

# On some problems and techniques in spectral geometry

Jean-Marc Bouclet  
Institut de Mathématiques de Toulouse

January 8, 2016

## Outline of the talk

1. Background on Riemannian geometry
2. Examples of problems in spectral geometry
3. Basic strategy to handle these questions
4. Relationship with the geodesic flow
5. An idea of proof

## Background on Riemannian geometry

Let  $(M^n, g)$  be a Riemannian manifold.

- ▶ Laplacian:

## Background on Riemannian geometry

Let  $(M^n, g)$  be a Riemannian manifold.

- ▶ Laplacian: writing  $g = \sum_{j,k} g_{jk}(x) dx_j dx_k$

## Background on Riemannian geometry

Let  $(M^n, g)$  be a Riemannian manifold.

- ▶ Laplacian: writing  $g = \sum_{j,k} g_{jk}(x) dx_j dx_k$  and  $(g^{jk}(x)) = (g_{jk}(x))^{-1}$

## Background on Riemannian geometry

Let  $(M^n, g)$  be a Riemannian manifold.

- ▶ Laplacian: writing  $g = \sum_{j,k} g_{jk}(x) dx_j dx_k$  and  $(g^{jk}(x)) = (g_{jk}(x))^{-1}$

$$\Delta_g = \sum_{j,k} g^{jk}(x) \partial_{x_j} \partial_{x_k} - \sum_{i,j,k} g^{jk}(x) \Gamma_{jk}^i(x) \partial_{x_i},$$

## Background on Riemannian geometry

Let  $(M^n, g)$  be a Riemannian manifold.

- ▶ Laplacian: writing  $g = \sum_{j,k} g_{jk}(x) dx_j dx_k$  and  $(g^{jk}(x)) = (g_{jk}(x))^{-1}$

$$\Delta_g = \sum_{j,k} g^{jk}(x) \partial_{x_j} \partial_{x_k} - \sum_{i,j,k} g^{jk}(x) \Gamma_{jk}^i(x) \partial_{x_i},$$

- ▶ Riemannian measure:

$$dv_g = |g(x)| dx_1 \cdots dx_n, \quad |g(x)| := \det(g_{jk}(x))^{1/2}$$

## Background on Riemannian geometry

Let  $(M^n, g)$  be a Riemannian manifold.

- ▶ Laplacian: writing  $g = \sum_{j,k} g_{jk}(x) dx_j dx_k$  and  $(g^{jk}(x)) = (g_{jk}(x))^{-1}$

$$\Delta_g = \sum_{j,k} g^{jk}(x) \partial_{x_j} \partial_{x_k} - \sum_{i,j,k} g^{jk}(x) \Gamma_{jk}^i(x) \partial_{x_i},$$

- ▶ Riemannian measure:

$$dv_g = |g(x)| dx_1 \cdots dx_n, \quad |g(x)| := \det(g_{jk}(x))^{1/2}$$

- ▶ Principal symbol:

$$p(x, \xi) = \sum_{j,k} g^{jk}(x) \xi_j \xi_k.$$

Invariantly defined as a function on  $T^*M$  (locally  $M \times \mathbb{R}^n$ )



## Background on Riemannian geometry

Let  $(M^n, g)$  be a Riemannian manifold.

- ▶ Laplacian: writing  $g = \sum_{j,k} g_{jk}(x) dx_j dx_k$  and  $(g^{jk}(x)) = (g_{jk}(x))^{-1}$

$$\Delta_g = \sum_{j,k} g^{jk}(x) \partial_{x_j} \partial_{x_k} - \sum_{i,j,k} g^{jk}(x) \Gamma_{jk}^i(x) \partial_{x_i},$$

- ▶ Riemannian measure:

$$dv_g = |g(x)| dx_1 \cdots dx_n, \quad |g(x)| := \det(g_{jk}(x))^{1/2}$$

- ▶ Principal symbol:

$$p(x, \xi) = \sum_{j,k} g^{jk}(x) \xi_j \xi_k.$$

Invariantly defined as a function on  $T^*M$  (locally  $M \times \mathbb{R}^n$ )

- ▶ Measure on  $T^*M$ :  $dx_1 \cdots dx_n d\xi_1 \cdots d\xi_n$

## Background on Riemannian geometry

Let  $(M^n, g)$  be a Riemannian manifold.

- ▶ Laplacian: writing  $g = \sum_{j,k} g_{jk}(x) dx_j dx_k$  and  $(g^{jk}(x)) = (g_{jk}(x))^{-1}$

$$\Delta_g = \sum_{j,k} g^{jk}(x) \partial_{x_j} \partial_{x_k} - \sum_{i,j,k} g^{jk}(x) \Gamma_{jk}^i(x) \partial_{x_i},$$

- ▶ Riemannian measure:

$$dv_g = |g(x)| dx_1 \cdots dx_n, \quad |g(x)| := \det(g_{jk}(x))^{1/2}$$

- ▶ Principal symbol:

$$p(x, \xi) = \sum_{j,k} g^{jk}(x) \xi_j \xi_k.$$

Invariantly defined as a function on  $T^*M$  (locally  $M \times \mathbb{R}^n$ )

- ▶ Measure on  $T^*M$ :  $dx_1 \cdots dx_n d\xi_1 \cdots d\xi_n$

**Rem:** setting  $\partial_x = (\partial_{x_1}, \dots, \partial_{x_n})$ ,

$$\Delta_g = p(x, \partial_x) + \text{lower order terms (in } \partial_x)$$

## Background on Riemannian geometry

- ▶ Geodesic flow:  $\phi^t : T^*M \rightarrow T^*M$ , i.e.  $\phi^t(x, \xi) := (x^t, \xi^t)$  with

$$\dot{x}_j^t = \frac{\partial p}{\partial \xi_j}(x^t, \xi^t), \quad \dot{\xi}_j^t = -\frac{\partial p}{\partial x_j}(x^t, \xi^t) \quad (x^t, \xi^t)|_{t=0} = (x, \xi)$$

The curves  $t \mapsto x^t$  are the geodesics (i.e. locally length minimizing) on  $M$

## Background on Riemannian geometry

- ▶ Geodesic flow:  $\phi^t : T^*M \rightarrow T^*M$ , i.e.  $\phi^t(x, \xi) := (x^t, \xi^t)$  with

$$\dot{x}_j^t = \frac{\partial p}{\partial \xi_j}(x^t, \xi^t), \quad \dot{\xi}_j^t = -\frac{\partial p}{\partial x_j}(x^t, \xi^t) \quad (x^t, \xi^t)|_{t=0} = (x, \xi)$$

The curves  $t \mapsto x^t$  are the geodesics (i.e. locally length minimizing) on  $M$

Let

$$L := -\Delta_g$$

**Fact:** If  $M$  is compact, there is an orthonormal basis of  $L^2(M, dv_g)$  of eigenfunctions of  $L$ :

$$Le_j = \lambda_j e_j, \quad 0 \leq \lambda_0 \leq \lambda_1 \leq \lambda_2 \leq \dots \quad \lambda_j \uparrow +\infty.$$

## Background on Riemannian geometry

- ▶ Geodesic flow:  $\phi^t : T^*M \rightarrow T^*M$ , i.e.  $\phi^t(x, \xi) := (x^t, \xi^t)$  with

$$\dot{x}_j^t = \frac{\partial p}{\partial \xi_j}(x^t, \xi^t), \quad \dot{\xi}_j^t = -\frac{\partial p}{\partial x_j}(x^t, \xi^t) \quad (x^t, \xi^t)|_{t=0} = (x, \xi)$$

The curves  $t \mapsto x^t$  are the geodesics (i.e. locally length minimizing) on  $M$

Let

$$L := -\Delta_g$$

**Fact:** If  $M$  is compact, there is an orthonormal basis of  $L^2(M, dv_g)$  of eigenfunctions of  $L$ :

$$Le_j = \lambda_j e_j, \quad 0 \leq \lambda_0 \leq \lambda_1 \leq \lambda_2 \leq \dots \quad \lambda_j \uparrow +\infty.$$

**Purpose of spectral geometry/semiclassical analysis:** to get information on

- ▶ Quantum data:  $\lambda_j$  and/or  $e_j$

## Background on Riemannian geometry

- ▶ Geodesic flow:  $\phi^t : T^*M \rightarrow T^*M$ , i.e.  $\phi^t(x, \xi) := (x^t, \xi^t)$  with

$$\dot{x}_j^t = \frac{\partial p}{\partial \xi_j}(x^t, \xi^t), \quad \dot{\xi}_j^t = -\frac{\partial p}{\partial x_j}(x^t, \xi^t) \quad (x^t, \xi^t)|_{t=0} = (x, \xi)$$

The curves  $t \mapsto x^t$  are the geodesics (i.e. locally length minimizing) on  $M$

Let

$$L := -\Delta_g$$

**Fact:** If  $M$  is compact, there is an orthonormal basis of  $L^2(M, dv_g)$  of eigenfunctions of  $L$ :

$$Le_j = \lambda_j e_j, \quad 0 \leq \lambda_0 \leq \lambda_1 \leq \lambda_2 \leq \dots \quad \lambda_j \uparrow +\infty.$$

**Purpose of spectral geometry/semiclassical analysis:** to get information on

- ▶ Quantum data:  $\lambda_j$  and/or  $e_j$

in term of

- ▶ Classical data: principal symbol, geodesic flow,...

in the semiclassical (or high energy) limit  $\lambda_j \rightarrow \infty$ .

# Examples of problems

## 1- **Distribution of eigenvalues:** the Weyl law

## Examples of problems

### 1- Distribution of eigenvalues: the Weyl law

$$\#\{j \mid \lambda_j \leq \lambda\} = (2\pi)^{-n} \text{vol}_g(M) \text{vol}(\mathbb{B}^n) \lambda^{\frac{n}{2}} (1 + o(1)), \quad \lambda \rightarrow \infty$$



## Examples of problems

**1- Distribution of eigenvalues:** the Weyl law

$$\#\{j \mid \lambda_j \leq \lambda\} = (2\pi)^{-n} \text{vol}_g(M) \text{vol}(\mathbb{B}^n) \lambda^{\frac{n}{2}} (1 + o(1)), \quad \lambda \rightarrow \infty$$

Using the semiclassical normalization

$$L_h := h^2 L, \quad \lambda =: h^{-2}$$

( $h \sim$  effective Planck's constant), this is a special case of

$$\#\text{spec}(L_h) \cap [a, b] = (2\pi h)^{-n} \text{vol}_{T^*M} (p^{-1}([a, b])) + h^{-n} \epsilon(h)$$

with  $a = 0$  and  $b = 1$  (global/macroscopic law).

