On the stabilization of the incompressible Navier-Stokes equations in a 2d channel with a normal control

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Outline









Outline



2 Strategy





Incompressible Navier-Stokes equations in a 2-d channel:

(0, 1)

$$\begin{split} &\Omega = \mathbb{I} \times (0,1), \text{where } \mathbb{I} = \mathbb{R}/2\pi\mathbb{Z}. \\ &\left\{ \begin{array}{ll} \partial_t u + (u \cdot \nabla)u - \nu\Delta u + \nabla p = 0, & \text{in } (0,\infty) \times \Omega, \\ \text{div } u = 0, & \text{in } (0,\infty) \times \Omega, \\ u(t,x_1,0) = (0,0), & \text{on } (0,\infty) \times \mathbb{T}, \\ u(t,x_1,1) = (0, \textbf{v}(t,x_1)), & \text{on } (0,\infty) \times \mathbb{T}, \\ u(0,x_1,x_2) = u^0(x_1,x_2), & \text{in } \Omega. \end{split} \right. \end{split}$$

PTTN

 m/α

•
$$u = u(t, x_1, x_2) \in \mathbb{R}^2$$
 is the velocity.

•
$$p = p(t, x_1, x_2)$$
 is the pressure.

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•
$$\nu > 0$$
 is the viscosity.

• $v = v(t, x_1)$ is the control function, acting on the normal component only.

Choose v to stabilize the state u.

Setting Result

Motivation and related topics

Motivation: Controllability/Stabilization of fluid-structure models with controls acting on the structure. See Lions Zuazua '95, Osses Puel '99, '09, Lequeurre '13, ...

Related topics:

- Controllability of incompressible Navier-Stokes equations.... Fursikov Imanuvilov '96, Fernandez-Cara Guerrero Imanuvilov Puel '04, ...
- ... with controls having zero components: Coron Guerrero '09, Carreno Guerrero '13, Coron Lissy '15,...
- Coupled parabolic systems with one boundary control: Ammar-Khodja Benabdallah Gonzalez-Burgos de Teresa '11, Duprez Lissy '15...
- Stabilization for incompressible Navier-Stokes equations: Krstic et al '01, Raymond '06, Barbu '07, Triggiani '07, Vazquez Coron Trélat '08, Munteanu '12,...

To be more precise....

Our goal

Get a local stabilization result around the state (u, p) = (0, 0).

Linearized equations:

$$\begin{cases} \partial_t u - \nu \Delta u + \nabla p = 0, & \text{ in } (0, \infty) \times \Omega, \\ \text{div } u = 0, & \text{ in } (0, \infty) \times \Omega, \\ u(t, x_1, 0) = (0, 0), & \text{ on } (0, \infty) \times \mathbb{T}, \\ u(t, x_1, 1) = (0, v(t, x_1)), & \text{ on } (0, \infty) \times \mathbb{T}, \\ u(0, x_1, x_2) = u^0(x_1, x_2), & \text{ in } \Omega, \end{cases} \end{cases}$$

??? \rightsquigarrow The linearized equations are already stable! Taking v = 0,

$$\frac{d}{dt}\left(\frac{1}{2}\int_{\Omega}|u(t,x)|^{2}\,dx\right)+\nu\int_{\Omega}|\nabla u(t,x)|^{2}\,dx=0$$

 \rightsquigarrow Exponential decay like $t \mapsto \exp(-\nu \pi^2 t)$.

To be more precise....

Get a local stabilization result around the state (u, p) = (0, 0)At an exponential rate larger than $\nu \pi^2$.

Linearized equations:

$$\left\{ \begin{array}{ll} \partial_t u - \nu \Delta u + \nabla p = 0, & \text{ in } (0, \infty) \times \Omega, \\ \text{div } u = 0, & \text{ in } (0, \infty) \times \Omega, \\ u(t, x_1, 0) = (0, 0), & \text{ on } (0, \infty) \times \mathbb{T}, \\ u(t, x_1, 1) = (0, v(t, x_1)), & \text{ on } (0, \infty) \times \mathbb{T}, \\ u(0, x_1, x_2) = u^0(x_1, x_2), & \text{ in } \Omega, \end{array} \right.$$

Difficulty:

div
$$u = 0$$
 in $(0,\infty) \times \Omega \Rightarrow \int_{\mathbb{T}} v(t,x_1) dx_1 = 0$ for all $t > 0$.

