

On Some Variational Optimization Problems in Classical Fluids and Superfluids

Bartosz Protas

Department of Mathematics & Statistics
McMaster University, Hamilton, Ontario, Canada
URL: <http://www.math.mcmaster.ca/bprotas>

Funded by Early Researcher Award (ERA) and NSERC
Computational Time Provided by SHARCNET

New Challenges in Mathematical Modelling and Numerical
Simulation of Superfluids, CIRM 2016

Agenda

Minimization of the Gross-Pitaevskii Energy Functional

- Formulation of the Problem
- Gradient Minimization
- Sobolev Gradients

Riemannian Optimization

- First-Order Geometry
- Second-Order Geometry
- Riemannian Conjugate Gradients

Probing Fundamental Bounds in Hydrodynamics

- Sharpness of Estimates as Optimization Problem
- Bounds for 2D Navier-Stokes Problem
- Computational Approach & Results

Collaborators

- ▶ **Ionut Danaila**
(Université de Rouen)
- ▶ **Diego Ayala**
(former Ph.D. student, now at University of Michigan)
- ▶ **Charles Doering**
(University of Michigan)

PROBLEM I

RIEMANNIAN OPTIMIZATION FOR COMPUTATION GROUND STATES IN BOSE-EINSTEIN CONDENSATES

(joint work with I. Danaila)

► Gross-Pitaevskii Free Energy Functional (non-dimensional form)

$$E(u) = \int_{\mathcal{D}} \left[\frac{1}{2} |\nabla u|^2 + C_{\text{trap}} |u|^2 + \frac{1}{2} C_g |u|^4 - i C_{\Omega} u^* A^t \nabla u \right] d\mathbf{x},$$

$$\|u\|_2^2 = \int_{\mathcal{D}} |u(\mathbf{x})|^2 d\mathbf{x} = 1, \quad \mathcal{D} \subseteq \mathbb{R}^d$$

where

$$u = \frac{\psi}{\sqrt{N} x_s^{-d/2}}, \quad \psi \text{ — wavefunction, } \psi : \mathcal{D} \rightarrow \mathbb{C}$$

N — number of atoms in the condensate

x_s — characteristic length scale

$$A^t = [y, -x, 0], \quad C_{\text{trap}}(x, y, z) \text{ — trapping potential}$$

C_g, C_{Ω} — constants

► C_{Ω} characterizes the effect of rotation

- ▶ Dirichlet boundary conditions: $u = 0$ on $\partial\mathcal{D}$
- ▶ Variational optimization, $E : H_0^1(D) \rightarrow \mathbb{R}$

$$\begin{aligned} & \min_{u \in H_0^1(\mathcal{D})} E(u) \\ & \text{subject to } \|u\|_{L_2(\mathcal{D})} = 1 \end{aligned}$$

- ▶ Minimizers constrained to a nonlinear manifold \mathcal{M} in $H_0^1(\mathcal{D})$

$$\mathcal{M} := \{u \in H_0^1(\mathcal{D}) : \|u\|_{L_2(\mathcal{D})} = 1\}$$

- ▶ Computational approaches:
 - ▶ Euler-Lagrange equation for $E(u) \implies$ nonlinear eigenvalue problem
 - ▶ Direct minimization of $E(u)$ via a gradient method

► Steepest-gradient approach

$$\begin{aligned} u^{(n+1)} &= u^{(n)} - \tau_n \nabla E(u^{(n)}), & n = 0, 1, \dots, \\ u^{(0)} &= u_0, & \text{(initial guess),} \end{aligned}$$

where:

$$\tilde{u} = \lim_{n \rightarrow \infty} u^{(n)} \quad \text{— the minimizer (“ground state”)}$$

$$\nabla E(u^{(n)}) \quad \text{— gradient of } E(u) \text{ at } u^{(n)}$$

$$\tau_n = \operatorname{argmin}_{\tau > 0} E(u^{(n)} - \tau \nabla E(u^{(n)})) \quad \text{— optimal step size}$$

► Key issues:

- Regularity of the minimizers $\tilde{u} \in H_0^1(\mathcal{D}) \implies$ Sobolev gradients
- Enforcement of the constraint $\tilde{u} \in \mathcal{M} \implies$ Riemannian optimization

- Gâteaux differential of the Gross-Pitaevskii Energy Functional

$$E'(u; v) = \lim_{\epsilon \rightarrow 0} \epsilon^{-1} [E(u + \epsilon v) - E(u)], \quad u, v \in \mathcal{X}$$

\mathcal{X} — some function space

- Riesz Representation Theorem:

$$\begin{aligned} E'(u; \cdot) \text{ bounded linear functional on } \mathcal{X} \\ \implies \forall v \in \mathcal{X} \quad E'(u; v) = \langle \nabla^{\mathcal{X}} E(u), v \rangle_{\mathcal{X}} \end{aligned}$$

- Relevant inner products (Danaila & Kazemi 2010)

$$\langle u, v \rangle_{L_2} = \int_{\mathcal{D}} \langle u, v \rangle \, d\mathbf{x}, \quad \text{where } \langle u, v \rangle = uv^*$$

$$\langle u, v \rangle_{H^1} = \int_{\mathcal{D}} \langle u, v \rangle + \langle \nabla u, \nabla v \rangle \, d\mathbf{x}$$

$$\langle u, v \rangle_{H_A} = \int_{\mathcal{D}} \langle u, v \rangle + \langle \nabla_A u, \nabla_A v \rangle \, d\mathbf{x}, \quad \nabla_A = \nabla + iC_{\Omega} A^t$$