## Examples of problems

1- **Distribution of eigenvalues:** the Weyl law

$$\#\{j \mid \lambda_j \leq \lambda\} = (2\pi)^{-n} \text{vol}_g(M) \text{vol}(\mathbb{B}^n) \lambda^{\frac{n}{2}} (1 + o(1)), \quad \lambda \rightarrow \infty$$

Using the semiclassical normalization

$$L_h := h^2 L, \quad \lambda =: h^{-2}$$

( $h \sim$  effective Planck's constant), this is a special case of

$$\#\text{spec}(L_h) \cap [a, b] = (2\pi h)^{-n} \text{vol}_{T^*M} (p^{-1}([a, b])) + h^{-n} \epsilon(h)$$

with  $a = 0$  and  $b = 1$  (global/macroscopic law).

**Rem:** Can be pushed to a microscopic law ( $b - a \ll 1$ ) as long as

$$\text{vol}_{T^*M} (p^{-1}([a, b])) \gg \epsilon(h)$$

## Examples of problems

**1- Distribution of eigenvalues:** the Weyl law

$$\#\{j \mid \lambda_j \leq \lambda\} = (2\pi)^{-n} \text{vol}_g(M) \text{vol}(\mathbb{B}^n) \lambda^{\frac{n}{2}} (1 + o(1)), \quad \lambda \rightarrow \infty$$

Using the semiclassical normalization

$$L_h := h^2 L, \quad \lambda =: h^{-2}$$

( $h \sim$  effective Planck's constant), this is a special case of

$$\#\text{spec}(L_h) \cap [a, b] = (2\pi h)^{-n} \text{vol}_{T^*M} (p^{-1}([a, b])) + h^{-n} \epsilon(h)$$

with  $a = 0$  and  $b = 1$  (global/macroscopic law).

**Rem:** Can be pushed to a microscopic law ( $b - a \ll 1$ ) as long as

$$\text{vol}_{T^*M} (p^{-1}([a, b])) \gg \epsilon(h)$$

**Question** How small is  $\epsilon(h)$  ?

## Examples of problems

- in general (Avakumovic-Hormander bound - optimal on the sphere)

$$\epsilon(h) = O(h)$$

## Examples of problems

- in general (Avakumovic-Hormander bound - optimal on the sphere)

$$\epsilon(h) = O(h)$$

- if the set of periodic geodesics has zero measure (Duistermaat-Guillemin bound)

$$\epsilon(h) = o(h)$$

## Examples of problems

- in general (Avakumovic-Hormander bound - optimal on the sphere)

$$\epsilon(h) = O(h)$$

- if the set of periodic geodesics has zero measure (Duistermaat-Guillemin bound)

$$\epsilon(h) = o(h)$$

- if the curvature is negative (Bérard bound)

$$\epsilon(h) = O\left(\frac{h}{\log(1/h)}\right)$$

## Examples of problems

- in general (Avakumovic-Hormander bound - optimal on the sphere)

$$\epsilon(h) = O(h)$$

- if the set of periodic geodesics has zero measure (Duistermaat-Guillemin bound)

$$\epsilon(h) = o(h)$$

- if the curvature is negative (Bérard bound)

$$\epsilon(h) = O\left(\frac{h}{\log(1/h)}\right)$$

**Open question:** can one improve on the last bound ?

## Examples of problems

- in general (Avakumovic-Hormander bound - optimal on the sphere)

$$\epsilon(h) = O(h)$$

- if the set of periodic geodesics has zero measure (Duistermaat-Guillemin bound)

$$\epsilon(h) = o(h)$$

- if the curvature is negative (Bérard bound)

$$\epsilon(h) = O\left(\frac{h}{\log(1/h)}\right)$$

**Open question:** can one improve on the last bound ?

**Comparison with RMT:** if  $N_h = \#\text{spec}(L_h) \cap [0, 1]$  (i.e. exactly  $N_h$  eigenvalues in  $[0, 1]$ )

$$\frac{1}{N_h} \#\text{spec}(L_h) \cap [a, b]$$



## Examples of problems

- in general (Avakumovic-Hormander bound - optimal on the sphere)

$$\epsilon(h) = O(h)$$

- if the set of periodic geodesics has zero measure (Duistermaat-Guillemin bound)

$$\epsilon(h) = o(h)$$

- if the curvature is negative (Bérard bound)

$$\epsilon(h) = O\left(\frac{h}{\log(1/h)}\right)$$

**Open question:** can one improve on the last bound ?

**Comparison with RMT:** if  $N_h = \#\text{spec}(L_h) \cap [0, 1]$  (i.e. exactly  $N_h$  eigenvalues in  $[0, 1]$ )

$$\frac{1}{N_h} \#\text{spec}(L_h) \cap [a, b] = \frac{\text{vol}_{T^*M}(\rho^{-1}([a, b]))}{\text{vol}_{T^*M}(\rho^{-1}([0, 1]))} + \tilde{\epsilon}(N_h)$$

## Examples of problems

- in general (Avakumovic-Hormander bound - optimal on the sphere)

$$\epsilon(h) = O(h)$$

- if the set of periodic geodesics has zero measure (Duistermaat-Guillemin bound)

$$\epsilon(h) = o(h)$$

- if the curvature is negative (Bérard bound)

$$\epsilon(h) = O\left(\frac{h}{\log(1/h)}\right)$$

**Open question:** can one improve on the last bound ?

**Comparison with RMT:** if  $N_h = \#\text{spec}(L_h) \cap [0, 1]$  (i.e. exactly  $N_h$  eigenvalues in  $[0, 1]$ )

$$\frac{1}{N_h} \#\text{spec}(L_h) \cap [a, b] = \frac{\text{vol}_{T^*M}(\rho^{-1}([a, b]))}{\text{vol}_{T^*M}(\rho^{-1}([0, 1]))} + \tilde{\epsilon}(N_h) = \frac{n}{2} \int_a^b \lambda^{\frac{n}{2}-1} d\lambda + \tilde{\epsilon}(N_h)$$

## Examples of problems

- in general (Avakumovic-Hormander bound - optimal on the sphere)

$$\epsilon(h) = O(h)$$

- if the set of periodic geodesics has zero measure (Duistermaat-Guillemin bound)

$$\epsilon(h) = o(h)$$

- if the curvature is negative (Bérard bound)

$$\epsilon(h) = O\left(\frac{h}{\log(1/h)}\right)$$

**Open question:** can one improve on the last bound ?

**Comparison with RMT:** if  $N_h = \#\text{spec}(L_h) \cap [0, 1]$  (i.e. exactly  $N_h$  eigenvalues in  $[0, 1]$ )

$$\frac{1}{N_h} \#\text{spec}(L_h) \cap [a, b] = \frac{\text{vol}_{T^*M}(\rho^{-1}([a, b]))}{\text{vol}_{T^*M}(\rho^{-1}([0, 1]))} + \tilde{\epsilon}(N_h) = \frac{n}{2} \int_a^b \lambda^{\frac{n}{2}-1} d\lambda + \tilde{\epsilon}(N_h)$$

as long as  $b - a \gg \tilde{\epsilon}(N_h)$  and with  $\tilde{\epsilon}(N_h)$  either

$$O(N_h^{-\frac{1}{n}}) \quad o(N_h^{-\frac{1}{n}}) \quad O\left(\frac{N_h^{-\frac{1}{n}}}{\log N_h}\right)$$

## Examples of problems

2-  $L^q$  norms of eigenfunctions: Sogge's estimates (see also Seger, Zelditch,...)

## Examples of problems

2-  $L^q$  norms of eigenfunctions: Sogge's estimates (see also Seger, Zelditch,...)

$$L_h e_h = e_h \quad \implies \quad \|e_h\|_q \lesssim h^{-\mu(q)} \|e_h\|_2$$

(here  $h^{-2} = \lambda_j$  is a large eigenvalue of  $L$  and  $e_h = e_j$ )

## Examples of problems

2-  $L^q$  norms of eigenfunctions: Sogge's estimates (see also Seger, Zelditch,...)

$$L_h e_h = e_h \quad \implies \quad \|e_h\|_q \lesssim h^{-\mu(q)} \|e_h\|_2$$

(here  $h^{-2} = \lambda_j$  is a large eigenvalue of  $L$  and  $e_h = e_j$ ) with

$$\mu(q) = \max\left(\frac{n-1}{4} - \frac{n-1}{2q}, \frac{n}{2} - \frac{n}{q} - \frac{1}{2}\right), \quad q \in [2, \infty]$$

## Examples of problems

2-  $L^q$  norms of eigenfunctions: Sogge's estimates (see also Seger, Zelditch,...)

$$L_h e_h = e_h \quad \implies \quad \|e_h\|_q \lesssim h^{-\mu(q)} \|e_h\|_2$$

(here  $h^{-2} = \lambda_j$  is a large eigenvalue of  $L$  and  $e_h = e_j$ ) with

$$\mu(q) = \max\left(\frac{n-1}{4} - \frac{n-1}{2q}, \frac{n}{2} - \frac{n}{q} - \frac{1}{2}\right), \quad q \in [2, \infty]$$

- Optimal on the sphere

## Examples of problems

2-  $L^q$  norms of eigenfunctions: Sogge's estimates (see also Seger, Zelditch,...)

$$L_h e_h = e_h \quad \implies \quad \|e_h\|_q \lesssim h^{-\mu(q)} \|e_h\|_2$$

(here  $h^{-2} = \lambda_j$  is a large eigenvalue of  $L$  and  $e_h = e_j$ ) with

$$\mu(q) = \max\left(\frac{n-1}{4} - \frac{n-1}{2q}, \frac{n}{2} - \frac{n}{q} - \frac{1}{2}\right), \quad q \in [2, \infty]$$

- Optimal on the sphere
- Don't worry about the precise numerology (in this talk): just remember this is better than (trivial) Sobolev bounds

$$\|\chi(L_h)\|_{2 \rightarrow q} \lesssim_\chi h^{-\sigma(q)}, \quad \sigma(q) = \frac{n}{2} - \frac{n}{q}, \quad \chi \in C_0^\infty(\mathbb{R}).$$

(think of  $\chi(-h^2 \Delta)$  on  $\mathbb{R}^n$  whose kernel is of the form  $h^{-n} f\left(\frac{x-y}{h}\right)$  with  $f \in \mathcal{S}(\mathbb{R})$ )



## Examples of problems

2-  $L^q$  norms of eigenfunctions: Sogge's estimates (see also Seger, Zelditch,...)

$$L_h e_h = e_h \quad \implies \quad \|e_h\|_q \lesssim h^{-\mu(q)} \|e_h\|_2$$

(here  $h^{-2} = \lambda_j$  is a large eigenvalue of  $L$  and  $e_h = e_j$ ) with

$$\mu(q) = \max\left(\frac{n-1}{4} - \frac{n-1}{2q}, \frac{n}{2} - \frac{n}{q} - \frac{1}{2}\right), \quad q \in [2, \infty]$$

- Optimal on the sphere
- Don't worry about the precise numerology (in this talk): just remember this is better than (trivial) Sobolev bounds

$$\|\chi(L_h)\|_{2 \rightarrow q} \lesssim_\chi h^{-\sigma(q)}, \quad \sigma(q) = \frac{n}{2} - \frac{n}{q}, \quad \chi \in C_0^\infty(\mathbb{R}).$$