To be more precise....

The 0-mode:

$$u_0(t,x_2) = \int_{\mathbb{T}} u(t,x_1,x_2) \, dx_1$$

satisfies the uncontrolled heat equation

$$\begin{cases} \partial_t u_{0,1} - \nu \partial_{22} u_{0,1} = 0, & \text{ in } (0,\infty) \times (0,1), \\ u_{0,1}(t,0) = u_{0,1}(t,1) = 0, & \text{ on } (0,\infty), \\ u_{0,2}(t,x_2) = 0, & \text{ in } (0,\infty) \times (0,1). \end{cases}$$

Consequence

The solutions of the linearized equations decay like $\exp(-\nu\pi^2 t)$ and, considering

$$u(t,x) = e^{-\nu \pi^2 t} \Psi_0(x_2)$$
 with $\Psi_0 = \Psi_0(x_2) = \sqrt{\frac{2}{\pi}} \begin{pmatrix} \sin(\pi x_2) \\ 0 \end{pmatrix}$,

this decay estimate is sharp whatever the control v is.

Main result

Theorem (S. Chowdhury, S.E., J.-P. Raymond)

Let $\omega_0 > 0$ be such that $0 < \omega_0 < 4\nu\pi^2$. There exists $\gamma > 0$ such that for all $u_0 \in \mathbf{V}_0^1(\Omega)$ with $\|u_0\|_{\mathbf{V}_0^1(\Omega)} \le \gamma$, there exists $v \in L^2((0,\infty) \times \mathbb{T})$ satisfying $\int_{\mathbb{T}} v(t, x_1) dx_1 = 0$ for all t > 0 such that the solution (u, p) of the incompressible Navier-Stokes equation satisfies, for some constant C > 0 independent of t,

$$orall t \geq 0, \quad \|u(t)\|_{\mathbf{V}^1(\Omega)} \leq C e^{-\omega_0 t}.$$

$$\begin{aligned} \mathbf{V}^{1}(\Omega) &= \left\{ u = (u_{1}, u_{2}) \in H^{1}(\Omega) \times H^{1}(\Omega) \mid \text{div } u = 0 \right\}, \\ \mathbf{V}^{1}_{0}(\Omega) &= \left\{ u \in \mathbf{V}^{1}(\Omega) \mid u(x_{1}, 0) = u(x_{1}, 1) = 0 \text{ for } x_{1} \in \mathbb{T} \right\}. \end{aligned}$$

Comments

- Straightforward when $\omega < \nu \pi^2$ \rightsquigarrow Interesting case $\omega \in (\nu \pi^2, 4\nu \pi^2)$.
- $4\nu\pi^2$ is the second eigenvalue of the elliptic operator generating the heat equation satisfied by the 0-mode:

$$\begin{cases} \partial_t u_{0,1} - \nu \partial_{22} u_{0,1} = 0, & \text{ in } (0,\infty) \times (0,1), \\ u_{0,1}(t,0) = u_{0,1}(t,1) = 0, & \text{ on } (0,\infty), \\ u_{0,2}(t,x_2) = 0, & \text{ in } (0,\infty) \times (0,1). \end{cases}$$

The stabilization result cannot be true for the linearized model
 ⇒ We have to use the non-linearity to improve the
 exponential decay.
 Strategy based on the so-called Power Series Expansion:
 see Coron Crépeau '04, Cerpa '07, Cerpa Crépeau '09, Coron
 Rivas '15.