- Different Sobolev gradients ($\mathcal{X} = L_2, H^1, H_A$)

$$E'(u; v) = \Re \left\langle \nabla^{L_2} E(u), v \right\rangle_{L_2} = \Re \left\langle \nabla^{H^1} E(u), v \right\rangle_{H^1} = \Re \left\langle \nabla^{H_A} E(u), v \right\rangle_{H_A}$$

- The L_2 gradient

$$\nabla^{L_2} E(u) = 2 \left(-\frac{1}{2} \nabla^2 u + C_{\text{trap}} u + C_g |u|^2 u - i C_\Omega A^t \nabla u \right),$$

- The Sobolev gradient $G = \nabla_{H_A} E(u)$ obtained from the L_2 gradient via an elliptic boundary-value problem

$$\begin{aligned} \forall_{v \in H_0^1(\mathcal{D})} \int_{\mathcal{D}} \left[(1 + C_\Omega^2(x^2 + y^2)) Gv + \nabla G \cdot \nabla v - 2i C_\Omega A^t \nabla Gv \right] d\mathbf{x} \\ = \int_{\mathcal{D}} \frac{1}{2} \nabla u \cdot \nabla v + [C_{\text{trap}} u + C_g |u|^2 u - i C_\Omega A^t \nabla u] v d\mathbf{x} \end{aligned}$$

- ▶ Riemannian Optimization is general approach based on differential geometry
 - ▶ here made simple by the constraint $\|u\|_{L_2(\mathcal{D})} = 1$
- ▶ “Intrinsic” approach with optimization performed directly on the manifold \mathcal{M} without reference to the space $H^1(\mathcal{D})$
 - ▶ optimization problem becomes *unconstrained*
 - ▶ can apply more efficient optimization algorithms (conjugate gradients, Newton’s method)
- ▶ Reference: P.-A. Absil, R. Mahony and R. Sepulchre, “Optimization Algorithms on Matrix Manifolds”, Princeton University Press, (2008).

- ▶ Projection of the gradient G on the tangent subspace $\mathcal{T}_u\mathcal{M}$

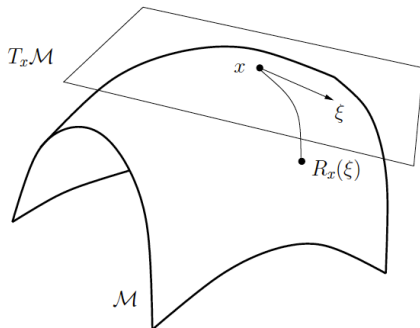
$$P_{u_n, H_A} G = G - \frac{\Re(\langle u_n, G \rangle_{L^2})}{\Re(\langle u_n, v_{H_A} \rangle_{L^2})} v_{H_A}, \quad \text{where}$$
$$\langle v_{H_A}, v \rangle_{H_A} = \langle u_n, v \rangle_{L^2}, \quad \forall v \in H_A$$

- ▶ There is some freedom in choosing the subtracted field (v_{H_A})
- ▶ Approach equivalent to constraint enforcement via Lagrange multipliers
 - ▶ Error in constraint satisfaction $\mathcal{O}(\tau_n)$

► RETRACTION

$$\mathcal{R}_u : \mathcal{T}_u\mathcal{M} \rightarrow \mathcal{M}$$

maps a tangent vector $\xi \in \mathcal{T}_u\mathcal{M}$ back to the manifold \mathcal{M}



- For our constraint manifold \mathcal{M}

$$\mathcal{R}_u(\xi) = \frac{u + \xi}{\|u + \xi\|_{L_2(\mathcal{D})}}$$

retraction = normalization

- Riemannian steepest descent approach

$$\begin{aligned} u_{n+1} &= \mathcal{R}_{u_n}(\tau_n P_{u_n, H_A} G(u_n)), \quad n = 0, 1, 2, \dots \\ u_0 &= u^0 \end{aligned}$$

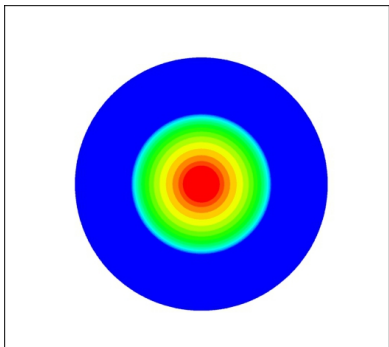
where

$$\tau_n = \operatorname{argmin}_{\tau > 0} E(\mathcal{R}_{u_n}(\tau P_{u_n, H_A} G(u_n)))$$

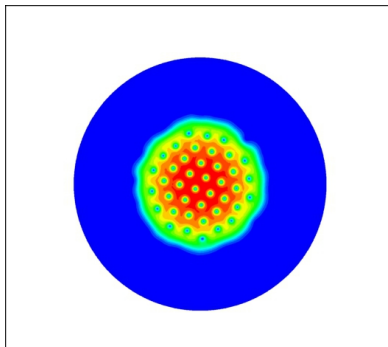
Results: Abrikosov Lattice ($C_g = 1000$, $C_\Omega = 0.9$)

- ▶ Finite-element approximation of energy functional and PDE operators
 - ▶ implementation in FreeFEM++ (with 10^5 P1 elements)