(think of  $\chi(-h^2 \Delta)$  on  $\mathbb{R}^n$  whose kernel is of the form  $h^{-n} f\left(\frac{x-y}{h}\right)$  with  $f \in \mathcal{S}(\mathbb{R})$ )

- There are log improvements in negative curvature

## Examples of problems

2-  $L^q$  norms of eigenfunctions: Sogge's estimates (see also Seger, Zelditch,...)

$$L_h e_h = e_h \quad \implies \quad \|e_h\|_q \lesssim h^{-\mu(q)} \|e_h\|_2$$

(here  $h^{-2} = \lambda_j$  is a large eigenvalue of  $L$  and  $e_h = e_j$ ) with

$$\mu(q) = \max\left(\frac{n-1}{4} - \frac{n-1}{2q}, \frac{n}{2} - \frac{n}{q} - \frac{1}{2}\right), \quad q \in [2, \infty]$$

- Optimal on the sphere
- Don't worry about the precise numerology (in this talk): just remember this is better than (trivial) Sobolev bounds

$$\|\chi(L_h)\|_{2 \rightarrow q} \lesssim_\chi h^{-\sigma(q)}, \quad \sigma(q) = \frac{n}{2} - \frac{n}{q}, \quad \chi \in C_0^\infty(\mathbb{R}).$$

(think of  $\chi(-h^2 \Delta)$  on  $\mathbb{R}^n$  whose kernel is of the form  $h^{-n} f\left(\frac{x-y}{h}\right)$  with  $f \in \mathcal{S}(\mathbb{R})$ )

- There are log improvements in negative curvature and stronger ones in arithmetic cases (Iwaniec-Sarnak): if  $M = \Gamma \backslash \mathbb{H}^2$  and  $e_h$  are Hecke eigenfunctions, then

$$\|e_h\|_\infty \lesssim_\epsilon h^{-\frac{5}{12} - \epsilon} \|e_h\|_2$$

(in dimension 2,  $\mu(\infty) = \frac{1}{2}$ )

## Examples of problems

2-  $L^q$  norms of eigenfunctions: Sogge's estimates (see also Seger, Zelditch,...)

$$L_h e_h = e_h \quad \implies \quad \|e_h\|_q \lesssim h^{-\mu(q)} \|e_h\|_2$$

(here  $h^{-2} = \lambda_j$  is a large eigenvalue of  $L$  and  $e_h = e_j$ ) with

$$\mu(q) = \max\left(\frac{n-1}{4} - \frac{n-1}{2q}, \frac{n}{2} - \frac{n}{q} - \frac{1}{2}\right), \quad q \in [2, \infty]$$

- Optimal on the sphere
- Don't worry about the precise numerology (in this talk): just remember this is better than (trivial) Sobolev bounds

$$\|\chi(L_h)\|_{2 \rightarrow q} \lesssim_\chi h^{-\sigma(q)}, \quad \sigma(q) = \frac{n}{2} - \frac{n}{q}, \quad \chi \in C_0^\infty(\mathbb{R}).$$

(think of  $\chi(-h^2 \Delta)$  on  $\mathbb{R}^n$  whose kernel is of the form  $h^{-n} f\left(\frac{x-y}{h}\right)$  with  $f \in \mathcal{S}(\mathbb{R})$ )

- There are log improvements in negative curvature and stronger ones in arithmetic cases (Iwaniec-Sarnak): if  $M = \Gamma \backslash \mathbb{H}^2$  and  $e_h$  are Hecke eigenfunctions, then

$$\|e_h\|_\infty \lesssim_\epsilon h^{-\frac{5}{12} - \epsilon} \|e_h\|_2$$

(in dimension 2,  $\mu(\infty) = \frac{1}{2}$ )

**Open question:** can one get similar improvement in non arithmetic cases ?

## Basic strategy

For the previous problems, we need to analyze spectral projections  $\mathbb{1}_{[a,b]}(L_h)$

## Basic strategy

For the previous problems, we need to analyze spectral projections  $\mathbb{1}_{[a,b]}(L_h)$

- Weyl's law:

$$\#\text{spec}(L_h) \cap [a, b] = \text{tr} (\mathbb{1}_{[a,b]}(L_h))$$

## Basic strategy

For the previous problems, we need to analyze spectral projections  $\mathbb{1}_{[a,b]}(L_h)$

- Weyl's law:

$$\#\text{spec}(L_h) \cap [a, b] = \text{tr} (\mathbb{1}_{[a,b]}(L_h))$$

- Sogge's bounds:

$$\|e_h\|_q \leq \|\mathbb{1}_{[a,b]}(L_h)\|_{2 \rightarrow q} \|e_h\|_{L^2}, \quad a \leq 1 \leq b$$

## Basic strategy

For the previous problems, we need to analyze spectral projections  $\mathbb{1}_{[a,b]}(L_h)$

- Weyl's law:

$$\#\text{spec}(L_h) \cap [a, b] = \text{tr} (\mathbb{1}_{[a,b]}(L_h))$$

- Sogge's bounds:

$$\|e_h\|_q \leq \|\mathbb{1}_{[a,b]}(L_h)\|_{2 \rightarrow q} \|e_h\|_{L^2}, \quad a \leq 1 \leq b$$

**Pbm:** the fine structure of the projections is hard to touch.

## Basic strategy

For the previous problems, we need to analyze spectral projections  $\mathbb{1}_{[a,b]}(L_h)$

- Weyl's law:

$$\#\text{spec}(L_h) \cap [a, b] = \text{tr}(\mathbb{1}_{[a,b]}(L_h))$$

- Sogge's bounds:

$$\|e_h\|_q \leq \|\mathbb{1}_{[a,b]}(L_h)\|_{2 \rightarrow q} \|e_h\|_{L^2}, \quad a \leq 1 \leq b$$

**Pbm:** the fine structure of the projections is hard to touch.

→ Rather than  $\mathbb{1}_{[a,b]}(\lambda)$  (which is not smooth), we will consider

$$\rho_{\epsilon(h)}(\lambda) := \rho\left(\frac{\lambda}{\epsilon(h)}\right)$$

with *suitable*  $\rho \in \mathcal{S}(\mathbb{R})$



## Basic strategy

For the previous problems, we need to analyze spectral projections  $\mathbb{1}_{[a,b]}(L_h)$

- Weyl's law:

$$\#\text{spec}(L_h) \cap [a, b] = \text{tr} (\mathbb{1}_{[a,b]}(L_h))$$

- Sogge's bounds:

$$\|e_h\|_q \leq \|\mathbb{1}_{[a,b]}(L_h)\|_{2 \rightarrow q} \|e_h\|_{L^2}, \quad a \leq 1 \leq b$$

**Pbm:** the fine structure of the projections is hard to touch.

→ Rather than  $\mathbb{1}_{[a,b]}(\lambda)$  (which is not smooth), we will consider

$$\rho_{\epsilon(h)}(\lambda) := \rho\left(\frac{\lambda}{\epsilon(h)}\right)$$

with suitable  $\rho \in \mathcal{S}(\mathbb{R})$

**Principle of the proof of Sogge's bounds :** Choose  $\rho$  s.t.  $\rho(0) = 1$  so that

$$\rho_{\epsilon(h)}(L_h - 1)e_h = e_h.$$

and suitably so that, with  $\epsilon(h) = h$ ,

$$\|\rho_{\epsilon(h)}(L_h - 1)\|_{2 \rightarrow q} \lesssim h^{-\mu(q)}$$

## Basic strategy

For the previous problems, we need to analyze spectral projections  $\mathbb{1}_{[a,b]}(L_h)$

- Weyl's law:

$$\#\text{spec}(L_h) \cap [a, b] = \text{tr} (\mathbb{1}_{[a,b]}(L_h))$$

- Sogge's bounds:

$$\|e_h\|_q \leq \|\mathbb{1}_{[a,b]}(L_h)\|_{2 \rightarrow q} \|e_h\|_{L^2}, \quad a \leq 1 \leq b$$

**Pbm:** the fine structure of the projections is hard to touch.

→ Rather than  $\mathbb{1}_{[a,b]}(\lambda)$  (which is not smooth), we will consider

$$\rho_{\epsilon(h)}(\lambda) := \rho\left(\frac{\lambda}{\epsilon(h)}\right)$$

with *suitable*  $\rho \in \mathcal{S}(\mathbb{R})$

**Principle of the proof of Sogge's bounds :** Choose  $\rho$  s.t.  $\rho(0) = 1$  so that

$$\rho_{\epsilon(h)}(L_h - 1)e_h = e_h.$$

and suitably so that, with  $\epsilon(h) = h$ ,

$$\|\rho_{\epsilon(h)}(L_h - 1)\|_{2 \rightarrow q} \lesssim h^{-\mu(q)}$$

Suitably means:  $\hat{\rho}$  is compactly supported in some  $[\delta, 2\delta]$ ,  $0 < \delta \ll 1$

## Basic strategy

**For the Weyl law:** Let  $\chi \in C_0^\infty$  be equal to 1 on  $[a, b]$ . Then

$$\operatorname{tr} [\mathbb{1}_{[a,b]}(L_h)] = \operatorname{tr} [\chi(L_h)\mathbb{1}_{[a,b]}(L_h)]$$

## Basic strategy

**For the Weyl law:** Let  $\chi \in C_0^\infty$  be equal to 1 on  $[a, b]$ . Then

$$\begin{aligned}\mathrm{tr} [\mathbb{1}_{[a,b]}(L_h)] &= \mathrm{tr} [\chi(L_h)\mathbb{1}_{[a,b]}(L_h)] \\ &= \frac{1}{\epsilon(h)} \mathrm{tr} \left[ \chi(L_h) \left( \frac{1}{\epsilon(h)} \rho_{\epsilon(h)} * \mathbb{1}_{[a,b]} \right) (L_h) \right] + R(h)\end{aligned}$$

## Basic strategy

**For the Weyl law:** Let  $\chi \in C_0^\infty$  be equal to 1 on  $[a, b]$ . Then

$$\begin{aligned}\mathrm{tr} [\mathbb{1}_{[a,b]}(L_h)] &= \mathrm{tr} [\chi(L_h)\mathbb{1}_{[a,b]}(L_h)] \\ &= \frac{1}{\epsilon(h)} \mathrm{tr} \left[ \chi(L_h) \left( \frac{1}{\epsilon(h)} \rho_{\epsilon(h)} * \mathbb{1}_{[a,b]} \right) (L_h) \right] + R(h) \\ &= \int_a^b \frac{1}{\epsilon(h)} \mathrm{tr} \left[ \frac{1}{\epsilon(h)} \chi(L_h) \rho_{\epsilon(h)}(L_h - \lambda) \right] d\lambda + R(h)\end{aligned}$$

