# Weakly asymmetric bridges and the KPZ equation

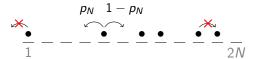
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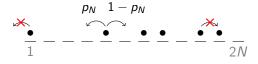
## The model

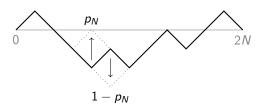
N particles on 2N sites



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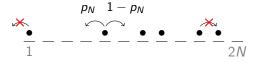
## N particles on 2N sites

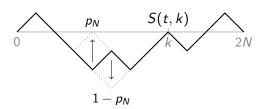




## The model

## N particles on 2N sites





## A simple fact

The unique invariant, reversible probability measure is

$$\mu_N(S) = \frac{1}{Z_N} \left( \frac{p_N}{1 - p_N} \right)^{\frac{1}{2}A(S)},$$

where we define the (signed) area under the interface S

$$A(S) = \sum_{k=1}^{2N} S(k) .$$

# Objective of this work

 $\rightarrow$  Understand the behaviour of the interface according to the asymmetry  $p_N-(1-p_N)$ .

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 $\rightarrow$  Understand the behaviour of the interface according to the asymmetry  $p_N - (1 - p_N)$ .

## Choice of parametrisation

$$rac{p_N}{1-p_N}=e^{rac{4\sigma}{(2N)^{lpha}}}\;,\qquad \sigma,lpha>0\;.$$

This means:

$$p_N = \frac{1}{2} + \frac{\sigma}{(2N)^{\alpha}} + \dots, \qquad 1 - p_N = \frac{1}{2} - \frac{\sigma}{(2N)^{\alpha}} + \dots.$$

# Scaling limit of the invariant measure

We present a Central Limit Theorem for the interface under  $\mu_N$ .

To that end, we rescale the interface in the following generic way:

$$u^{N}(x) := \frac{S(x N^{\cdots}) - \frac{\sum_{\alpha}^{N}(x)}{N^{\cdots}}}{N^{\cdots}},$$

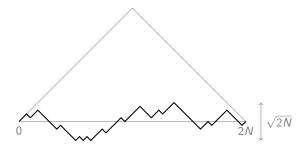
where  $\Sigma_{\alpha}^{N}$  is the mean under  $\mu_{N}$ .

# Scaling limit of the invariant measure: $\alpha > 3/2$

## Theorem (L.)

For all  $\alpha \in (3/2, \infty)$ ,

$$\left(\frac{S(x2N)}{\sqrt{2N}}, x \in [0,1]\right) \Rightarrow \text{ Brownian Bridge}$$

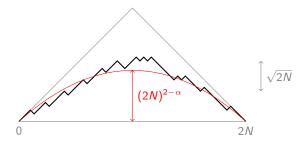


# Scaling limit of the invariant measure: $\alpha \in (1, 3/2]$

## Theorem (L.)

For all  $\alpha \in (1, 3/2]$ ,

$$\left( \frac{S(x2N) - \sigma x (1-x)(2N)^{2-\alpha}}{\sqrt{2N}} \right), x \in [0,1] \Rightarrow \text{ Brownian Bridge }.$$



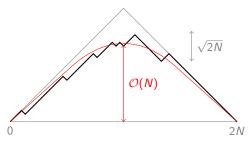
# Scaling limit of the invariant measure: $\alpha = 1$

#### Theorem (L.)

For all  $\alpha = 1$ ,

$$\left(\frac{S(x2N)-\sum_1^N(x)}{\sqrt{2N}}$$
,  $x\in[0,1]\right)\Rightarrow$  Time changed Brownian Bridge .

Sim. to Dobrushin-Hryniv (PTRF 96), Derrida-Enaud-Landim-Olla (JSP 05).