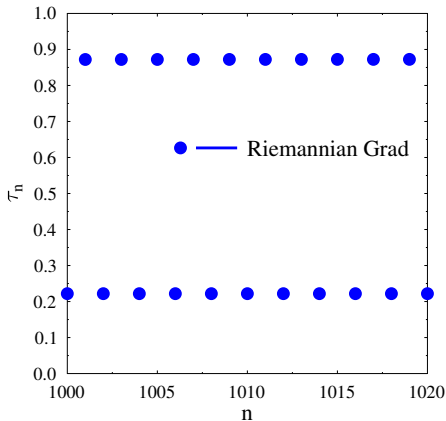
▶



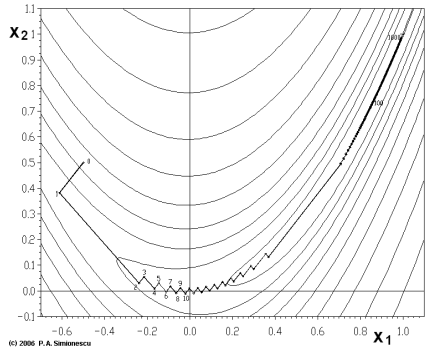
$|u|$, initial guess
(Thomas-Fermi approximation)



$|u|$, final state



step size τ_n



steepest descent for
 “banana function”

► Consider $\min_{\mathbf{x} \in \mathbb{R}^N} f(\mathbf{x})$, where $f : \mathbb{R}^N \rightarrow \mathbb{R}$

► **Nonlinear Conjugate Gradients Method**

$$\begin{aligned}\mathbf{x}_{n+1} &= \mathbf{x}_n + \tau_n \mathbf{d}_n, & n = 0, 1, \dots \\ \mathbf{x}_0 &= \mathbf{x}^0\end{aligned}$$

► descent direction \mathbf{d}_n is defined as

$$\begin{aligned}\mathbf{d}_n &= -\mathbf{g}_n + \beta_n \mathbf{d}_{n-1}, & n = 1, 2, \dots \\ \mathbf{d}_0 &= -\mathbf{g}_0, & \mathbf{g}_n = \nabla f(\mathbf{x}_n)\end{aligned}$$

► “momentum” coefficients β_n ensure conjugacy of decent directions

$$\beta_n = \beta_n^{FR} := \frac{\langle \mathbf{g}_n, \mathbf{g}_n \rangle_{\mathcal{X}}}{\langle \mathbf{g}_{n-1}, \mathbf{g}_{n-1} \rangle_{\mathcal{X}}} \quad (\text{Fletcher-Reeves}),$$

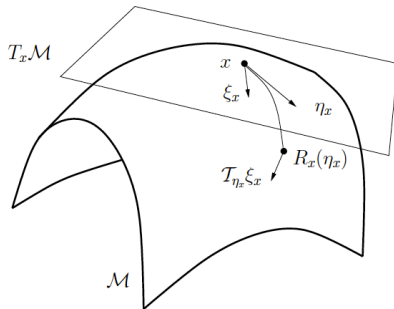
$$\beta_n = \beta_n^{PR} := \frac{\langle \mathbf{g}_n, (\mathbf{g}_n - \mathbf{g}_{n-1}) \rangle_{\mathcal{X}}}{\langle \mathbf{g}_{n-1}, \mathbf{g}_{n-1} \rangle_{\mathcal{X}}} \quad (\text{Polak-Ribière})$$

- In the Riemannian setting

$$\mathbf{g}_{n-1}, \mathbf{d}_{n-1} \in \mathcal{T}_{\mathbf{x}_{n-1}} \quad \text{and} \quad \mathbf{g}_n, \mathbf{d}_n \in \mathcal{T}_{\mathbf{x}_n},$$

hence cannot be added or multiplied ...

- Need a mapping between the tangent spaces $\mathcal{T}_{u_{n-1}}\mathcal{M}$ and $\mathcal{T}_{u_n}\mathcal{M}$
- **VECTOR TRANSPORT** $\mathcal{T}_\eta(\xi) : \mathcal{T}\mathcal{M} \times \mathcal{T}\mathcal{M} \rightarrow \mathcal{T}\mathcal{M}$, $\xi, \eta \in \mathcal{T}\mathcal{M}$
describing how the vector field ξ is transported along the manifold \mathcal{M}
by the field η



Conclusions (I)

- ▶ Riemannian optimization accelerates the computation of BEC ground states
- ▶ Key enablers for Riemannian Conjugate Gradients:
 - ▶ projections onto $\mathcal{T}_{u_n}\mathcal{M}$
 - ▶ retractions from $\mathcal{T}_{u_n}\mathcal{M}$ onto \mathcal{M} ,
 - ▶ vector transport between $\mathcal{T}_{u_{n-1}}$ and \mathcal{T}_{u_n}
- ▶ Riemannian Newton's method?