## Basic strategy

**For the Weyl law:** Let  $\chi \in C_0^\infty$  be equal to 1 on  $[a, b]$ . Then

$$\begin{aligned}\mathrm{tr} [\mathbb{1}_{[a,b]}(L_h)] &= \mathrm{tr} [\chi(L_h)\mathbb{1}_{[a,b]}(L_h)] \\ &= \frac{1}{\epsilon(h)} \mathrm{tr} \left[ \chi(L_h) \left( \frac{1}{\epsilon(h)} \rho_{\epsilon(h)} * \mathbb{1}_{[a,b]} \right) (L_h) \right] + R(h) \\ &= \int_a^b \frac{1}{\epsilon(h)} \mathrm{tr} \left[ \frac{1}{\epsilon(h)} \chi(L_h) \rho_{\epsilon(h)}(L_h - \lambda) \right] d\lambda + R(h)\end{aligned}$$

where

$$\frac{1}{\epsilon(h)} \rho_{\epsilon(h)}(L_h - \lambda)$$

## Basic strategy

**For the Weyl law:** Let  $\chi \in C_0^\infty$  be equal to 1 on  $[a, b]$ . Then

$$\begin{aligned}\operatorname{tr} [\mathbb{1}_{[a,b]}(L_h)] &= \operatorname{tr} [\chi(L_h)\mathbb{1}_{[a,b]}(L_h)] \\ &= \frac{1}{\epsilon(h)} \operatorname{tr} \left[ \chi(L_h) \left( \frac{1}{\epsilon(h)} \rho_{\epsilon(h)} * \mathbb{1}_{[a,b]} \right) (L_h) \right] + R(h) \\ &= \int_a^b \frac{1}{\epsilon(h)} \operatorname{tr} \left[ \frac{1}{\epsilon(h)} \chi(L_h) \rho_{\epsilon(h)}(L_h - \lambda) \right] d\lambda + R(h)\end{aligned}$$

where

$$\frac{1}{\epsilon(h)} \rho_{\epsilon(h)}(L_h - \lambda) = \frac{1}{2\pi h} \int_{\mathbb{R}} e^{i\frac{t}{h}(L_h - \lambda)} \hat{\rho} \left( \frac{t}{T(h)} \right) dt$$

with

$$T(h) = \frac{h}{\epsilon(h)}$$

## Basic strategy

**For the Weyl law:** Let  $\chi \in C_0^\infty$  be equal to 1 on  $[a, b]$ . Then

$$\begin{aligned}\mathrm{tr} [\mathbb{1}_{[a,b]}(L_h)] &= \mathrm{tr} [\chi(L_h)\mathbb{1}_{[a,b]}(L_h)] \\ &= \frac{1}{\epsilon(h)} \mathrm{tr} \left[ \chi(L_h) \left( \frac{1}{\epsilon(h)} \rho_{\epsilon(h)} * \mathbb{1}_{[a,b]} \right) (L_h) \right] + R(h) \\ &= \int_a^b \frac{1}{\epsilon(h)} \mathrm{tr} \left[ \frac{1}{\epsilon(h)} \chi(L_h) \rho_{\epsilon(h)}(L_h - \lambda) \right] d\lambda + R(h)\end{aligned}$$

where

$$\frac{1}{\epsilon(h)} \rho_{\epsilon(h)}(L_h - \lambda) = \frac{1}{2\pi h} \int_{\mathbb{R}} e^{i\frac{t}{h}(L_h - \lambda)} \hat{\rho} \left( \frac{t}{T(h)} \right) dt$$

with

$$T(h) = \frac{h}{\epsilon(h)}$$

Here we use the **propagator**  $e^{i\frac{t}{h}L_h}$ .



## Basic strategy

**For the Weyl law:** Let  $\chi \in C_0^\infty$  be equal to 1 on  $[a, b]$ . Then

$$\begin{aligned}\mathrm{tr} [\mathbb{1}_{[a,b]}(L_h)] &= \mathrm{tr} [\chi(L_h)\mathbb{1}_{[a,b]}(L_h)] \\ &= \frac{1}{\epsilon(h)} \mathrm{tr} \left[ \chi(L_h) \left( \frac{1}{\epsilon(h)} \rho_{\epsilon(h)} * \mathbb{1}_{[a,b]} \right) (L_h) \right] + R(h) \\ &= \int_a^b \frac{1}{\epsilon(h)} \mathrm{tr} \left[ \frac{1}{\epsilon(h)} \chi(L_h) \rho_{\epsilon(h)}(L_h - \lambda) \right] d\lambda + R(h)\end{aligned}$$

where

$$\frac{1}{\epsilon(h)} \rho_{\epsilon(h)}(L_h - \lambda) = \frac{1}{2\pi h} \int_{\mathbb{R}} e^{i\frac{t}{h}(L_h - \lambda)} \hat{\rho} \left( \frac{t}{T(h)} \right) dt$$

with

$$T(h) = \frac{h}{\epsilon(h)}$$

Here we use the **propagator**  $e^{i\frac{t}{h}L_h}$ . Note that it solves the Schroedinger equation

$$ih\partial_t e^{i\frac{t}{h}L_h} = h^2 \Delta_g e^{i\frac{t}{h}L_h}, \quad e^{i\frac{t}{h}L_h} \Big|_{t=0} = I$$

## Basic strategy

**For the Weyl law:** Let  $\chi \in C_0^\infty$  be equal to 1 on  $[a, b]$ . Then

$$\begin{aligned}\mathrm{tr} [\mathbb{1}_{[a,b]}(L_h)] &= \mathrm{tr} [\chi(L_h)\mathbb{1}_{[a,b]}(L_h)] \\ &= \frac{1}{\epsilon(h)} \mathrm{tr} \left[ \chi(L_h) \left( \frac{1}{\epsilon(h)} \rho_{\epsilon(h)} * \mathbb{1}_{[a,b]} \right) (L_h) \right] + R(h) \\ &= \int_a^b \frac{1}{\epsilon(h)} \mathrm{tr} \left[ \frac{1}{\epsilon(h)} \chi(L_h) \rho_{\epsilon(h)}(L_h - \lambda) \right] d\lambda + R(h)\end{aligned}$$

where

$$\frac{1}{\epsilon(h)} \rho_{\epsilon(h)}(L_h - \lambda) = \frac{1}{2\pi h} \int_{\mathbb{R}} e^{i\frac{t}{h}(L_h - \lambda)} \hat{\rho} \left( \frac{t}{T(h)} \right) dt$$

with

$$T(h) = \frac{h}{\epsilon(h)}$$

Here we use the **propagator**  $e^{i\frac{t}{h}L_h}$ . Note that it solves the Schroedinger equation

$$i\hbar \partial_t e^{i\frac{t}{h}L_h} = \hbar^2 \Delta_g e^{i\frac{t}{h}L_h}, \quad e^{i\frac{t}{h}L_h} \Big|_{t=0} = I$$

**Core of the proof:** try to get a precise description over times  $|t| \lesssim T(h)$  of

$$e^{i\frac{t}{h}(L_h - \lambda)} \chi(L_h)$$

## Basic strategy

**For the Weyl law:** Let  $\chi \in C_0^\infty$  be equal to 1 on  $[a, b]$ . Then

$$\begin{aligned}\mathrm{tr} [\mathbb{1}_{[a,b]}(L_h)] &= \mathrm{tr} [\chi(L_h)\mathbb{1}_{[a,b]}(L_h)] \\ &= \frac{1}{\epsilon(h)} \mathrm{tr} \left[ \chi(L_h) \left( \frac{1}{\epsilon(h)} \rho_{\epsilon(h)} * \mathbb{1}_{[a,b]} \right) (L_h) \right] + R(h) \\ &= \int_a^b \frac{1}{\epsilon(h)} \mathrm{tr} \left[ \frac{1}{\epsilon(h)} \chi(L_h) \rho_{\epsilon(h)}(L_h - \lambda) \right] d\lambda + R(h)\end{aligned}$$

where

$$\frac{1}{\epsilon(h)} \rho_{\epsilon(h)}(L_h - \lambda) = \frac{1}{2\pi h} \int_{\mathbb{R}} e^{i\frac{t}{h}(L_h - \lambda)} \hat{\rho} \left( \frac{t}{T(h)} \right) dt$$

with

$$T(h) = \frac{h}{\epsilon(h)}$$

Here we use the **propagator**  $e^{i\frac{t}{h}L_h}$ . Note that it solves the Schroedinger equation

$$i\hbar \partial_t e^{i\frac{t}{h}L_h} = \hbar^2 \Delta_g e^{i\frac{t}{h}L_h}, \quad e^{i\frac{t}{h}L_h} \Big|_{t=0} = I$$

**Core of the proof:** try to get a precise description over times  $|t| \lesssim T(h)$  of

$$e^{i\frac{t}{h}(L_h - \lambda)} \chi(L_h)$$

**Rem 1:** We use the propagator  $e^{i\frac{t}{h}L_h}$  rather than the resolvent  $(L_h - z)^{-1}$

## Basic strategy

**For the Weyl law:** Let  $\chi \in C_0^\infty$  be equal to 1 on  $[a, b]$ . Then

$$\begin{aligned}\mathrm{tr} [\mathbb{1}_{[a,b]}(L_h)] &= \mathrm{tr} [\chi(L_h)\mathbb{1}_{[a,b]}(L_h)] \\ &= \frac{1}{\epsilon(h)} \mathrm{tr} \left[ \chi(L_h) \left( \frac{1}{\epsilon(h)} \rho_{\epsilon(h)} * \mathbb{1}_{[a,b]} \right) (L_h) \right] + R(h) \\ &= \int_a^b \frac{1}{\epsilon(h)} \mathrm{tr} \left[ \frac{1}{\epsilon(h)} \chi(L_h) \rho_{\epsilon(h)}(L_h - \lambda) \right] d\lambda + R(h)\end{aligned}$$

where

$$\frac{1}{\epsilon(h)} \rho_{\epsilon(h)}(L_h - \lambda) = \frac{1}{2\pi h} \int_{\mathbb{R}} e^{i\frac{t}{h}(L_h - \lambda)} \hat{\rho} \left( \frac{t}{T(h)} \right) dt$$

with

$$T(h) = \frac{h}{\epsilon(h)}$$

Here we use the **propagator**  $e^{i\frac{t}{h}L_h}$ . Note that it solves the Schroedinger equation

$$i\hbar \partial_t e^{i\frac{t}{h}L_h} = \hbar^2 \Delta_g e^{i\frac{t}{h}L_h}, \quad e^{i\frac{t}{h}L_h} \Big|_{t=0} = I$$

**Core of the proof:** try to get a precise description over times  $|t| \lesssim T(h)$  of

$$e^{i\frac{t}{h}(L_h - \lambda)} \chi(L_h)$$

**Rem 1:** We use the propagator  $e^{i\frac{t}{h}L_h}$  rather than the resolvent  $(L_h - z)^{-1}$

**Rem 2:** the control of  $R(h)$  ( $= O(h^{-n}\epsilon(h))$ ) is also based on this description.

## Basic strategy

We then approach  $e^{i\frac{t}{h}L_h}\chi(L_h)$  using **Fourier Integral Operators**,