Outline

Introduction







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Strategy

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Write
$$u = \varepsilon \alpha + \varepsilon^2 \beta$$
, $v = \varepsilon v_1 + \varepsilon^2 v_2$, with

$$\begin{cases} \partial_t \alpha - \nu \Delta \alpha + \nabla p_1 = 0, & \text{ in } (0, \infty) \times \Omega, \\ \operatorname{div} \alpha = 0, & \operatorname{in } (0, \infty) \times \Omega, \\ \alpha(t, x_1, 0) = (0, 0), & \text{ on } (0, \infty) \times \mathbb{T}, \\ \alpha(t, x_1, 1) = (0, \nu_1(t, x_1)), & \text{ on } (0, \infty) \times \mathbb{T}, \\ \alpha(0, x_1, x_2) = \alpha^0(x_1, x_2), & \text{ in } \Omega, \end{cases}$$

$$\begin{array}{ll} \left(\begin{array}{ll} \partial_t\beta - \nu\Delta\beta + \nabla p_2 = -(\alpha + \varepsilon\beta) \cdot \nabla(\alpha + \varepsilon\beta), & \text{ in } (0, \infty) \times \Omega, \\ \text{div } \beta = 0, & \text{ in } (0, \infty) \times \Omega, \\ \beta(t, x_1, 0) = (0, 0), & \text{ on } (0, \infty) \times \mathbb{T}, \\ \beta(t, x_1, 1) = (0, \nu_2(t, x_1)), & \text{ on } (0, \infty) \times \mathbb{T}, \\ \beta(0, x_1, x_2) = \beta^0(x_1, x_2), & \text{ in } \Omega, \end{array} \right)$$

Strategy

Part 2

- α satisfies the linearized incompressible Navier-Stokes equations.
- ⇒ If α contains 0-modes decaying slower than exp $(-\omega_0 t)$, one cannot achieve an exponential decay rate ω_0 .
- ⇒ The component of the solution *u* on the eigenfunction $\Psi_0 = \Psi_0(x_2) = \sqrt{\frac{2}{\pi}} \begin{pmatrix} \sin(\pi x_2) \\ 0 \end{pmatrix}$ • Is in β .
 - Should be handled by constructing a suitable α .

Preliminaries

• The Stokes operator A is self-adjoint, positive definite, with compact resolvent on the space $\mathbf{V}_n^0(\Omega) = \left\{ u \in (L^2(\Omega))^2 \mid \text{div}(u) = 0 \text{ on } \Omega \text{ and } u \cdot n = 0 \text{ on } \Gamma \right\}$ \Rightarrow Sequences of positive eigenvalues $\lambda_1 \leq \lambda_2 \leq \cdots \rightarrow \infty$ and corresponding orthonormal basis of eigenvectors (Ψ_j) .

$$A\Psi = \lambda \Psi \Leftrightarrow \begin{cases} -\nu \Delta \Psi + \nabla q = \lambda \Psi, & \text{in } \Omega, \\ \text{div } \Psi = 0, & \text{in } \Omega, \\ \Psi = 0, & \text{on } \Gamma, \end{cases}$$

Adjoint of the control operator: $B^*\Psi = q(x_1, 1) - \frac{1}{2\pi} \int_{\mathbb{T}} q(x_1, 1) dx_1.$

Lemma $A\Psi = \lambda \Psi$ and $B^*\Psi = 0$ imply $\Psi(x) = \Psi(x_2)$.

Decomposition of the space $\mathbf{V}_{n}^{0}(\Omega)$:

- A stable space: $\mathbf{Z}_s = \text{Span} \{ \Phi \mid A\Phi = \lambda \Phi, \text{ with } \lambda > \omega \}.$
- An unstable space: $\mathbf{Z}_u = \mathbf{Z}_s^{\perp}$, itself decomposed as
 - An unstable uncontrollable space $\mathbf{Z}_{uu} = \operatorname{Span} \Psi_0$.
 - An unstable detectable space $\mathbf{Z}_{ud} = \mathbf{Z}_{uu}^{\perp} \cap \mathbf{Z}_{u}$.

Corresponding (orthogonal) projections: \mathbb{P}_s , \mathbb{P}_u , \mathbb{P}_{ud} and \mathbb{P}_{uu} .

Roughly Notations Precise Lemmata

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Part 3

Strategy

Iterative strategy: $(0, \infty) = \bigcup_{n \in \mathbb{N}} [nT, (n+1)T]$ for some T > 0. (T = 1).