PROBLEM II

PROBING FUNDAMENTAL BOUNDS IN HYDRODYNAMICS

(joint work with D. Ayala and Ch. Doering)

- ▶ Navier-Stokes equation ($\Omega = [0, L]^d$, $d = 2, 3$)

$$\begin{cases} \frac{\partial \mathbf{v}}{\partial t} + (\mathbf{v} \cdot \nabla) \mathbf{v} + \nabla p - \nu \Delta \mathbf{v} = \mathbf{0}, & \text{in } \Omega \times (0, T] \\ \nabla \cdot \mathbf{v} = 0, & \text{in } \Omega \times (0, T] \\ \mathbf{v} = \mathbf{v}_0 & \text{in } \Omega \text{ at } t = 0 \\ \text{Boundary Condition} & \text{on } \Gamma \times (0, T] \end{cases}$$

- ▶ 2D Case

- ▶ Existence Theory Complete — smooth and unique solutions exist for arbitrary times and arbitrarily large data

- ▶ 3D Case

- ▶ Weak solutions (possibly nonsmooth) exist for arbitrary times
- ▶ Classical (smooth) solutions (possibly nonsmooth) exist for *finite* times only
- ▶ Possibility of “blow-up” (finite-time singularity formation)
- ▶ One of the Clay Institute “Millennium Problems” (\$ 1M!)
http://www.claymath.org/millennium/Navier-Stokes_Equations

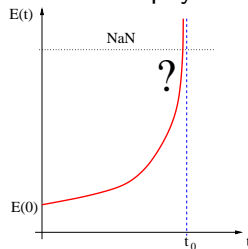
What is known? — Available Estimates

- ▶ A Key Quantity — Enstrophy

$$\mathcal{E}(t) \triangleq \int_{\Omega} |\nabla \times \mathbf{v}|^2 d\Omega \quad (= \|\nabla \mathbf{v}\|_2^2)$$

- ▶ Smoothness of Solutions \iff Bounded Enstrophy
 (Foias & Temam, 1989)

$$\max_{t \in [0, T]} \mathcal{E}(t) < \infty \quad ???$$



- ▶ Can estimate $\frac{d\mathcal{E}(t)}{dt}$ using the momentum equation, Sobolev's embeddings, Young and Cauchy-Schwartz inequalities, ...
 - ▶ REMARK: incompressibility not used in these estimates

► 2D Case:

$$\frac{d\mathcal{E}(t)}{dt} \leq \frac{C^2}{\nu} \mathcal{E}(t)^2$$

- Grönwall's lemma and energy equation yield $\forall_t \mathcal{E}(t) < \infty$
- smooth solutions exist for all times

► 3D Case:

$$\frac{d\mathcal{E}(t)}{dt} \leq \frac{27C^2}{128\nu^3} \mathcal{E}(t)^3$$

- corresponding estimate not available
- upper bound on $\mathcal{E}(t)$ blows up in finite time

$$\mathcal{E}(t) \leq \frac{\mathcal{E}(0)}{\sqrt{1 - 4 \frac{C\mathcal{E}(0)^2}{\nu^3} t}}$$

- singularity in finite time cannot be ruled out!

Problem of Lu & Doering (2008), I

- ▶ Can we actually find solutions which “saturate” a given estimate?
- ▶ Estimate $\frac{d\mathcal{E}(t)}{dt} \leq c\mathcal{E}(t)^3$ at a *fixed* instant of time t

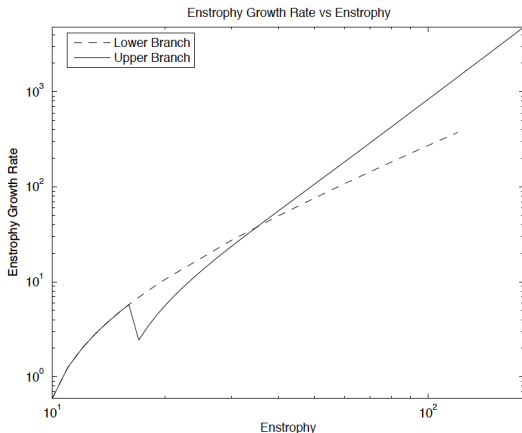
$$\max_{\mathbf{v} \in H^1(\Omega), \nabla \cdot \mathbf{v} = 0} \frac{d\mathcal{E}(t)}{dt}$$

subject to $\mathcal{E}(t) = \mathcal{E}_0$

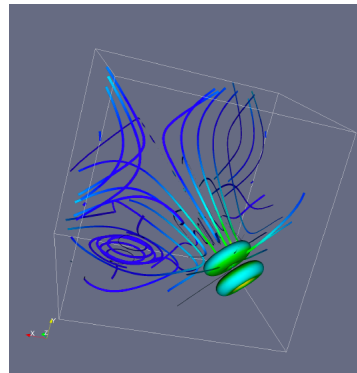
where

- ▶
$$\frac{d\mathcal{E}(t)}{dt} = -\nu \|\Delta \mathbf{v}\|_2^2 + \int_{\Omega} \mathbf{v} \cdot \nabla \mathbf{v} \cdot \Delta \mathbf{v} \, d\Omega$$
- ▶ \mathcal{E}_0 is a parameter
- ▶ Solution using a gradient-based descent method

Problem of Lu & Doering (2008), II



$$\left[\frac{d\mathcal{E}(t)}{dt} \right]_{\max} = 8.97 \times 10^{-4} \mathcal{E}_0^{2.997}$$



vorticity field (top branch)

- ▶ How about solutions which saturate $\frac{d\mathcal{E}(t)}{dt} \leq c\mathcal{E}(t)^3$ over a *finite* time window $[0, T]$?

$$\begin{aligned} & \max_{\mathbf{v}_0 \in H^1(\Omega), \nabla \cdot \mathbf{v} = 0} \mathcal{E}(T) \\ & \text{subject to } \mathcal{E}(0) = \mathcal{E}_0 \end{aligned}$$

where



$$\mathcal{E}(t) = \int_0^t \frac{d\mathcal{E}(\tau)}{d\tau} d\tau + \mathcal{E}_0$$

- ▶ \mathcal{E}_0 and T are parameters

- ▶ In principle doable, but will try something simpler first ...