## Basic strategy

We then approach  $e^{i\frac{t}{h}L_h}\chi(L_h)$  using **Fourier Integral Operators**, i.e. operators with integral kernels of the form

$$K_h(t, x, y) = h^{-d(D)} \int_{\mathbb{R}^D} e^{i\frac{1}{h}F(t, x, y, Z)} A(t, x, y, Z, h) dZ, \quad x, y \in M$$

## Basic strategy

We then approach  $e^{i\frac{t}{h}L_h}\chi(L_h)$  using **Fourier Integral Operators**, i.e. operators with integral kernels of the form

$$K_h(t, x, y) = h^{-d(D)} \int_{\mathbb{R}^D} e^{i\frac{1}{h}F(t, x, y, Z)} A(t, x, y, Z, h) dZ, \quad x, y \in M$$

where

- ▶ the phase  $F$  is given by the geodesic flow

## Basic strategy

We then approach  $e^{i\frac{t}{h}L_h}\chi(L_h)$  using **Fourier Integral Operators**, i.e. operators with integral kernels of the form

$$K_h(t, x, y) = h^{-d(D)} \int_{\mathbb{R}^D} e^{i\frac{1}{h}F(t, x, y, Z)} A(t, x, y, Z, h) dZ, \quad x, y \in M$$

where

- ▶ the phase  $F$  is given by the geodesic flow (also  $F \in C^\infty$ ,  $\text{Im}(F) \geq 0$ )



## Basic strategy

We then approach  $e^{i\frac{\hbar}{h}L_h}\chi(L_h)$  using **Fourier Integral Operators**, i.e. operators with integral kernels of the form

$$K_h(t, x, y) = h^{-d(D)} \int_{\mathbb{R}^D} e^{i\frac{1}{h}F(t, x, y, Z)} A(t, x, y, Z, h) dZ, \quad x, y \in M$$

where

- ▶ the phase  $F$  is given by the geodesic flow (also  $F \in C^\infty$ ,  $\text{Im}(F) \geq 0$ )
- ▶ the amplitude has an asymptotic expansion

$$A \sim A_0 + h^{1/2}A_1 + hA_2 + \dots$$

## Basic strategy

We then approach  $e^{i\frac{t}{h}L_h}\chi(L_h)$  using **Fourier Integral Operators**, i.e. operators with integral kernels of the form

$$K_h(t, x, y) = h^{-d(D)} \int_{\mathbb{R}^D} e^{i\frac{1}{h}F(t, x, y, Z)} A(t, x, y, Z, h) dZ, \quad x, y \in M$$

where

- ▶ the phase  $F$  is given by the geodesic flow (also  $F \in C^\infty$ ,  $\text{Im}(F) \geq 0$ )
- ▶ the amplitude has an asymptotic expansion

$$A \sim A_0 + h^{1/2}A_1 + hA_2 + \dots$$

One can then retrieve information on  $\rho_{\epsilon(h)}(L_h - \lambda)\chi(L_h)$  by using stationary phase asymptotics in

$$h^{-d(D)} \int_{\mathbb{R}} \int_{\mathbb{R}^D} e^{i\frac{1}{h}(F(t, x, y, Z) - t\lambda)} A(t, x, y, Z, h) \hat{\rho}\left(\frac{t}{T(h)}\right) dZ dt$$

or for the trace

$$h^{-1-d(D)} \int_M \int_{\mathbb{R}} \int_{\mathbb{R}^D} e^{i\frac{1}{h}(F(t, x, x, Z) - t\lambda)} A(t, x, x, Z, h) \hat{\rho}\left(\frac{t}{T(h)}\right) dZ dt dv_g(x)$$

## Basic strategy

In practice

$$A_k = O(e^{k\gamma|t|})$$

with  $e^{\gamma|t|}$  an upper bound bound on  $D\phi^t$ ,

## Basic strategy

In practice

$$A_k = O(e^{k\gamma|t|})$$

with  $e^{\gamma|t|}$  an upper bound bound on  $D\phi^t$ , so we get an asymptotic expansion provided

$$h^{\frac{1}{2}} e^{\gamma|t|} \ll 1$$

## Basic strategy

In practice

$$A_k = O(e^{k\gamma|t|})$$

with  $e^{\gamma|t|}$  an upper bound bound on  $D\phi^t$ , so we get an asymptotic expansion provided

$$h^{\frac{1}{2}} e^{\gamma|t|} \ll 1$$

**Consequence:** such approximations are in general limited to  $|t| \leq c \log(1/h)$

## Relationship with the geodesic flow

**Theorem [Propagator approximation]** (B. 2016+)

## Relationship with the geodesic flow

**Theorem [Propagator approximation]** (B. 2016+) *Let  $\psi \in C^\infty(M)$  be supported in a coordinate patch  $U$ . The integral kernel of  $e^{-i\frac{\epsilon}{\hbar}L_h}\chi(L_h)\psi$*

## Relationship with the geodesic flow

**Theorem [Propagator approximation]** (B. 2016+) *Let  $\psi \in C^\infty(M)$  be supported in a coordinate patch  $U$ . The integral kernel of  $e^{-i\frac{t}{h}L_h}\chi(L_h)\psi$  is well approximated by*

$$K_h(t, x, y) = h^{-\frac{3n}{2}} \int \int_{T^*U} A(t, x, y, z, \zeta, h) \exp \frac{i}{h} F(t, x, y, z, \zeta) dz d\zeta$$

for times  $|t| \leq c |\log h|$ .



## Relationship with the geodesic flow

**Theorem [Propagator approximation]** (B. 2016+) *Let  $\psi \in C^\infty(M)$  be supported in a coordinate patch  $U$ . The integral kernel of  $e^{-i\frac{t}{h}L_h}\chi(L_h)\psi$  is well approximated by*

$$K_h(t, x, y) = h^{-\frac{3n}{2}} \int \int_{T^*U} A(t, x, y, z, \zeta, h) \exp \frac{i}{h} F(t, x, y, z, \zeta) dz d\zeta$$

for times  $|t| \leq c|\log h|$ . The **phase** reads

$$F = tp(z, \zeta) + \zeta^t \cdot W_{z^t}^x + \frac{1}{2} \left\langle \Gamma_{(z, \zeta)}^t W_{z^t}^x, W_{z^t}^x \right\rangle_{z^t} - \zeta \cdot W_z^y + \frac{1}{2} \left\langle \widetilde{\Gamma_{(z, \zeta)}^0} W_z^y, W_z^y \right\rangle_z$$

## Relationship with the geodesic flow

**Theorem [Propagator approximation]** (B. 2016+) *Let  $\psi \in C^\infty(M)$  be supported in a coordinate patch  $U$ . The integral kernel of  $e^{-i\frac{t}{h}L_h}\chi(L_h)\psi$  is well approximated by*

$$K_h(t, x, y) = h^{-\frac{3n}{2}} \int \int_{T^*U} A(t, x, y, z, \zeta, h) \exp \frac{i}{h} F(t, x, y, z, \zeta) dz d\zeta$$

for times  $|t| \leq c|\log h|$ . The phase reads

$$F = tp(z, \zeta) + \zeta^t \cdot W_{z^t}^x + \frac{1}{2} \left\langle \Gamma_{(z, \zeta)}^t W_{z^t}^x, W_{z^t}^x \right\rangle_{z^t} - \zeta \cdot W_z^y + \frac{1}{2} \left\langle \widetilde{\Gamma_{(z, \zeta)}^0} W_z^y, W_z^y \right\rangle_z$$

where  $\Gamma_{(z, \zeta)}^t : T_{z^t}M^{\mathbb{C}} \rightarrow T_{z^t}M^{\mathbb{C}}$  is complex linear, symmetric with positive definite imaginary part, and solves

$$\nabla_{\dot{z}^t} \Gamma_{(z, \zeta)}^t = -R_{z^t}(\cdot, \dot{z}^t) \dot{z}^t - (\Gamma_{(z, \zeta)}^t)^2,$$

## Relationship with the geodesic flow

**Theorem [Propagator approximation]** (B. 2016+) *Let  $\psi \in C^\infty(M)$  be supported in a coordinate patch  $U$ . The integral kernel of  $e^{-i\frac{t}{h}L_h}\chi(L_h)\psi$  is well approximated by*

$$K_h(t, x, y) = h^{-\frac{3n}{2}} \int \int_{T^*U} A(t, x, y, z, \zeta, h) \exp \frac{i}{h} F(t, x, y, z, \zeta) dz d\zeta$$

for times  $|t| \leq c|\log h|$ . The phase reads

$$F = tp(z, \zeta) + \zeta^t \cdot W_{z^t}^x + \frac{1}{2} \left\langle \Gamma_{(z, \zeta)}^t W_{z^t}^x, W_{z^t}^x \right\rangle_{z^t} - \zeta \cdot W_z^y + \frac{1}{2} \left\langle \widetilde{\Gamma_{(z, \zeta)}^0} W_z^y, W_z^y \right\rangle_z$$

where  $\Gamma_{(z, \zeta)}^t : T_{z^t}M^{\mathbb{C}} \rightarrow T_{z^t}M^{\mathbb{C}}$  is complex linear, symmetric with positive definite imaginary part, and solves

$$\nabla_{\dot{z}^t} \Gamma_{(z, \zeta)}^t = -R_{z^t}(\cdot, \dot{z}^t) \dot{z}^t - (\Gamma_{(z, \zeta)}^t)^2, \quad \Gamma_{(z, \zeta)}^0 = i(g^{jk}(z)) + \text{real correction}$$

where  $R_{z^t}$  is the Riemann tensor at  $z^t$

## Relationship with the geodesic flow

**Theorem [Propagator approximation]** (B. 2016+) *Let  $\psi \in C^\infty(M)$  be supported in a coordinate patch  $U$ . The integral kernel of  $e^{-i\frac{t}{h}L_h}\chi(L_h)\psi$  is well approximated by*

$$K_h(t, x, y) = h^{-\frac{3n}{2}} \int \int_{T^*U} A(t, x, y, z, \zeta, h) \exp \frac{i}{h} F(t, x, y, z, \zeta) dz d\zeta$$

for times  $|t| \leq c|\log h|$ . The phase reads

$$F = tp(z, \zeta) + \zeta^t \cdot W_{z^t}^x + \frac{1}{2} \left\langle \Gamma_{(z, \zeta)}^t W_{z^t}^x, W_{z^t}^x \right\rangle_{z^t} - \zeta \cdot W_z^y + \frac{1}{2} \left\langle \widetilde{\Gamma_{(z, \zeta)}^0} W_z^y, W_z^y \right\rangle_z$$

where  $\Gamma_{(z, \zeta)}^t : T_{z^t}M^{\mathbb{C}} \rightarrow T_{z^t}M^{\mathbb{C}}$  is complex linear, symmetric with positive definite imaginary part, and solves

$$\nabla_{\dot{z}^t} \Gamma_{(z, \zeta)}^t = -R_{z^t}(\cdot, \dot{z}^t) \dot{z}^t - (\Gamma_{(z, \zeta)}^t)^2, \quad \Gamma_{(z, \zeta)}^0 = i(g^{jk}(z)) + \text{real correction}$$

where  $R_{z^t}$  is the Riemann tensor at  $z^t$

$$\left\langle \widetilde{\Gamma_{(z, \zeta)}^0} W_z^y, W_z^y \right\rangle_z = -\text{Re} \left\langle \Gamma_{(z, \zeta)}^0 W_z^y, W_z^y \right\rangle_z + i \text{Im} \left\langle \Gamma_{(z, \zeta)}^0 W_z^y, W_z^y \right\rangle_z$$

