbb{Z})^2$$

where β denotes the minimum of the distance between Σ_1 and Σ_2 , B the maximum of the distance between Σ_1 and Σ_2 and L the length of the outer boundary.

It is clear that need to take account of L and B .

It is clear that need to take account of L and B .

We also need to take account of β .

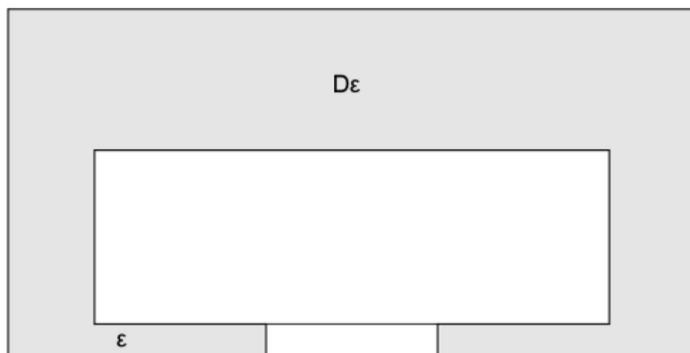


Figure : $\lambda_1 \rightarrow 0$ as $\epsilon \rightarrow 0$

We need the convexity

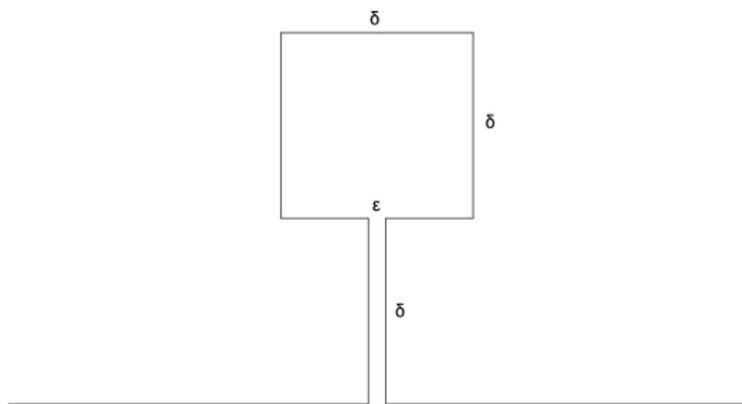


Figure : A local deformation implying $\lambda_1 \rightarrow 0$

Sketch of the proofs.

Upper bounds for λ_1 :

Case 1: if $A = h$ with $dh = \delta h = 0$ and, moreover, $\frac{1}{2\pi} \int_c h \in \mathbb{Z}$ for all closed curved c , we have

$$A = h = n_1 A_1 + \dots + n_k A_k$$

where (A_1, \dots, A_k) is the dual basis of harmonic forms and $n_1, \dots, n_k \in \mathbb{Z}$.

Case 1: if $A = h$ with $dh = \delta h = 0$ and, moreover, $\frac{1}{2\pi} \int_c h \in \mathbb{Z}$ for all closed curved c , we have

$$A = h = n_1 A_1 + \dots + n_k A_k$$

where (A_1, \dots, A_k) is the dual basis of harmonic forms and $n_1, \dots, n_k \in \mathbb{Z}$.

It is known that $\lambda_1(\Delta_A) = 0$, and it is easy to exhibit the eigenfunction:

Case 1: if $A = h$ with $dh = \delta h = 0$ and, moreover, $\frac{1}{2\pi} \int_c h \in \mathbb{Z}$ for all closed curved c , we have

$$A = h = n_1 A_1 + \dots + n_k A_k$$

where (A_1, \dots, A_k) is the dual basis of harmonic forms and $n_1, \dots, n_k \in \mathbb{Z}$.

It is known that $\lambda_1(\Delta_A) = 0$, and it is easy to exhibit the eigenfunction:

If x_0 is a given point of Ω , let

$$\phi(x) = \int_{x_0}^x h.$$

Case 1: if $A = h$ with $dh = \delta h = 0$ and, moreover, $\frac{1}{2\pi} \int_c h \in \mathbb{Z}$ for all closed curved c , we have

$$A = h = n_1 A_1 + \dots + n_k A_k$$

where (A_1, \dots, A_k) is the dual basis of harmonic forms and $n_1, \dots, n_k \in \mathbb{Z}$.

It is known that $\lambda_1(\Delta_A) = 0$, and it is easy to exhibit the eigenfunction:

If x_0 is a given point of Ω , let

$$\phi(x) = \int_{x_0}^x h.$$

Then, $\phi(x)$ does depend on the path between x_0 and x only up to a factor 2π , and

$$u(x) = e^{i\phi(x)}$$

is an eigenfunction.