Relevant Estimates

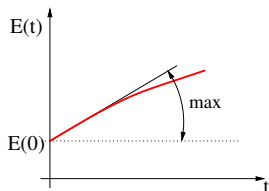
	BEST ESTIMATE	SHARP?
1D Burgers instantaneous	$\frac{d\mathcal{E}(t)}{dt} \leq \frac{3}{2} \left(\frac{1}{\pi^2\nu} \right)^{1/3} \mathcal{E}(t)^{5/3}$	
1D Burgers finite-time	$\max_{t \in [0, T]} \mathcal{E}(t) \leq \left[\mathcal{E}_0^{1/3} + \left(\frac{L}{4} \right)^2 \left(\frac{1}{\pi^2\nu} \right)^{4/3} \mathcal{E}_0 \right]^3$	
2D Navier-Stokes instantaneous	$\begin{aligned} \frac{d\mathcal{P}(t)}{dt} &\leq - \left(\frac{\nu}{\mathcal{E}} \right) \mathcal{P}^2 + C_1 \left(\frac{\mathcal{E}}{\nu} \right) \mathcal{P} \\ \frac{d\mathcal{P}(t)}{dt} &\leq \frac{C_2}{\nu} \mathcal{K}^{1/2} \mathcal{P}^{3/2} \end{aligned}$	
2D Navier-Stokes finite-time	$\max_{t>0} \mathcal{P}(t) \leq \left[\mathcal{P}_0^{1/2} + \frac{C_2}{4\nu^2} \mathcal{K}_0^{1/2} \mathcal{E}_0 \right]^2$	
3D Navier-Stokes instantaneous	$\frac{d\mathcal{E}(t)}{dt} \leq \frac{27C^2}{128\nu^3} \mathcal{E}(t)^3$	YES Lu & Doering (2008)
3D Navier-Stokes finite-time	$\mathcal{E}(t) \leq \frac{\mathcal{E}(0)}{\sqrt{1 - 4 \frac{C\mathcal{E}(0)^2}{\nu^3} t}}$	

► QUESTION #1 (“SMALL”)

Sharpness of *instantaneous* estimates
 (at some *fixed* t)

$$\max_{\mathbf{u}} \frac{d\mathcal{E}}{dt} \quad (1D, 3D)$$

$$\max_{\mathbf{u}} \frac{d\mathcal{P}}{dt} \quad (2D)$$

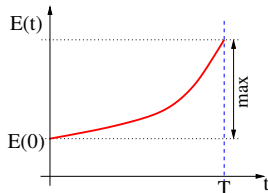


► QUESTION #2 (“BIG”)

Sharpness of *finite-time* estimates
 (at some time window $[0, T]$, $T > 0$)

$$\max_{\mathbf{u}_0} \left[\max_{t \in [0, T]} \mathcal{E}(t) \right] \quad (1D, 3D)$$

$$\max_{\mathbf{u}_0} \left[\max_{t \in [0, T]} \mathcal{P}(t) \right] \quad (2D)$$



Relevant Estimates

	BEST ESTIMATE	SHARP?
1D Burgers instantaneous	$\frac{d\mathcal{E}(t)}{dt} \leq \frac{3}{2} \left(\frac{1}{\pi^2\nu}\right)^{1/3} \mathcal{E}(t)^{5/3}$	YES Lu & Doering (2008)
1D Burgers finite-time	$\max_{t \in [0, T]} \mathcal{E}(t) \leq \left[\mathcal{E}_0^{1/3} + \left(\frac{L}{4}\right)^2 \left(\frac{1}{\pi^2\nu}\right)^{4/3} \mathcal{E}_0 \right]^3$	No Ayala & P. (2011)
2D Navier-Stokes instantaneous	$\frac{d\mathcal{P}(t)}{dt} \leq -\left(\frac{\nu}{\mathcal{E}}\right) \mathcal{P}^2 + C_1 \left(\frac{\mathcal{E}}{\nu}\right) \mathcal{P}$ $\frac{d\mathcal{P}(t)}{dt} \leq \frac{C_2}{\nu} \mathcal{K}^{1/2} \mathcal{P}^{3/2}$	
2D Navier-Stokes finite-time	$\max_{t>0} \mathcal{P}(t) \leq \left[\mathcal{P}_0^{1/2} + \frac{C_2}{4\nu^2} \mathcal{K}_0^{1/2} \mathcal{E}_0 \right]^2$	
3D Navier-Stokes instantaneous	$\frac{d\mathcal{E}(t)}{dt} \leq \frac{27C^2}{128\nu^3} \mathcal{E}(t)^3$	YES Lu & Doering (2008)
3D Navier-Stokes finite-time	$\mathcal{E}(t) \leq \frac{\mathcal{E}(0)}{\sqrt{1 - 4 \frac{C\mathcal{E}(0)^2}{\nu^3} t}}$	

► 2D VORTICITY EQUATION IN A PERIODIC BOX ($\omega = \mathbf{e}_z \cdot \boldsymbol{\omega}$)

$$\frac{\partial \omega}{\partial t} + J(\omega, \psi) = \nu \Delta \omega \quad \text{where } J(f, g) = f_x g_y - f_y g_x$$

$$- \Delta \psi = \omega$$

► Enstrophy uninteresting in 2D flows (w/o boundaries)

$$\frac{1}{2} \frac{d}{dt} \int_{\Omega} \omega^2 d\Omega = -\nu \int_{\Omega} (\nabla \omega)^2 d\Omega < 0$$