## Relationship with the geodesic flow

**Theorem [Propagator approximation]** (B. 2016+) *Let  $\psi \in C^\infty(M)$  be supported in a coordinate patch  $U$ . The integral kernel of  $e^{-i\frac{t}{h}L_h}\chi(L_h)\psi$  is well approximated by*

$$K_h(t, x, y) = h^{-\frac{3n}{2}} \int \int_{T^*U} A(t, x, y, z, \zeta, h) \exp \frac{i}{h} F(t, x, y, z, \zeta) dz d\zeta$$

for times  $|t| \leq c|\log h|$ . The phase reads

$$F = tp(z, \zeta) + \zeta^t \cdot W_{z^t}^x + \frac{1}{2} \left\langle \Gamma_{(z, \zeta)}^t W_{z^t}^x, W_{z^t}^x \right\rangle_{z^t} - \zeta \cdot W_z^y + \frac{1}{2} \left\langle \widetilde{\Gamma_{(z, \zeta)}^0} W_z^y, W_z^y \right\rangle_z$$

where  $\Gamma_{(z, \zeta)}^t : T_{z^t}M^{\mathbb{C}} \rightarrow T_{z^t}M^{\mathbb{C}}$  is complex linear, symmetric with positive definite imaginary part, and solves

$$\nabla_{\dot{z}^t} \Gamma_{(z, \zeta)}^t = -R_{z^t}(\cdot, \dot{z}^t) \dot{z}^t - (\Gamma_{(z, \zeta)}^t)^2, \quad \Gamma_{(z, \zeta)}^0 = i(g^{jk}(z)) + \text{real correction}$$

where  $R_{z^t}$  is the Riemann tensor at  $z^t$

$$\left\langle \widetilde{\Gamma_{(z, \zeta)}^0} W_z^y, W_z^y \right\rangle_z = -\text{Re} \left\langle \Gamma_{(z, \zeta)}^0 W_z^y, W_z^y \right\rangle_z + i \text{Im} \left\langle \Gamma_{(z, \zeta)}^0 W_z^y, W_z^y \right\rangle_z$$

and

$$y = \exp_z(W_z^y)$$

## Relationship with the geodesic flow

**Theorem [Propagator approximation]** (B. 2016+) *Let  $\psi \in C^\infty(M)$  be supported in a coordinate patch  $U$ . The integral kernel of  $e^{-i\frac{t}{h}L_h}\chi(L_h)\psi$  is well approximated by*

$$K_h(t, x, y) = h^{-\frac{3n}{2}} \int \int_{T^*U} A(t, x, y, z, \zeta, h) \exp \frac{i}{h} F(t, x, y, z, \zeta) dz d\zeta$$

for times  $|t| \leq c|\log h|$ . The phase reads

$$F = tp(z, \zeta) + \zeta^t \cdot W_{z^t}^x + \frac{1}{2} \left\langle \Gamma_{(z, \zeta)}^t W_{z^t}^x, W_{z^t}^x \right\rangle_{z^t} - \zeta \cdot W_z^y + \frac{1}{2} \left\langle \widetilde{\Gamma_{(z, \zeta)}^0} W_z^y, W_z^y \right\rangle_z$$

where  $\Gamma_{(z, \zeta)}^t : T_{z^t}M^{\mathbb{C}} \rightarrow T_{z^t}M^{\mathbb{C}}$  is complex linear, symmetric with positive definite imaginary part, and solves

$$\nabla_{\dot{z}^t} \Gamma_{(z, \zeta)}^t = -R_{z^t}(\cdot, \dot{z}^t) \dot{z}^t - (\Gamma_{(z, \zeta)}^t)^2, \quad \Gamma_{(z, \zeta)}^0 = i(g^{jk}(z)) + \text{real correction}$$

where  $R_{z^t}$  is the Riemann tensor at  $z^t$

$$\left\langle \widetilde{\Gamma_{(z, \zeta)}^0} W_z^y, W_z^y \right\rangle_z = -\text{Re} \left\langle \Gamma_{(z, \zeta)}^0 W_z^y, W_z^y \right\rangle_z + i \text{Im} \left\langle \Gamma_{(z, \zeta)}^0 W_z^y, W_z^y \right\rangle_z$$

and

$$y = \exp_z(W_z^y)$$

## An idea of proof: the Bérard bound

### Stationary phase in the integral

$$h^{-\frac{3n}{2}-1} \int \int \int \hat{\rho} \left( \frac{t}{T(h)} \right) A(t, x, x, z, \zeta, h) \exp \frac{i}{h} (t\lambda - F(t, x, x, z, \zeta)) dt dz d\zeta dv_g(x)$$

## An idea of proof: the Bérard bound

### Stationary phase in the integral

$$h^{-\frac{3n}{2}-1} \int \int \int \hat{\rho} \left( \frac{t}{T(h)} \right) A(t, x, x, z, \zeta, h) \exp \frac{i}{h} (t\lambda - F(t, x, x, z, \zeta)) dt dz d\zeta dv_g(x)$$

Critical points given by the conditions:

$$\operatorname{Im}(F(t, x, x, z, \zeta)) = 0 \quad \text{and} \quad d_{t,z,\zeta,x}(t\lambda - F(t, x, x, z, \zeta)) = 0$$



## An idea of proof: the Bérard bound

### Stationary phase in the integral

$$h^{-\frac{3n}{2}-1} \int \int \int \hat{\rho} \left( \frac{t}{T(h)} \right) A(t, x, x, z, \zeta, h) \exp \frac{i}{h} (t\lambda - F(t, x, x, z, \zeta)) dt dz d\zeta dv_g(x)$$

Critical points given by the conditions:

$$\operatorname{Im}(F(t, x, x, z, \zeta)) = 0 \quad \text{and} \quad d_{t,z,\zeta,x}(t\lambda - F(t, x, x, z, \zeta)) = 0$$

i.e.

$$z = z^t = x, \quad p(z, \zeta) = \lambda, \quad \zeta^t = \zeta$$

## An idea of proof: the Bérard bound

### Stationary phase in the integral

$$h^{-\frac{3n}{2}-1} \int \int \int \hat{\rho} \left( \frac{t}{T(h)} \right) A(t, x, x, z, \zeta, h) \exp \frac{i}{h} (t\lambda - F(t, x, x, z, \zeta)) dt dz d\zeta dv_g(x)$$

Critical points given by the conditions:

$$\operatorname{Im}(F(t, x, x, z, \zeta)) = 0 \quad \text{and} \quad d_{t,z,\zeta,x}(t\lambda - F(t, x, x, z, \zeta)) = 0$$

i.e.

$$z = z^t = x, \quad p(z, \zeta) = \lambda, \quad \zeta^t = \zeta$$

In particular

$$\phi^t(z, \zeta) = (z, \zeta)$$

## An idea of proof: the Bérard bound

### Stationary phase in the integral

$$h^{-\frac{3n}{2}-1} \int \int \int \hat{\rho} \left( \frac{t}{T(h)} \right) A(t, x, x, z, \zeta, h) \exp \frac{i}{h} (t\lambda - F(t, x, x, z, \zeta)) dt dz d\zeta dv_g(x)$$

Critical points given by the conditions:

$$\operatorname{Im}(F(t, x, x, z, \zeta)) = 0 \quad \text{and} \quad d_{t,z,\zeta,x}(t\lambda - F(t, x, x, z, \zeta)) = 0$$

i.e.

$$z = z^t = x, \quad \rho(z, \zeta) = \lambda, \quad \zeta^t = \zeta$$

In particular

$$\phi^t(z, \zeta) = (z, \zeta) \quad \implies \quad t \text{ is the length of a closed geodesic (if } \lambda = 1)$$

## An idea of proof: the Bérard bound

### Stationary phase in the integral

$$h^{-\frac{3n}{2}-1} \int \int \int \hat{\rho} \left( \frac{t}{T(h)} \right) A(t, x, x, z, \zeta, h) \exp \frac{i}{h} (t\lambda - F(t, x, x, z, \zeta)) dt dz d\zeta dv_g(x)$$

Critical points given by the conditions:

$$\operatorname{Im}(F(t, x, x, z, \zeta)) = 0 \quad \text{and} \quad d_{t,z,\zeta,x}(t\lambda - F(t, x, x, z, \zeta)) = 0$$

i.e.

$$z = z^t = x, \quad \rho(z, \zeta) = \lambda, \quad \zeta^t = \zeta$$

In particular

$$\phi^t(z, \zeta) = (z, \zeta) \quad \implies \quad t \text{ is the length of a closed geodesic (if } \lambda = 1)$$

Main contribution of the trivial period  $t = 0$ :

$$h^{-n} \int_{p=\lambda} d\Sigma$$

## An idea of proof: the Bérard bound

### Stationary phase in the integral

$$h^{-\frac{3n}{2}-1} \int \int \int \hat{\rho} \left( \frac{t}{T(h)} \right) A(t, x, x, z, \zeta, h) \exp \frac{i}{h} (t\lambda - F(t, x, x, z, \zeta)) dt dz d\zeta dv_g(x)$$

Critical points given by the conditions:

$$\operatorname{Im}(F(t, x, x, z, \zeta)) = 0 \quad \text{and} \quad d_{t,z,\zeta,x}(t\lambda - F(t, x, x, z, \zeta)) = 0$$

i.e.

$$z = z^t = x, \quad \rho(z, \zeta) = \lambda, \quad \zeta^t = \zeta$$

In particular

$$\phi^t(z, \zeta) = (z, \zeta) \quad \implies \quad t \text{ is the length of a closed geodesic (if } \lambda = 1)$$

Main contribution of the trivial period  $t = 0$ :

$$h^{-n} \int_{p=\lambda} d\Sigma = O(h^{-n})$$

## An idea of proof: the Bérard bound

### Stationary phase in the integral

$$h^{-\frac{3n}{2}-1} \int \int \int \hat{\rho} \left( \frac{t}{T(h)} \right) A(t, x, x, z, \zeta, h) \exp \frac{i}{h} (t\lambda - F(t, x, x, z, \zeta)) dt dz d\zeta dv_g(x)$$

Critical points given by the conditions:

$$\operatorname{Im}(F(t, x, x, z, \zeta)) = 0 \quad \text{and} \quad d_{t,z,\zeta,x}(t\lambda - F(t, x, x, z, \zeta)) = 0$$

i.e.

$$z = z^t = x, \quad \rho(z, \zeta) = \lambda, \quad \zeta^t = \zeta$$

In particular

$$\phi^t(z, \zeta) = (z, \zeta) \quad \implies \quad t \text{ is the length of a closed geodesic (if } \lambda = 1)$$

