# Hamiltonian stability

### over discrete spaces

#### Franco Vivaldi Queen Mary, University of London





integrable

# foliation by invariant curves



integrable

## foliation by invariant curves



integrable near-integrable: stable



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near-integrable: stable

strong perturbation: unstable

integrable



What happens if the space is discrete?



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Truncation (computer arithmetic).

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- Geometric discretization.
- Reduction to a finite field.

Astron. & Astrophys. 31, 289-301 (1974)

#### Numerical Study of Discrete Plane Area-preserving Mappings

F. Rannou Observatoire de Nice

Received August, 10, 1973

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$$x_{t+1} = x_t + y_t + \frac{m}{2\pi} \left( 1 - \cos \frac{2\pi}{m} y_t \right) \pmod{m}$$

$$y_{t+1} = y_t - \frac{\lambda m}{2\pi} \left( \sin \frac{2\pi}{m} x_{t+1} + 1 - \cos \frac{2\pi}{m} x_{t+1} \right) \pmod{m}$$

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$$\begin{aligned} & \text{lattice size} \\ x_{t+1} &= x_t + y_t + \frac{m}{2\pi} \left( 1 - \cos \frac{2\pi}{m} y_t \right) & (\text{mod } m) \\ y_{t+1} &= y_t - \frac{\lambda m}{2\pi} \left( \sin \frac{2\pi}{m} x_{t+1} + 1 - \cos \frac{2\pi}{m} x_{t+1} \right) & (\text{mod } m) \end{aligned}$$

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### First study of an invertible lattice map on the torus



All orbits are periodic.

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lattice size

rounding to

nearest integer

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PPPL-1938 (Nov. 1982) Physica 8D(3), 360-380 (Sept. 1983)

#### LONG-TIME CORRELATIONS IN THE STOCHASTIC REGIME

Charles F. F. KARNEY Plasma Physics Laboratory, Princeton University, Princeton, New Jersey 08544, USA

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$$x_{t+1} = x_t + y_{t+1}$$
  
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An algebraic mapping forced on a lattice via truncation

$$x_{t+1} = x_t + y_{t+1}$$
  
 $y_{t+1} = y_t + 2(x_t^2 - \kappa)$ 

Detailed study of the boundary between order and chaos
Very large lattice (2<sup>48</sup> x 2<sup>48</sup> points).
90 hours of CPU on the Cray-I computer.
Detected non-exponential correlation decay caused by structure of chaos boundary.

Physica D 56 (1992) 1-22 North-Holland

#### PHYSICA 🛛

#### Exact numerical studies of Hamiltonian maps: Iterating without roundoff error

David J.D. Earn

Institute of Astronomy, Madingley Road, Cambridge CB3 6HA, UK

and

Scott Tremaine

Canadian Institute for Theoretical Astrophysics, University of Toronto, Toronto, Canada M5S 1A1<sup>2</sup> and Theoretical Astrophysics 130-33, California Institute of Technology, Pasadena, CA 91125, USA

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### Floating-point numbers

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sign exponent

#### mantissa

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 $\pm 2^{exponent-q} \times mantissa$ 

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Same number of points between any two consecutive powers of 2.



Highly non-uniform distribution.



Injective maps become strongly non-injective.



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- Irreversible behaviour: long transients, short limit cycles.



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- Irreversible behaviour: long transients, short limit cycles.
- Points of a limit cycle arrange themselves according to the invariant measure.



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**Standard model:**  $F_{\lambda} : \mathbb{Z}^2 \to \mathbb{Z}^2$   $(x, y) \mapsto (\lfloor \lambda x \rfloor - y, x)$   $|\lambda| < 2$ 

(fv, Bosio, Lowenstein, Vladimirov, Akiyama)

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#### What do we know?

If  $\lambda$  is rational with a single prime divisor p at denominator, then the round-off map may be embedded into a positive-entropy map of the ring of p-adic integers. [Bosio, fv]

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**Observations:** 



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Difficult to compute.



#### **Observations:**

- Difficult to compute.
- Very large fluctuations, within a highly organised structure.











- The departure of round-off orbits from exact orbits obeys a central limit theorem. [Vladimirov, fv]
  - The period function is defined at infinitely many points on the symmetry axis. [Akiyama]



# ?

Look for different discrete structures, for which bounding invariant sets exist.



- Look for different discrete structures, for which bounding invariant sets exist.
- Look for systems with more manageable (zero entropy) fluctuations.

A flow with a piecewise-constant vector field is diffracted by a line







replace the flow by a map, using the same vectors

flow and map differ within a strip Linked strip/maps: outer billiards of polygons

Linked strip maps: outer billiards of polygons
Linked strip maps: outer billiards of polygons

x

Linked strip maps: outer billiards of polygons

x

Linked strip maps: outer billiards of polygons

















PLOT: 0.849 FEET.

LOD. BOOK SCALE.

83/06/18. 19.37.34.





100.0001 SCALE.

83/06/18. 19.37.34.







Under suitable rationality conditions, connected invariant chains of polygons exist ("necklages"), which bound all orbits inside.





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Invariant chains of polygons form countable families, ensuring global stability.

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If the rationality conditions are not satisfied, then there are unbounded orbits (Schwartz, 2007).

(Zhang, Alwani, fv)

$$y_{t+1} \equiv y_t + \varepsilon(x_t) \pmod{N}$$
  
 $x_{t+1} \equiv x_t + y_{t+1} \pmod{N}$ 

$$\varepsilon(x) = \begin{cases} +1 & 0 \leq x < \lfloor N/2 \rfloor \\ -1 & \lfloor N/2 \rfloor \leq x < N \end{cases}$$

#### (Zhang, Alwani, fv)





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what is the motion inside an island?

There is a small number of very long cycles.

Islands of stability exist, for selected rotation numbers.

$$\mathscr{F}: \mathbb{Z}^2 \to \mathbb{Z}^2 \qquad \begin{array}{rcl} y_{t+1} &=& y_t - \operatorname{sign}(x_t) \\ x_{t+1} &=& x_t + \alpha y_{t+1} + \beta \end{array} \qquad 0 \leqslant \beta < \alpha$$

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The parameters  $\alpha$  and  $\beta$  depend on the rotation number of the island.





**Conjecture** (Zhang, fv). If  $gcd(\alpha, 2\beta)$  is odd, then all orbits are periodic, and their period, sufficiently far from the origin, eventually becomes constant. Otherwise all orbits are unbounded.



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The conjecture is true for the scaled system (from symmetry and uniform distribution modulo  $\alpha$ ).



0.015


## Thank you for your attention