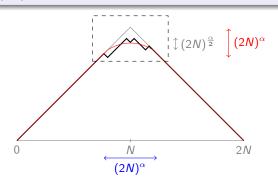
$$\Sigma_1^N(x) = 2N \int_0^x L'(\sigma(1-2y))dy + \mathcal{O}(1), \quad x \in [0,1].$$

# Scaling limit of the invariant measure: $\alpha < 1$

#### Theorem (L.)

For all  $\alpha < 1$ ,

$$\left(\frac{S(N+x(2N)^{\alpha})-\sum_{\alpha}^{N}(x)}{(2N)^{\frac{\alpha}{2}}}\text{ , }x\in\mathbf{R}\right)\Rightarrow\text{ Time changed Brownian Bridge }.$$



$$\sum_{\alpha}^{N}(x) = N + (2N)^{\alpha} \left( x + \int_{-x}^{\infty} (L'(2\sigma y) - 1) dy \right) + \mathcal{O}(1) , \quad x \in \mathbb{R} .$$

# Dynamic at equilibrium

What does the dynamical interface  $(S(t, k), t \ge 0, k \in [0, 2N])$  look like when it starts from  $\mu_N$  ?

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What does the dynamical interface  $(S(t, k), t \ge 0, k \in [0, 2N])$  look like when it starts from  $\mu_N$ ?

- Keep the same scalings in space and height as in the previous results.
- Speed up time appropriately.

## Dynamic at equilibrium: $\alpha > 1$

$$u^{N}(t,x) := \frac{S(t(2N)^{2},x2N) - \sigma(2N)^{2-\alpha}x(1-x)}{\sqrt{2N}}, \quad x \in [0,1].$$

#### Theorem (L.)

Start from  $\mu_N$ . Then,  $u^N \Rightarrow u$  where

$$\begin{cases} \partial_t u = \frac{1}{2} \partial_x^2 u + \xi, & x \in [0, 1] \\ u(t, 0) = u(t, 1) = 0, \end{cases}$$

and  $\xi$  space-time white noise.

# Dynamic at equilibrium: $\alpha = 1$

$$u^{N}(t,x) := \frac{S(t(2N)^{2}, x2N) - \frac{\sum_{1}^{N}(x)}{\sqrt{2N}}}{\sqrt{2N}}, \quad x \in [0,1].$$

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Start from  $\mu_N$ . Then,  $u^N \Rightarrow u$  where

$$\begin{cases} \partial_t u = \frac{1}{2} \partial_x^2 u - 2\sigma \, \partial_x \Sigma_1 \partial_x u + \sqrt{1 - (\partial_x \Sigma_1)^2} \, \xi \ , \quad x \in [0, 1] \\ u(t, 0) = u(t, 1) = 0 \ , \end{cases}$$

and  $\xi$  space-time white noise, and  $\Sigma_1(\cdot) = \lim_N \frac{\Sigma_1^N(\cdot)}{2N}$ .

Similar to De Masi-Presutti-Scacciatelli (Ann. IHP 89), Dittrich-Gärtner (MN 91).

# Dynamic at equilibrium: $\alpha < 1$

$$u^{N}(t,x):=\frac{S(t(2N)^{2\alpha},N+x(2N)^{\alpha})-\frac{\sum_{\alpha}^{N}(x)}{(2N)^{\frac{\alpha}{2}}},\quad x\in\mathbb{R}.$$

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Start from  $\mu_N$ . Then,  $u^N \Rightarrow u$  where

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.

We look at:

- the hydrodynamic limit,
- the fluctuations.

## Hydrodynamic limit

Set

$$m^{N}(t,x) = \begin{cases} \frac{1}{(2N)^{2-\alpha}} S(t(2N)^{2}, x(2N)), & \alpha \in [1,3/2), \\ \frac{1}{2N} S(t(2N)^{1+\alpha}, x(2N)), & \alpha < 1. \end{cases}$$

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#### Theorem (L.)