Case 2: $A = h + \delta\psi$.

Case 2: $A = h + \delta\psi$.

Recall that

$$\lambda_1(\Delta_A) = \min\{R(u) : u \neq 0\},$$

where

$$R(u) = \frac{\int_{\Omega} |\nabla^A u|^2 d\text{vol}_g}{\int_{\Omega} |u|^2 d\text{vol}_g}.$$

Case 2: $A = h + \delta\psi$.

Recall that

$$\lambda_1(\Delta_A) = \min\{R(u) : u \neq 0\},$$

where

$$R(u) = \frac{\int_{\Omega} |\nabla^A u|^2 d\text{vol}_g}{\int_{\Omega} |u|^2 d\text{vol}_g}.$$

We have to choose a *good* test function:

Case 2: $A = h + \delta\psi$.

Recall that

$$\lambda_1(\Delta_A) = \min\{R(u) : u \neq 0\},$$

where

$$R(u) = \frac{\int_{\Omega} |\nabla^A u|^2 d\text{vol}_g}{\int_{\Omega} |u|^2 d\text{vol}_g}.$$

We have to choose a *good* test function:

In general, $h \notin \mathcal{L}_Z$. We choose $\omega \in \mathcal{L}_Z$ minimizing $d(A, \mathcal{L}_Z)$, and consider the same function as in case 1:

We take

$$u(x) = e^{i\phi(x)} \text{ with } \phi(x) = \int_{x_0}^x \omega.$$

We take

$$u(x) = e^{i\phi(x)} \text{ with } \phi(x) = \int_{x_0}^x \omega.$$

We get

$$\nabla^A(u) = du - iAu = i\omega u - ihu - i\delta\psi u = iu((\omega - h) - \delta\psi).$$

We take

$$u(x) = e^{i\phi(x)} \text{ with } \phi(x) = \int_{x_0}^x \omega.$$

We get

$$\nabla^A(u) = du - iAu = i\omega u - ihu - i\delta\psi u = iu((\omega - h) - \delta\psi).$$

We have, using $|u| = 1$,

$$R(u) = \frac{\int_{\Omega} |\nabla^A u|^2 d\text{vol}_g}{\int_{\Omega} |u|^2 d\text{vol}_g} = \frac{\|(\omega - h) - \delta\psi\|^2}{|\Omega|}$$

and

We take

$$u(x) = e^{i\phi(x)} \text{ with } \phi(x) = \int_{x_0}^x \omega.$$

We get

$$\nabla^A(u) = du - iAu = i\omega u - ihu - i\delta\psi u = iu((\omega - h) - \delta\psi).$$

We have, using $|u| = 1$,

$$R(u) = \frac{\int_{\Omega} |\nabla^A u|^2 d\text{vol}_g}{\int_{\Omega} |u|^2 d\text{vol}_g} = \frac{\|(\omega - h) - \delta\psi\|^2}{|\Omega|}$$

and

$$\lambda_1(\Delta_A) \leq \frac{\|w - h\|^2 + \|\delta\psi\|^2}{|\Omega|}.$$

By the choice of ω , we have $\|w - h\|^2 = d^2(h, \mathcal{L}_Z)$.

By the choice of ω , we have $\|w - h\|^2 = d^2(h, \mathcal{L}_Z)$.

By the min-max,

$$\lambda_{1,1}(\Omega) \leq \frac{\int_{\Omega} |d(\delta\psi)|^2 d\text{vol}_g}{\int_{\Omega} |\delta\psi|^2 d\text{vol}_g},$$

and

By the choice of ω , we have $\|w - h\|^2 = d^2(h, \mathcal{L}_Z)$.

By the min-max,

$$\lambda_{1,1}(\Omega) \leq \frac{\int_{\Omega} |d(\delta\psi)|^2 d\text{vol}_g}{\int_{\Omega} |\delta\psi|^2 d\text{vol}_g},$$

and

$$\int_{\Omega} |\delta\psi|^2 d\text{vol}_g \leq \frac{\|B\|^2}{\lambda_{1,1}(\Omega)}.$$

By the choice of ω , we have $\|w - h\|^2 = d^2(h, \mathcal{L}_Z)$.