Starting point: $u^0 = \varepsilon \alpha^0 + \varepsilon^2 \beta^0$ with

$$\left\|\alpha^0\right\|_{\mathbf{V}_0^1(\Omega)}^2+\left\|\beta^0\right\|_{\mathbf{V}_0^1(\Omega)}\leq 1,\quad \text{ and }\quad \mathbb{P}_{uu}\alpha^0=0.$$

Initialization Step: On [0, T], choose

- v_1 such that $\mathbb{P}_u \alpha(T) = 0$.
- $v_2 = 0$.

Part 4

Strategy

Iteration step: In each time interval [nT, (n+1)T], we design controls v_1 and v_2 such that

 $\mathbb{P}_{u}\alpha((n+1)T) = 0$, and $\mathbb{P}_{u}\beta^{*}((n+1)T) = 0$,

where β^{\ast} is the solution of

$$\begin{cases} \partial_t \beta^* - \nu \Delta \beta^* + \nabla p^* = -\alpha \cdot \nabla \alpha, & \text{in } (nT, (n+1)T) \times \Omega, \\ \text{div } \beta^* = 0, & \text{in } (nT, (n+1)T) \times \Omega, \\ \beta^*(t, x_1, 0) = (0, 0), & \text{on } (nT, (n+1)T) \times \mathbb{T}, \\ \beta^*(t, x_1, 1) = (0, \nu_2(t, x_1)), & \text{on } (nT, (n+1)T) \times \mathbb{T}, \\ \beta^*(nT^+, x) = \beta(nT^-, x), & \text{in } \Omega. \end{cases}$$

Remark: $\beta^* \neq \beta$, but $\|\beta - \beta^*\| \leq \varepsilon$

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Lemma for the initialization step

Given $\alpha^0 \in \mathbf{V}_0^1(\Omega)$ and $\alpha^f \in \mathbf{Z}_{ud}$, there exists a control function

$$v_1 \in H^1_0(0, T; L^2(\mathbb{T})) \cap L^2(0, T; H^2(\mathbb{T})) \text{ with } \int_{\mathbb{T}} v_1(t, x_1) \, dx_1 = 0$$

such that the solution α of

$$\begin{cases} \begin{array}{ll} \partial_t \alpha - \nu \Delta \alpha + \nabla p_1 = 0, & \text{ in } (0, \infty) \times \Omega, \\ \text{div } \alpha = 0, & \text{ in } (0, \infty) \times \Omega, \\ \alpha(t, x_1, 0) = (0, 0), & \text{ on } (0, \infty) \times \mathbb{T}, \\ \alpha(t, x_1, 1) = (0, \nu_1(t, x_1)), & \text{ on } (0, \infty) \times \mathbb{T}, \\ \alpha(0, x_1, x_2) = \alpha^0(x_1, x_2), & \text{ in } \Omega, \end{cases} \end{cases}$$

satisfies $\mathbb{P}_{ud}\alpha(T) = \alpha^f$. The projection $\mathbb{P}_{uu}\alpha(T)$ cannot be controlled, but satisfies $\mathbb{P}_{uu}\alpha(T) = \exp(-\nu\pi^2 T)\mathbb{P}_{uu}\alpha^0$.

Key Lemma for the iteration process

Let $\tilde{\beta}^0 \in \mathbf{Z}_{uu}$ and $f \in L^2(0, T; \mathbf{H}^2(\Omega)) \cap H^1(0, T; \mathbf{V}^0(\Omega))$. Then there exists a control function $\tilde{v}_1 \in H^1_0(0, T; H^2(\mathbb{T}) \cap L^2_0(\mathbb{T}))$ such that the solution $\tilde{\alpha}$ of

$$\begin{cases} \partial_t \tilde{\alpha} - \nu \Delta \tilde{\alpha} + \nabla \tilde{p}_1 = 0, & \text{in } (0, T) \times \Omega, \\ \text{div } \tilde{\alpha} = 0, & \text{in } (0, T) \times \Omega, \\ \tilde{\alpha}(t, x_1, 0) = (0, 0), & \text{on } (0, T) \times \mathbb{T}, \\ \tilde{\alpha}(t, x_1, 1) = (0, \tilde{\nu}_1(t, x_1)), & \text{on } (0, T) \times \mathbb{T}, \\ \tilde{\alpha}(0, x) = 0, & \text{in } \Omega, \end{cases}$$