Main contribution of the trivial period  $t = 0$ :

$$h^{-n} \int_{p=\lambda} d\Sigma = O(h^{-n})$$

## An idea of proof: the Bérard bound

### Stationary phase in the integral

$$h^{-\frac{3n}{2}-1} \int \int \int \hat{\rho} \left( \frac{t}{T(h)} \right) A(t, x, x, z, \zeta, h) \exp \frac{i}{h} (t\lambda - F(t, x, x, z, \zeta)) dt dz d\zeta dv_g(x)$$

Critical points given by the conditions:

$$\operatorname{Im}(F(t, x, x, z, \zeta)) = 0 \quad \text{and} \quad d_{t,z,\zeta,x}(t\lambda - F(t, x, x, z, \zeta)) = 0$$

i.e.

$$z = z^t = x, \quad \rho(z, \zeta) = \lambda, \quad \zeta^t = \zeta$$

In particular

$$\phi^t(z, \zeta) = (z, \zeta) \quad \implies \quad t \text{ is the length of a closed geodesic (if } \lambda = 1)$$

Main contribution of the trivial period  $t = 0$ :

$$h^{-n} \int_{p=\lambda} d\Sigma = O(h^{-n})$$

Main contribution of non trivial periods: if  $T(h) = \varepsilon \log(1/h)$  it is (at most)  
 $O(h^{-\frac{n}{2}-O(\varepsilon)-1})$

## An idea of proof: the Bérard bound

### Stationary phase in the integral

$$h^{-\frac{3n}{2}-1} \int \int \int \hat{\rho} \left( \frac{t}{T(h)} \right) A(t, x, x, z, \zeta, h) \exp \frac{i}{h} (t\lambda - F(t, x, x, z, \zeta)) dt dz d\zeta dv_g(x)$$

Critical points given by the conditions:

$$\operatorname{Im}(F(t, x, x, z, \zeta)) = 0 \quad \text{and} \quad d_{t,z,\zeta,x}(t\lambda - F(t, x, x, z, \zeta)) = 0$$

i.e.

$$z = z^t = x, \quad p(z, \zeta) = \lambda, \quad \zeta^t = \zeta$$

In particular

$$\phi^t(z, \zeta) = (z, \zeta) \quad \implies \quad t \text{ is the length of a closed geodesic (if } \lambda = 1)$$

Main contribution of the trivial period  $t = 0$ :

$$h^{-n} \int_{p=\lambda} d\Sigma = O(h^{-n})$$

Main contribution of non trivial periods: if  $T(h) = \varepsilon \log(1/h)$  it is (at most)  
 $O(h^{-\frac{n}{2}-O(\varepsilon)-1}) \ll h^{-n}$  (at least if  $n \geq 3$ )



**Thank you for your attention**

**Thank you for your attention and happy new year**



## Some problems in spectral geometry

## Some problems in spectral geometry

Let  $(M, g)$  be a Riemannian manifold of dimension  $n$ . Let  $dv_g$  be the natural measure and  $\Delta_g$  be the Laplace-Beltrami operator, in local coordinates

$$dv_g = |g(x)| dx_1 \cdots dx_n, \quad |g(x)| := \det(g_{jk}(x))^{1/2} \quad (1)$$

$$\Delta_g = \sum_{j,k} g^{jk}(x) \partial_{x_j} \partial_{x_k} - \sum_{i,j,k} g^{jk}(x) \Gamma_{jk}^i(x) \partial_{x_i}$$

where  $(g^{jk}(x)) = (g_{jk}(x))^{-1}$  and  $g = \sum_{j,k} g_{jk}(x) dx_j dx_k$ .

**Proposition.** *If  $M$  is closed (and connected), there exists an orthonormal basis  $(\varphi_j)_{j \geq 0}$  of eigenfunctions of  $\Delta_g$*

$$-\Delta_g \varphi_j = \lambda_j \varphi_j, \quad 0 = \lambda_0 < \lambda_1 \leq \lambda_2 \leq \cdots, \quad \lambda_j \uparrow +\infty.$$

Simple example: on  $\mathbb{R}^n$

$$\chi(x)f(-h^2\Delta)v(x) = (2\pi h)^{-n} \int \int e^{\frac{i}{h}(x-y)\cdot\xi} \chi(x)f(|\xi|^2)v(y)dyd\xi$$

# Wave packets: a review of basic facts

## Wave packets: a review of basic facts

**Canonical example:** gaussian wave packet



## Wave packets: a review of basic facts

**Canonical example:** gaussian wave packet

$$\psi_{z,\zeta}(x) = \pi^{-\frac{n}{4}} \exp\left(i\zeta \cdot (x - z) - \frac{|x - z|^2}{2}\right), \quad x \in \mathbb{R}^n$$

## Wave packets: a review of basic facts

**Canonical example:** gaussian wave packet

$$\psi_{z,\zeta}(x) = \pi^{-\frac{n}{4}} \exp\left(i\zeta \cdot (x - z) - \frac{|x - z|^2}{2}\right), \quad x \in \mathbb{R}^n$$

localized around (or centered at)  $z$  in space,

## Wave packets: a review of basic facts

**Canonical example:** gaussian wave packet

$$\psi_{z,\zeta}(x) = \pi^{-\frac{n}{4}} \exp\left(i\zeta \cdot (x - z) - \frac{|x - z|^2}{2}\right), \quad x \in \mathbb{R}^n$$

localized around (or centered at)  $z$  in space, and near  $\zeta$  in momentum

$$(\mathcal{F}\psi_{z,\zeta})(\xi) = \pi^{-\frac{n}{4}} \exp\left(-iz \cdot \xi - \frac{|\xi - \zeta|^2}{2}\right)$$

## Wave packets: a review of basic facts

**Canonical example:** gaussian wave packet

$$\psi_{z,\zeta}(x) = \pi^{-\frac{n}{4}} \exp\left(i\zeta \cdot (x - z) - \frac{|x - z|^2}{2}\right), \quad x \in \mathbb{R}^n$$

localized around (or centered at)  $z$  in space, and near  $\zeta$  in momentum

$$(\mathcal{F}\psi_{z,\zeta})(\xi) = \pi^{-\frac{n}{4}} \exp\left(-iz \cdot \xi - \frac{|\xi - \zeta|^2}{2}\right)$$

Think of  $(z, \zeta)$  as a point in phase space  $T^*\mathbb{R}^n$

## Wave packets: a review of basic facts

**Canonical example:** gaussian wave packet

$$\psi_{z,\zeta}(x) = \pi^{-\frac{n}{4}} \exp\left(i\zeta \cdot (x - z) - \frac{|x - z|^2}{2}\right), \quad x \in \mathbb{R}^n$$

localized around (or centered at)  $z$  in space, and near  $\zeta$  in momentum

$$(\mathcal{F}\psi_{z,\zeta})(\xi) = \pi^{-\frac{n}{4}} \exp\left(-iz \cdot \xi - \frac{|\xi - \zeta|^2}{2}\right)$$

Think of  $(z, \zeta)$  as a point in phase space  $T^*\mathbb{R}^n$

We call  $\psi_{z,\zeta}$  a **Gaussian wave packet centered at**  $(z, \zeta)$

## Wave packets: a review of basic facts

**Canonical example:** gaussian wave packet

$$\psi_{z,\zeta}(x) = \pi^{-\frac{n}{4}} \exp\left(i\zeta \cdot (x - z) - \frac{|x - z|^2}{2}\right), \quad x \in \mathbb{R}^n$$

localized around (or centered at)  $z$  in space, and near  $\zeta$  in momentum

$$(\mathcal{F}\psi_{z,\zeta})(\xi) = \pi^{-\frac{n}{4}} \exp\left(-iz \cdot \xi - \frac{|\xi - \zeta|^2}{2}\right)$$

Think of  $(z, \zeta)$  as a point in phase space  $T^*\mathbb{R}^n$

We call  $\psi_{z,\zeta}$  a **Gaussian wave packet centered at**  $(z, \zeta)$

**Main interests (for us) :**

## Wave packets: a review of basic facts

**Canonical example:** gaussian wave packet

$$\psi_{z,\zeta}(x) = \pi^{-\frac{n}{4}} \exp\left(i\zeta \cdot (x - z) - \frac{|x - z|^2}{2}\right), \quad x \in \mathbb{R}^n$$

localized around (or centered at)  $z$  in space, and near  $\zeta$  in momentum

$$(\mathcal{F}\psi_{z,\zeta})(\xi) = \pi^{-\frac{n}{4}} \exp\left(-iz \cdot \xi - \frac{|\xi - \zeta|^2}{2}\right)$$

Think of  $(z, \zeta)$  as a point in phase space  $T^*\mathbb{R}^n$

We call  $\psi_{z,\zeta}$  a **Gaussian wave packet centered at**  $(z, \zeta)$

**Main interests (for us) :**

1. One can write "waves" (i.e. functions) as superposition of wave packets

## Wave packets: a review of basic facts

**Canonical example:** gaussian wave packet

$$\psi_{z,\zeta}(x) = \pi^{-\frac{n}{4}} \exp\left(i\zeta \cdot (x - z) - \frac{|x - z|^2}{2}\right), \quad x \in \mathbb{R}^n$$

localized around (or centered at)  $z$  in space, and near  $\zeta$  in momentum

$$(\mathcal{F}\psi_{z,\zeta})(\xi) = \pi^{-\frac{n}{4}} \exp\left(-iz \cdot \xi - \frac{|\xi - \zeta|^2}{2}\right)$$

Think of  $(z, \zeta)$  as a point in phase space  $T^*\mathbb{R}^n$

We call  $\psi_{z,\zeta}$  a **Gaussian wave packet centered at**  $(z, \zeta)$

**Main interests (for us) :**

1. One can write "waves" (i.e. functions) as superposition of wave packets
2. The evolution of a wave packet under a Schrödinger flow can be described rather explicitly (in a suitable regime)



# Wave packets: a review of basic facts

## 1. Wave packet decomposition

# Wave packets: a review of basic facts

## 1. Wave packet decomposition

Define the **Bargmann transform** of a function  $u$  by

$$Bu(z, \zeta) = \int_{\mathbb{R}^n} \overline{\psi_{z, \zeta}(x)} u(x) dx$$

# Wave packets: a review of basic facts

## 1. Wave packet decomposition

Define the **Bargmann transform** of a function  $u$  by

$$Bu(z, \zeta) = \int_{\mathbb{R}^n} \overline{\psi_{z, \zeta}(x)} u(x) dx$$

Then, one has the **inversion formula**

$$u = (2\pi)^{-n} B^* Bu$$

# Wave packets: a review of basic facts

## 1. Wave packet decomposition

Define the **Bargmann transform** of a function  $u$  by

$$Bu(z, \zeta) = \int_{\mathbb{R}^n} \overline{\psi_{z, \zeta}(x)} u(x) dx$$

Then, one has the **inversion formula**

$$u = (2\pi)^{-n} B^* Bu$$

In other words

$$u(x) = (2\pi)^{-n} \int \int_{T^*\mathbb{R}^n} (Bu)(z, \zeta) \psi_{z, \zeta}(x) dz d\zeta$$

is a decomposition of  $u$  as a (continuous) sum of wave packets

# Wave packets: a review of basic facts

## 2. Evolution of wave packets under the Schrödinger equation

# Wave packets: a review of basic facts

## 2. Evolution of wave packets under the Schrödinger equation

For quadratic potentials, one has exact formulas.