By the min-max,

$$\lambda_{1,1}(\Omega) \leq \frac{\int_{\Omega} |d(\delta\psi)|^2 d\text{vol}_g}{\int_{\Omega} |\delta\psi|^2 d\text{vol}_g},$$

and

$$\int_{\Omega} |\delta\psi|^2 d\text{vol}_g \leq \frac{\|B\|^2}{\lambda_{1,1}(\Omega)}.$$

This implies

$$\lambda_1(\Delta_A) \leq \frac{d(A, \mathcal{L}_Z)^2}{|\Omega|} + \frac{\|B\|^2}{\lambda_{1,1}(\Omega)|\Omega|}.$$

Case 3: the other eigenvalues.

Case 3: the other eigenvalues.

We use the same strategy as in the second step. We choose $\omega \in \mathcal{L}_{\mathbf{Z}}$ minimizing $d(A, \mathcal{L}_{\mathbf{Z}})$, and consider the function

$$u(x) = e^{i\phi(x)} \text{ with } \phi(x) = \int_{x_0}^x \omega.$$

Case 3: the other eigenvalues.

We use the same strategy as in the second step. We choose $\omega \in \mathcal{L}_{\mathbf{Z}}$ minimizing $d(A, \mathcal{L}_{\mathbf{Z}})$, and consider the function

$$u(x) = e^{i\phi(x)} \text{ with } \phi(x) = \int_{x_0}^x \omega.$$

The test functions will be of the type fu where f is a real smooth function on Ω .

We have

$$\begin{aligned}(d - iA)(fu) &= udf + fdu - ihuf - iuf\delta\psi \\ &= udf + iuf(\omega - h - \delta\psi).\end{aligned}$$

We have

$$\begin{aligned}(d - iA)(fu) &= udf + fdu - ihuf - iuf\delta\psi \\ &= udf + iuf(\omega - h - \delta\psi).\end{aligned}$$

Since $|u| = 1$:

$$|(d - iA)(fu)|^2 \leq 2(|df|^2 + f^2|\omega - h - \delta\psi|^2).$$

We have

$$\begin{aligned}(d - iA)(fu) &= udf + fdu - ihuf - iuf\delta\psi \\ &= udf + iuf(\omega - h - \delta\psi).\end{aligned}$$

Since $|u| = 1$:

$$|(d - iA)(fu)|^2 \leq 2\left(|df|^2 + f^2|\omega - h - \delta\psi|^2\right).$$

We have to control the Rayleigh quotient

$$R(fu) \leq 2\left(\frac{\int_{\Omega} |df|^2}{\int_{\Omega} f^2} + \frac{\int_{\Omega} f^2|\omega - h - \delta\psi|^2}{\int_{\Omega} f^2}\right).$$

So, we are lead to control the Rayleigh quotient

$$R(f) = 2 \frac{\int_{\Omega} |df|^2 + Vf^2}{\int_{\Omega} f^2}, \quad \text{where } V = |\omega - h - \delta\psi|^2.$$

So, we are lead to control the Rayleigh quotient

$$R(f) = 2 \frac{\int_{\Omega} |df|^2 + Vf^2}{\int_{\Omega} f^2}, \quad \text{where } V = |\omega - h - \delta\psi|^2.$$

Thus, the problem is now to find an upper bound for the spectrum of the operator $\Delta + V$, where Δ is the usual Laplacian acting on functions and $V = |\omega - h - \delta\psi|^2$ is a nonnegative potential.

The proof follows word for word what is done in the case where the potential is equal to 0.

The proof follows word for word what is done in the case where the potential is equal to 0.

The idea is to construct disjointly supported domains $\Omega_1, \dots, \Omega_k$ on Ω and to associate a test function for the Rayleigh quotient to each of these domains.



The fact that the potential V is positif implies that we can choose the domains Ω_i with

$$\int_{\Omega_i} V d\text{vol}_g \leq c \frac{\int_{\Omega} V}{k}.$$

Lower bound on Riemannian cylinders.

Lower bound on Riemannian cylinders.

The proof are technical, in particular the equality case.

Lower bound on Riemannian cylinders.

The proof are technical, in particular the equality case.

We use the foliation as follow: we restrict the potential A on the cylinder to each circle of the foliation, and this allows us to estimate the spectrum of the Riemannian cylinder in comparison with the spectrum of circles.

Lower bound on Riemannian cylinders.

The proof are technical, in particular the equality case.

We use the foliation as follow: we restrict the potential A on the cylinder to each circle of the foliation, and this allows us to estimate the spectrum of the Riemannian cylinder in comparison with the spectrum of circles.

The presence of K reflects the complexity of the geometry of the cylinder.

Thank you!