satisfies $\tilde{\alpha}(T) = 0$ in Ω , and such that the solution $\tilde{\beta}$ of

$$\begin{cases} \partial_t \tilde{\beta} - \nu \Delta \tilde{\beta} + \nabla \tilde{p}_2 = -(f + \tilde{\alpha}) \cdot \nabla (f + \tilde{\alpha}), & \text{ in } (0, T) \times \Omega, \\ \text{div } \tilde{\beta} = 0, & \text{ in } (0, T) \times \Omega, \\ \tilde{\beta}(t, x_1, 0) = \tilde{\beta}(t, x_1, 1) = (0, 0), & \text{ on } (0, \infty) \times \mathbb{T}, \\ \tilde{\beta}(0, x_1, x_2) = \tilde{\beta}^0(x_1, x_2), & \text{ in } \Omega, \end{cases}$$

satisfies $\mathbb{P}_{uu}\tilde{\beta}(T) = 0$.

 \rightsquigarrow These lemmata come with estimates which are needed to conclude. For instance, we get

$$\begin{split} \|\tilde{\alpha}\|_{L^{2}(0,T;\mathbf{H}^{2}(\Omega))\cap H^{1}(0,T;\mathbf{L}^{2}(\Omega))}^{2} + \|\tilde{\beta}\|_{L^{2}(0,T;\mathbf{H}^{2}(\Omega))\cap H^{1}(0,T;\mathbf{L}^{2}(\Omega))} \\ & \leq C\left(\|f\|_{L^{2}(0,T;\mathbf{H}^{2}(\Omega))\cap H^{1}(0,T;\mathbf{V}^{0}(\Omega))}^{2} + \|\tilde{\beta}^{0}\|_{\mathbf{V}_{0}^{1}(\Omega)}\right) \end{split}$$

After tedious estimates, at each step,

$$\begin{aligned} \|\alpha((n+1)T)\|_{\mathbf{V}_{0}^{1}(\Omega)} &\leq e^{-\omega T} \|\alpha(nT)\|_{\mathbf{V}_{0}^{1}(\Omega)}, \\ \left(\|\mathbb{P}_{s}\beta((n+1)T)\|_{\mathbf{V}_{0}^{1}(\Omega)}\right) &\leq K_{\varepsilon} \left(\|\mathbb{P}_{s}\beta(nT)\|_{\mathbf{V}_{0}^{1}(\Omega)}\right) + \|\alpha(nT)\|_{\mathbf{V}_{0}^{1}(\Omega)}^{2} \left(\begin{array}{c} C\\ C \end{array}\right) \\ \|\mathbb{P}_{u}\beta((n+1)T)\|_{\mathbf{V}_{0}^{1}(\Omega)} \end{array}\right) &\leq K_{\varepsilon} \left(\begin{array}{c} \|\mathbb{P}_{s}\beta(nT)\|_{\mathbf{V}_{0}^{1}(\Omega)}\\ \|\mathbb{P}_{u}\beta(nT)\|_{\mathbf{V}_{0}^{1}(\Omega)} \end{array}\right) + \|\alpha(nT)\|_{\mathbf{V}_{0}^{1}(\Omega)}^{2} \left(\begin{array}{c} C\\ C \end{array}\right) \\ \text{where } K_{\varepsilon} &= \left(\begin{array}{c} e^{-\omega T} + C\sqrt{\varepsilon} & C\\ C\sqrt{\varepsilon} & C\sqrt{\varepsilon} \end{array}\right). \end{aligned}$$

 \Rightarrow Decay as exp $(-\omega_0 t)$.