► Evolution equation for the vorticity gradient $\nabla \omega$

$$\frac{\partial \nabla \omega}{\partial t} + (\mathbf{u} \cdot \nabla) \nabla \omega = \nu \Delta \nabla \omega + \underbrace{\nabla \omega \cdot \nabla \mathbf{u}}_{\text{"STRETCHING" TERM}}$$

► Palinstrophy

$$\mathcal{P}(t) \triangleq \int_{\Omega} (\nabla \omega(t, \mathbf{x}))^2 d\Omega = \int_{\Omega} (\nabla \Delta \psi(t, \mathbf{x}))^2 d\Omega$$

► Also of interest — Kinetic Energy

$$\mathcal{K}(t) \triangleq \int_{\Omega} \mathbf{u}(t, \mathbf{x})^2 d\Omega = \int_{\Omega} (\nabla \psi(t, \mathbf{x}))^2 d\Omega$$

► Poincaré's inequality

$$\mathcal{K} \leq (2\pi)^{-2} \mathcal{E} \leq (2\pi)^{-2} \mathcal{P}$$

- ▶ Estimates for the Rate of Growth of Palinstrophy

$$\frac{d\mathcal{P}(t)}{dt} = \int_{\Omega} J(\Delta\psi, \psi) \Delta^2\psi \, d\Omega - \nu \int_{\Omega} (\Delta^2\psi)^2 \, d\Omega \triangleq \mathcal{R}_{\mathcal{P}}(\psi)$$

$$\frac{d\mathcal{P}(t)}{dt} \leq -\left(\frac{\nu}{\mathcal{E}}\right) \mathcal{P}^2 + C_1 \left(\frac{\mathcal{E}}{\nu}\right) \mathcal{P} \quad (\text{Doering \& Lunasin, 2011})$$

$$\frac{d\mathcal{P}(t)}{dt} \leq \frac{C_2}{\nu} \mathcal{K}^{1/2} \mathcal{P}^{3/2} \quad (\text{Ayala, 2012})$$

- ▶ Using Poincaré's inequality (may not be sharp)

$$\frac{d\mathcal{P}(t)}{dt} \leq \frac{C}{\nu} \mathcal{P}^2,$$

- ▶ Bound on growth in finite time

$$\max_{t>0} \mathcal{P}(t) \leq \left[\mathcal{P}_0^{1/2} + \frac{C_2}{4\nu^2} \mathcal{K}_0^{1/2} \mathcal{E}_0 \right]^2 \quad (\text{Ayala, 2012})$$

- Maximum Growth of $\frac{d\mathcal{P}(t)}{dt}$ for fixed $\mathcal{E}_0 > 0, \mathcal{P}_0 > (2\pi)^2 \mathcal{E}_0$

$$\max_{\psi \in \mathcal{S}_{\mathcal{P}_0, \mathcal{E}_0}} \mathcal{R}_{\mathcal{P}_0}(\psi) \quad \text{where}$$

$$\mathcal{S}_{\mathcal{P}_0, \mathcal{E}_0} = \left\{ \psi \in H^4(\Omega) : \begin{array}{l} \frac{1}{2} \int_{\Omega} (\nabla \Delta \psi)^2 d\Omega = \mathcal{P}_0 \\ \frac{1}{2} \int_{\Omega} (\Delta \psi)^2 d\Omega = \mathcal{E}_0 \end{array} \right\}$$

- Maximum Growth of $\frac{d\mathcal{P}(t)}{dt}$ for fixed $\mathcal{K}_0 > 0, \mathcal{P}_0 > (2\pi)^4 \mathcal{K}_0$

$$\max_{\psi \in \mathcal{S}_{\mathcal{P}_0, \mathcal{K}_0}} \mathcal{R}_{\mathcal{P}_0}(\psi) \quad \text{where}$$

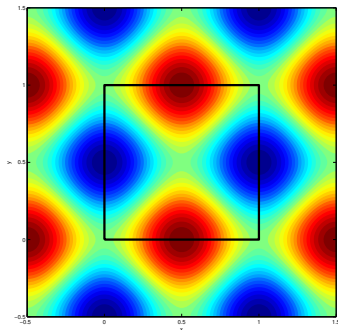
$$\mathcal{S}_{\mathcal{P}_0, \mathcal{K}_0} = \left\{ \psi \in H^4(\Omega) : \begin{array}{l} \frac{1}{2} \int_{\Omega} (\nabla \Delta \psi)^2 d\Omega = \mathcal{P}_0 \\ \frac{1}{2} \int_{\Omega} (\nabla \psi)^2 d\Omega = \mathcal{K}_0 \end{array} \right\}$$

- Small Palinstrophy Limit: $\mathcal{P}_0 \rightarrow (2\pi)^2 \mathcal{E}_0$

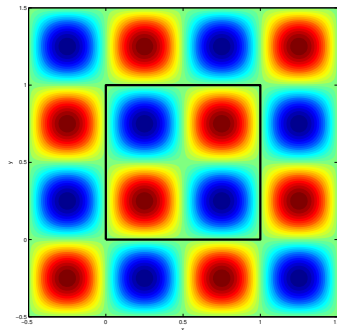
$$\tilde{\varphi}_0 = \arg \max_{\varphi \in \mathcal{S}_0} \mathcal{R}_0(\varphi), \quad \mathcal{R}_0(\varphi) = -\nu \int_{\Omega} (\Delta^2 \varphi)^2 d\Omega,$$

$$\mathcal{S}_0 = \left\{ \varphi \in H^4(\Omega) : \frac{1}{2} \int_{\Omega} (\nabla \Delta \psi)^2 d\Omega = \frac{(2\pi)^2}{2} \int_{\Omega} (\Delta \psi)^2 d\Omega \right\}$$