Start from a flat initial condition. Then  $m^N \Rightarrow m$  where

$$\partial_t m = \begin{cases} \frac{1}{2} \partial_x^2 m + \sigma , & \alpha \in (1, 3/2) , \\ \frac{1}{2} \partial_x^2 m + \sigma (1 - (\partial_x m)^2) , & \alpha = 1 , \\ \sigma (1 - (\partial_x m)^2) , & \alpha < 1 , \end{cases}$$

with Dirichlet boundary conditions m(t, 0) = m(t, 1) = 0.

Similar results: Gärtner (SPA 88), De Masi-Presutti-Scacciatelli (Ann. IHP 89), Kipnis-Olla-Varadhan (CPAM 89), Rezakhanlou (CMP 91), Bahadoran (CMP 12) ...

## More on the case $\alpha < 1$

$$egin{cases} \partial_t m = \sigma ig(1 - (\partial_x m)^2ig) \ , \ m(t=0,\cdot) \equiv 0 \ , \ m(t,0) = m(t,1) = 0 \ . \end{cases}$$

#### More on the case $\alpha < 1$

$$\begin{cases} \partial_t m = \sigma (1 - (\partial_x m)^2), \\ m(t = 0, \cdot) \equiv 0, \\ m(t, 0) = m(t, 1) = 0. \end{cases}$$

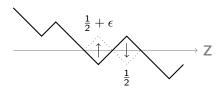
#### **Inviscid Burgers**

We actually prove convergence of the density of particles towards the (entropy!) solution of the inviscid Burgers equation:

$$\begin{cases} \partial_t \varrho = \sigma \partial_x \big( \varrho (1 - \varrho) \big) , \\ \varrho (t = 0, \cdot) = 1/2 , \\ \varrho (t, 0) = 1 , \quad \varrho (t, 1) = 0 . \end{cases}$$

Solution theory by Bardos-Le Roux-Nédélec (CPDE 79). CV result similar to Rezakhanlou (CMP 91), Bahadoran (CMP 12).

## A famous result of Bertini and Giacomin on KPZ



# Theorem (Bertini-Giacomin CMP 97)

Set 
$$h^{\epsilon}(t,x) = \epsilon \left(S\left(\frac{t}{\epsilon^4},\frac{x}{\epsilon^2}\right) - \frac{t}{\epsilon^3}\right)$$
. Then,  $h^{\epsilon} \Rightarrow h$  where

(KPZ) 
$$\partial_t h = \frac{1}{2} \partial_x^2 h - \frac{1}{2} (\partial_x h)^2 + \xi$$
,  $x \in \mathbf{R}$ .

Let's apply Bertini-Giacomin's scaling in our case:

• Height scaling:  $\epsilon \leftrightarrow \frac{1}{(2N)^{\alpha}}$ . Heights are smaller than N. So we need to take  $\alpha \leq 1$ .

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- Space scaling:  $\epsilon^{-2} \leftrightarrow (2N)^{2\alpha}$ . Lattice of order N. So we need to take  $\alpha \le 1/2$ .
- Time scaling:  $\epsilon^{-4} \leftrightarrow (2N)^{4\alpha}$ . Hydro. limit reaches equilibrium in finite time in the time scale  $(2N)^{1+\alpha}$ . So we need to take  $4\alpha \le 1+\alpha$ , or equivalently  $\alpha \le 1/3$ .

For  $\alpha \leq 1/3$ , set

$$h^{N}(t,x) = \frac{1}{(2N)^{\alpha}} \left( S\left(t(2N)^{4\alpha}, N + x(2N)^{2\alpha}\right) - \sigma t(2N)^{3\alpha} \right).$$

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## Theorem (L.)

We have  $h^N \Rightarrow h$  in  $\mathbb{D}([0, T), \mathcal{C}(R))$  where h is the solution of (KPZ) on the line and

$$T = egin{cases} +\infty & & lpha < 1/3 \ rac{1}{2\sigma} & & lpha = 1/3 \ . \end{cases}$$

Notice that the fluctuations vanish *suddenly* at time T in the case  $\alpha = 1/3$ .