Outline

1 Introduction







Key Lemma

Let $\tilde{\beta}^0 \in \mathbf{Z}_{uu}$ and $f \in L^2(0, T; \mathbf{H}^2(\Omega)) \cap H^1(0, T; \mathbf{V}^0(\Omega))$. Then there exists a control function $\tilde{v}_1 \in H^1_0(0, T; H^2(\mathbb{T}) \cap L^2_0(\mathbb{T}))$ such that the solution $\tilde{\alpha}$ of

$$\begin{cases} \partial_t \tilde{\alpha} - \nu \Delta \tilde{\alpha} + \nabla \tilde{p}_1 = 0, & \text{in } (0, T) \times \Omega, \\ \text{div } \tilde{\alpha} = 0, & \text{in } (0, T) \times \Omega, \\ \tilde{\alpha}(t, x_1, 0) = (0, 0), & \text{on } (0, T) \times \mathbb{T}, \\ \tilde{\alpha}(t, x_1, 1) = (0, \tilde{\nu}_1(t, x_1)), & \text{on } (0, T) \times \mathbb{T}, \\ \tilde{\alpha}(0, x) = 0, & \text{in } \Omega, \end{cases}$$

satisfies $\tilde{\alpha}(T) = 0$ in Ω , and such that the solution $\tilde{\beta}$ of

$$\begin{cases} \partial_t \tilde{\beta} - \nu \Delta \tilde{\beta} + \nabla \tilde{p}_2 = -(f + \tilde{\alpha}) \cdot \nabla (f + \tilde{\alpha}), & \text{ in } (0, T) \times \Omega, \\ \text{div } \tilde{\beta} = 0, & \text{ in } (0, T) \times \Omega, \\ \tilde{\beta}(t, x_1, 0) = \tilde{\beta}(t, x_1, 1) = (0, 0), & \text{ on } (0, \infty) \times \mathbb{T}, \\ \tilde{\beta}(0, x_1, x_2) = \tilde{\beta}^0(x_1, x_2), & \text{ in } \Omega, \end{cases}$$

satisfies $\mathbb{P}_{uu}\tilde{\beta}(T) = 0$.

A slightly stronger Result

There exist control functions v^a , $v^b \in H^1_0(0, T; H^2(\mathbb{T}) \cap L^2_0(\mathbb{T}))$ such that, for all $a, b \in \mathbb{R}$, the solution α of

$$\begin{array}{ll} & \left(\begin{array}{ll} \partial_t \alpha - \nu \Delta \alpha + \nabla p_1 = 0, & \text{in } (0, T) \times \Omega, \\ \text{div } \alpha = 0, & \text{in } (0, T) \times \Omega, \\ \alpha(t, x_1, 0) = (0, 0), & \text{on } (0, T) \times \mathbb{T} \\ \alpha(t, x_1, 1) = (0, (av^a + bv^b)(t, x_1)), & \text{on } (0, T) \times \mathbb{T} \\ \alpha(0, x) = 0, & \text{in } \Omega, \end{array} \right.$$

satisfies $\alpha(T) = 0$ in Ω , and such that the solution β of

$$\begin{cases} \partial_t \beta - \nu \Delta \beta + \nabla p_2 = -\alpha \cdot \nabla \alpha, & \text{in } (0, T) \times \Omega, \\ \text{div } \beta = 0, & \text{in } (0, T) \times \Omega, \\ \beta(t, x_1, 0) = \beta(t, x_1, 1) = (0, 0), & \text{on } (0, \infty) \times \mathbb{T}, \\ \beta(0, x) = 0, & \text{in } \Omega, \end{cases}$$

satisfies $\mathbb{P}_{uu}\beta(T) = ab\Psi_0$.

Ideas of the proof

Take

$$v^{a}(t, x_{1}) = v^{c}(t)\cos(x_{1}), \quad v^{b}(t, x_{1}) = v^{s}(t)\sin(x_{1}).$$

 $\rightsquigarrow \alpha^{a}$ and α^{b} are supported on the first mode of the equations:

$$\alpha^{a}(t, x_1, x_2) = \begin{pmatrix} \sin(x_1)\alpha_1^{c}(t, x_2) \\ \cos(x_1)\alpha_2^{c}(t, x_2) \end{pmatrix}$$

• The equation satisfied by the first modes of the linear incompressible Stokes equations is null-controllable. ~> Proof by spectral estimates and Müntz Theorem.