- Optimizers: **Eigenfunctions of the Laplacian** ($\tilde{\varphi}_0 \in \text{Ker}(\Delta)$)



$$\varphi(x, y) = \sin(\pi(y-x)) \sin(\pi(y+x))$$



$$\varphi(x, y) = \sin(2\pi x) \sin(2\pi y)$$

Numerical Solution of Maximization Problem

- Discretization of Gradient Flow

$$\begin{aligned}\frac{d\psi}{d\tau} &= -\nabla^{H^4} \mathcal{R}_\nu(\psi), & \psi(0) &= \psi_0 \\ \psi^{(n+1)} &= \psi^{(n)} - \Delta\tau^{(n)} \nabla^{H^4} \mathcal{R}_\nu(\psi^{(n)}), & \psi^{(0)} &= \psi_0\end{aligned}$$

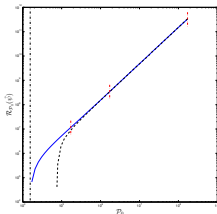
- Gradient in $H^4(\Omega)$ (via variational techniques)

$$\begin{aligned}[\text{Id} - L^8 \Delta^4] \nabla^{H^4} \mathcal{R}_\nu &= \nabla^{L^2} \mathcal{R}_\nu && (\text{Periodic BCs}) \\ \nabla^{L^2} \mathcal{R}_\nu(\psi) &= \Delta^2 J(\Delta\psi, \psi) + \Delta J(\psi, \Delta^2\psi) + J(\Delta^2\psi, \Delta\psi) - 2\nu\Delta^4\psi\end{aligned}$$

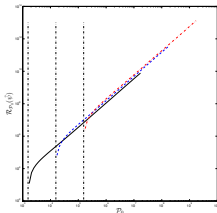
- Constraint satisfaction via arc minimization

Maximizers with Fixed $(\mathcal{K}_0, \mathcal{P}_0)$

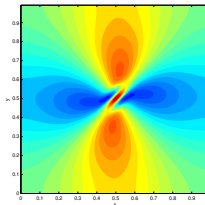
Estimate: $\frac{d\mathcal{P}(t)}{dt} \leq \frac{c_2}{\nu} \mathcal{K}_0^{1/2} \mathcal{P}_0^{3/2}$



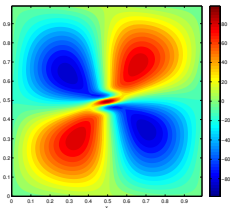
$\max \frac{d\mathcal{P}}{dt}$ vs. \mathcal{P}_0 , $\mathcal{K}_0 = 10$



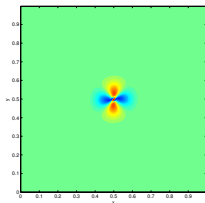
$\max \frac{d\mathcal{P}}{dt} \sim \mathcal{P}_0^{3/2}$ as $\mathcal{P}_0 \rightarrow \infty$



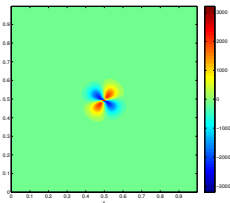
(a) $\mathcal{P}_0 \approx 10\mathcal{P}_c$



(b) $\mathcal{P}_0 \approx 10\mathcal{P}_c$



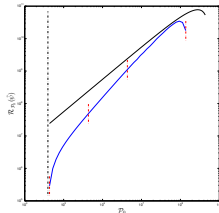
(c) $\mathcal{P}_0 \approx 10^4\mathcal{P}_c$



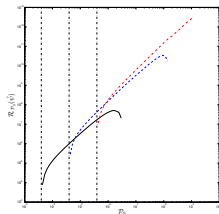
(d) $\mathcal{P}_0 \approx 10^4\mathcal{P}_c$

Maximizers with Fixed $(\mathcal{E}_0, \mathcal{P}_0)$

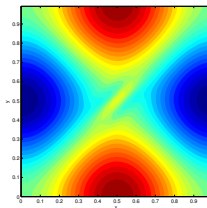
Estimate: $\frac{d\mathcal{P}(t)}{dt} \leq -(\frac{\nu}{\mathcal{E}_0}) \mathcal{P}_0^2 + C_1 (\frac{\mathcal{E}_0}{\nu}) \mathcal{P}_0$



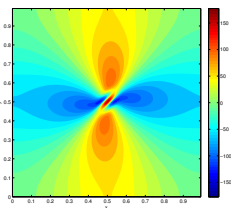
$\max \frac{d\mathcal{P}}{dt}$ vs. \mathcal{P}_0 , $\mathcal{E}_0 = 10^3$



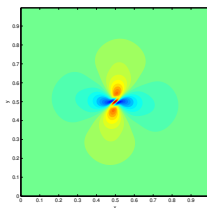
$\max \frac{d\mathcal{P}}{dt}$ vs. \mathcal{P}_0