• Writing $\langle \beta(T), \Psi_0 \rangle$ in a suitable way.

$$e^{
u\pi^2 T}\langle eta(T),\Psi_0
angle=\pi^{5/2}\int_0^T v^s(t)q(t,1)\,dt,$$

where q is obtained by solving

$$\begin{cases} -\partial_t Z + \nu Z - \nu \partial_{22} Z + \begin{pmatrix} q \\ \partial_2 q \end{pmatrix} = F(t, x_2), & \text{ in } (0, T) \times (0, 1), \\ -Z_1 + \partial_2 Z_2 = 0, & \text{ in } (0, T) \times (0, 1), \\ Z(t, 0) = Z(t, 1) = (0, 0), & \text{ in } (0, T), \\ Z(T, x_2) = 0, & \text{ in } (0, 1). \end{cases}$$

with
$$F(t, x_2) = \cos(\pi x_2)e^{\nu\pi^2 t} \begin{pmatrix} \alpha_2^c(t, x_2) \\ \alpha_1^s(t, x_2) \end{pmatrix}$$
, depending only on $v^a(t, x_1) = v^c(t)\cos(x_1)$.

 \rightsquigarrow Show the existence of v^a/v^c such that $\|q(t,1)\|_{L^2(0,T)} \neq 0$.

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Construction of v^c , generation of a suitable trajectory

For $\mu \in \mathbb{R}$, introduce $(\alpha^*(x_2), p^*(x_2))$ solving

$$\begin{array}{ll} & \mu \alpha_1^* + \nu \alpha_1^* - \nu \partial_{22} \alpha_1^* - p^* = 0, & \text{ in } (0,1), \\ & \mu \alpha_2^* + \nu \alpha_2^* - \nu \partial_{22} \alpha_2^* + \partial_2 p^* = 0, & \text{ in } (0,1), \\ & \alpha_1^* + \partial_2 \alpha_2^* = 0, & \text{ in } (0,1), \\ & \alpha_1^*(0) = \alpha_1^*(1) = \alpha_2^*(0) = 0, & \alpha_2^*(1) = 1. \end{array}$$

Then $\overline{\alpha}(t, x_1, x_2) = e^{\mu t} (\sin(x_1)\alpha_1^*(x_2), \cos(x_1)\alpha_2^*(x_2)), \ \overline{\nu}(t) = e^{\mu t},$ solves the linear Stokes equations.

Lemma

There exists a suitable $\mu \in \mathbb{R}$ such that if $\alpha(t) = \overline{\alpha}(t)$ on some time interval then the boundary pressure q(t, 1) given by the aforementioned process cannot be identically 0 on that time interval.

Reduction to the stationary case and numerically checked.

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Construction of v^a/v^c and v^b/v^s

Construction of v^a/v^c in 4 steps:

- On (0, T/4), control α^a to go from 0 to $\overline{\alpha}(T/4)$.
- On (T/4, T/2), take $v^a(t) = e^{\mu t}$ and $\alpha^a(t) = \overline{\alpha}(t)$.
- On (T/2, 3T/4), control α^a goes from $\overline{\alpha}(T/2)$ to 0.
- On (3T/4, T), take $v^a(t) = 0$, and $\alpha^a(t) = 0$, hence q(t) = 0.

Construction of v^b/v^s :

On (0, 3T/4), take v^s such that \$\int_{0}^{3T/4} v^{s}(t)q(t,1) dt = 1\$.
On (3T/4, T), control α^b to go from α^b(3T/4) to 0.

Outline

Introduction







Open Question

• Exponential stabilization result at a rate higher than $4\nu\pi^2$?

Difficulty: One has to guarantee that we can enter the space of missing directions

Span
$$\left\{ \begin{pmatrix} \sin(\pi x_2) \\ 0 \end{pmatrix}, \begin{pmatrix} \sin(2\pi x_2) \\ 0 \end{pmatrix} \right\}$$

in both directions independently.

Exponential stabilization at any given rate is open, so is the controllability of the system.

Thank you for your attention!

Comments Welcome

Reference:

Open loop stabilization of incompressible Navier-Stokes equations in a 2d channel using power series expansion.

S. Chowdhury, S. Ervedoza, and J.-P. Raymond, in preparation.