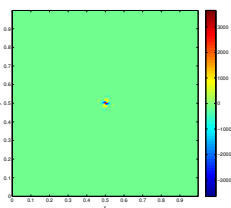
(a) $\mathcal{P}_0 \approx \mathcal{P}_c$



(b) $\mathcal{P}_0 \approx 10\mathcal{P}_c$



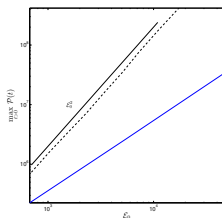
(c) $\mathcal{P}_0 \approx 10^2 \mathcal{P}_c$



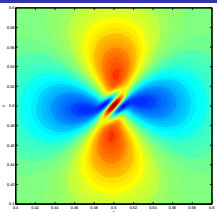
(d) $\mathcal{P}_0 \approx 10^{7/2} \mathcal{P}_c$

Maximizers with Fixed $(\mathcal{K}_0, \mathcal{P}_0)$

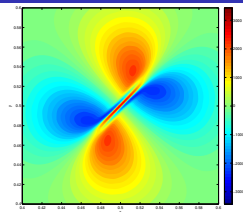
Finite-Time Estimate: $\max_{t>0} \mathcal{P}(t) \leq \left[\mathcal{P}_0^{1/2} + \frac{\mathcal{C}_2}{4\nu^2} \mathcal{K}_0^{1/2} \mathcal{E}_0 \right]^2$



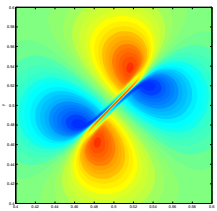
— \mathcal{P}_0 -constraint
 - - - $\{\mathcal{K}_0, \mathcal{P}_0\}$ -constraint



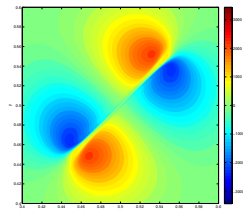
(a) $t = 0.000213$



(b) $t = 0.000458$



(c) $t = 0.000633$



(d) $t = 0.001265$

Relevant Estimates

	BEST ESTIMATE	SHARP?
1D Burgers instantaneous	$\frac{d\mathcal{E}(t)}{dt} \leq \frac{3}{2} \left(\frac{1}{\pi^2\nu}\right)^{1/3} \mathcal{E}(t)^{5/3}$	YES Lu & Doering (2008)
1D Burgers finite-time	$\max_{t \in [0, T]} \mathcal{E}(t) \leq \left[\mathcal{E}_0^{1/3} + \left(\frac{L}{4}\right)^2 \left(\frac{1}{\pi^2\nu}\right)^{4/3} \mathcal{E}_0 \right]^3$	No Ayala & P. (2011)
2D Navier-Stokes instantaneous	$\frac{d\mathcal{P}(t)}{dt} \leq -\left(\frac{\nu}{\mathcal{E}}\right) \mathcal{P}^2 + C_1 \left(\frac{\mathcal{E}}{\nu}\right) \mathcal{P}$ $\frac{d\mathcal{P}(t)}{dt} \leq \frac{C_2}{\nu} \mathcal{K}^{1/2} \mathcal{P}^{3/2}$	[YES] Ayala & P. (2013)
2D Navier-Stokes finite-time	$\max_{t>0} \mathcal{P}(t) \leq \left[\mathcal{P}_0^{1/2} + \frac{C_2}{4\nu^2} \mathcal{K}_0^{1/2} \mathcal{E}_0 \right]^2$	[YES] Ayala & P. (2013)
3D Navier-Stokes instantaneous	$\frac{d\mathcal{E}(t)}{dt} \leq \frac{27C^2}{128\nu^3} \mathcal{E}(t)^3$	YES Lu & Doering (2008)
3D Navier-Stokes finite-time	$\mathcal{E}(t) \leq \frac{\mathcal{E}(0)}{\sqrt{1 - 4\frac{C\mathcal{E}(0)^2}{\nu^3} t}}$???

Conclusions (II)

- ▶ Variational analysis and numerical optimization as tools for assessing sharpness of PDEs analysis
 - ▶ motivated by big open questions in mathematical fluid mechanics (including one of Clay “Millennium Problems”)
- ▶ 2D extreme vortex states saturate the worst-case bounds
 - ▶ analysis is sharp
- ▶ Ongoing work on the 3D case (the real problem)
 - ▶ so far, no evidence of blow-up, but results still from conclusive ...

References

- ▶ L. Lu and C. R. Doering, “Limits on Enstrophy Growth for Solutions of the Three-dimensional Navier-Stokes Equations” *Indiana University Mathematics Journal* **57**, 2693–2727, 2008.
-
- ▶ D. Ayala and B. Protas, “On Maximum Enstrophy Growth in a Hydrodynamic System”, *Physica D* **240**, 1553–1563, 2011.
 - ▶ D. Ayala and B. Protas, “Maximum Palinstrophy Growth in 2D Incompressible Flows: Instantaneous Case”, *Journal of Fluid Mechanics* **742** 340–367, 2014.
 - ▶ D. Ayala and B. Protas, “Vortices, Maximum Growth and the Problem of Finite-Time Singularity Formation”, *Fluid Dynamics Research (Special Issue for IUTAM Symposium on Vortex Dynamics)*, **46**, 031404, 2014.
 - ▶ D. Ayala and B. Protas, “Extreme Vortex States and the Growth of Enstrophy in 3D Incompressible Flows”, (submitted; see [arXiv:1605.05742](https://arxiv.org/abs/1605.05742)), 2016.