# Wave packets: a review of basic facts

## 2. Evolution of wave packets under the Schrödinger equation

For quadratic potentials, one has exact formulas. Set

$$p_\nu(x, \xi) = \frac{|\xi|^2}{2} + \nu \frac{|x|^2}{2}, \quad H_\nu = -\frac{\Delta}{2} + \nu \frac{|x|^2}{2}, \quad \nu = 0, +1, -1$$

## Wave packets: a review of basic facts

### 2. Evolution of wave packets under the Schrödinger equation

For quadratic potentials, one has exact formulas. Set

$$p_\nu(x, \xi) = \frac{|\xi|^2}{2} + \nu \frac{|x|^2}{2}, \quad H_\nu = -\frac{\Delta}{2} + \nu \frac{|x|^2}{2}, \quad \nu = 0, +1, -1$$

Then

$$e^{-itH_\nu} \psi_{z, \zeta}(x) = \pi^{-\frac{n}{4}} \gamma_\nu^t \exp i \left( S_\nu^t + \zeta_\nu^t \cdot (x - z_\nu^t) + \frac{\Gamma_\nu^t}{2} (x - z_\nu^t) \cdot (x - z_\nu^t) \right)$$



# Wave packets: a review of basic facts

## 2. Evolution of wave packets under the Schrödinger equation

For quadratic potentials, one has exact formulas. Set

$$p_\nu(x, \xi) = \frac{|\xi|^2}{2} + \nu \frac{|x|^2}{2}, \quad H_\nu = -\frac{\Delta}{2} + \nu \frac{|x|^2}{2}, \quad \nu = 0, +1, -1$$

Then

$$e^{-itH_\nu} \psi_{z, \zeta}(x) = \pi^{-\frac{n}{4}} \gamma_\nu^t \exp i \left( S_\nu^t + \zeta_\nu^t \cdot (x - z_\nu^t) + \frac{\Gamma_\nu^t}{2} (x - z_\nu^t) \cdot (x - z_\nu^t) \right)$$

where

$$(z_\nu^t, \zeta_\nu^t) = \Phi_{p_\nu}^t(z, \zeta), \quad S_\nu^t = \int_0^t \dot{z}_\nu^s \cdot \zeta_\nu^s - p_\nu(z_\nu^s, \zeta_\nu^s) ds$$

# Wave packets: a review of basic facts

## 2. Evolution of wave packets under the Schrödinger equation

For quadratic potentials, one has exact formulas. Set

$$p_\nu(x, \xi) = \frac{|\xi|^2}{2} + \nu \frac{|x|^2}{2}, \quad H_\nu = -\frac{\Delta}{2} + \nu \frac{|x|^2}{2}, \quad \nu = 0, +1, -1$$

Then

$$e^{-itH_\nu} \psi_{z, \zeta}(x) = \pi^{-\frac{n}{4}} \gamma_\nu^t \exp i \left( S_\nu^t + \zeta_\nu^t \cdot (x - z_\nu^t) + \frac{\Gamma_\nu^t}{2} (x - z_\nu^t) \cdot (x - z_\nu^t) \right)$$

where

$$(z_\nu^t, \zeta_\nu^t) = \Phi_{p_\nu}^t(z, \zeta), \quad S_\nu^t = \int_0^t \dot{z}_\nu^s \cdot \zeta_\nu^s - p_\nu(z_\nu^s, \zeta_\nu^s) ds$$

and  $\gamma_\nu^t, \Gamma_\nu^t$  are given in term of the differential of flow  $\Phi_{p_\nu}^t$ ,

$$D\Phi_{p_\nu}^t(z, \zeta) = \begin{pmatrix} A_\nu^t & B_\nu^t \\ C_\nu^t & D_\nu^t \end{pmatrix},$$

by

$$\Gamma_\nu^t = (C_\nu^t + iD_\nu^t)(A_\nu^t + iB_\nu^t)^{-1}, \quad \gamma_\nu^t = \det(A_\nu^t + iB_\nu^t)^{-1/2}.$$

## Wave packets: a review of basic facts

Explicitly, we obtain

$$\begin{aligned}\Gamma_0^t &= \frac{t+i}{1+t^2} l_n, & \gamma_0^t &= (1+it)^{-\frac{n}{2}} \\ \Gamma_1^t &= i l_n, & \gamma_1^t &= (\cos t + i \sin t)^{-\frac{n}{2}} \\ \Gamma_{-1}^t &= \frac{\sinh(2t) + i}{\cosh(2t)} l_n, & \gamma_{-1}^t &= (\cosh t + i \sinh t)^{-\frac{n}{2}}\end{aligned}$$

## Wave packets: a review of basic facts

Explicitly, we obtain

$$\begin{aligned}\Gamma_0^t &= \frac{t+i}{1+t^2} l_n, & \gamma_0^t &= (1+it)^{-\frac{n}{2}} \\ \Gamma_1^t &= i l_n, & \gamma_1^t &= (\cos t + i \sin t)^{-\frac{n}{2}} \\ \Gamma_{-1}^t &= \frac{\sinh(2t) + i}{\cosh(2t)} l_n, & \gamma_{-1}^t &= (\cosh t + i \sinh t)^{-\frac{n}{2}}\end{aligned}$$

This allows in particular to read the profile and spreading of the packets:

$$\begin{aligned}|e^{itH_0} \psi_{z,\zeta}(x)| &= \frac{1}{(\pi(1+t^2))^{\frac{n}{4}}} \exp\left(-\frac{|x-z_0^t|^2}{2(1+t^2)}\right) \\ |e^{itH_1} \psi_{z,\zeta}(x)| &= \frac{1}{\pi^{\frac{n}{4}}} \exp\left(-\frac{|x-z_1^t|^2}{2}\right) \\ |e^{itH_{-1}} \psi_{z,\zeta}(x)| &= \frac{1}{(\pi \cosh(2t))^{\frac{n}{4}}} \exp\left(-\frac{|x-z_{-1}^t|^2}{2 \cosh(2t)}\right)\end{aligned}$$

## Wave packets for semiclassical Schrödinger operators

From now on, we use a **semiclassical** normalization

$$\psi_{z,\zeta}^h(x) = (\pi h)^{-\frac{n}{4}} \exp\left(\frac{i}{h}\zeta \cdot (x - z) - \frac{|x - z|^2}{2h}\right)$$

## Wave packets for semiclassical Schrödinger operators

From now on, we use a **semiclassical** normalization

$$\psi_{z,\zeta}^h(x) = (\pi h)^{-\frac{n}{4}} \exp\left(\frac{i}{h}\zeta \cdot (x - z) - \frac{|x - z|^2}{2h}\right)$$

$\implies$  Localization around  $z$  on a scale  $h^{1/2}$

## Wave packets for semiclassical Schrödinger operators

From now on, we use a **semiclassical** normalization

$$\psi_{z,\zeta}^h(x) = (\pi h)^{-\frac{n}{4}} \exp\left(\frac{i}{h}\zeta \cdot (x - z) - \frac{|x - z|^2}{2h}\right)$$

⇒ Localization around  $z$  on a scale  $h^{1/2}$

Consider a semiclassical Schrödinger operator on  $\mathbb{R}^n$

$$H(h) = -\frac{h^2\Delta}{2} + V(x), \quad p(x, \xi) = \frac{|\xi|^2}{2} + V(x),$$

with  $V \in C^\infty(\mathbb{R}^n, \mathbb{R})$ .

# Wave packets for semiclassical Schrödinger operators

From now on, we use a **semiclassical** normalization

$$\psi_{z,\zeta}^h(x) = (\pi h)^{-\frac{n}{4}} \exp\left(\frac{i}{h}\zeta \cdot (x - z) - \frac{|x - z|^2}{2h}\right)$$

$\implies$  Localization around  $z$  on a scale  $h^{1/2}$

Consider a semiclassical Schrödinger operator on  $\mathbb{R}^n$

$$H(h) = -\frac{h^2\Delta}{2} + V(x), \quad p(x, \xi) = \frac{|\xi|^2}{2} + V(x),$$

with  $V \in C^\infty(\mathbb{R}^n, \mathbb{R})$ . Denote

$$(z^t, \zeta^t) = \Phi_p^t(z, \zeta), \quad \begin{pmatrix} A^t & B^t \\ C^t & D^t \end{pmatrix} := D\Phi_p^t(z, \zeta)$$



# Wave packets for semiclassical Schrödinger operators

From now on, we use a **semiclassical** normalization

$$\psi_{z,\zeta}^h(x) = (\pi h)^{-\frac{n}{4}} \exp\left(\frac{i}{h} \zeta \cdot (x - z) - \frac{|x - z|^2}{2h}\right)$$

⇒ Localization around  $z$  on a scale  $h^{1/2}$

Consider a semiclassical Schrödinger operator on  $\mathbb{R}^n$

$$H(h) = -\frac{h^2 \Delta}{2} + V(x), \quad p(x, \xi) = \frac{|\xi|^2}{2} + V(x),$$

with  $V \in C^\infty(\mathbb{R}^n, \mathbb{R})$ . Denote

$$(z^t, \zeta^t) = \Phi_p^t(z, \zeta), \quad \begin{pmatrix} A^t & B^t \\ C^t & D^t \end{pmatrix} := D\Phi_p^t(z, \zeta)$$

and

$$S^t = \int_0^t \dot{z}^s \cdot \zeta^s - p(z^s, \zeta^s) ds$$

# Wave packets for semiclassical Schrödinger operators

From now on, we use a **semiclassical** normalization

$$\psi_{z,\zeta}^h(x) = (\pi h)^{-\frac{n}{4}} \exp\left(\frac{i}{h}\zeta \cdot (x - z) - \frac{|x - z|^2}{2h}\right)$$

⇒ Localization around  $z$  on a scale  $h^{1/2}$

Consider a semiclassical Schrödinger operator on  $\mathbb{R}^n$

$$H(h) = -\frac{h^2\Delta}{2} + V(x), \quad p(x, \xi) = \frac{|\xi|^2}{2} + V(x),$$

with  $V \in C^\infty(\mathbb{R}^n, \mathbb{R})$ . Denote

$$(z^t, \zeta^t) = \Phi_p^t(z, \zeta), \quad \begin{pmatrix} A^t & B^t \\ C^t & D^t \end{pmatrix} := D\Phi_p^t(z, \zeta)$$

and

$$S^t = \int_0^t \dot{z}^s \cdot \zeta^s - p(z^s, \zeta^s) ds$$

**Proposition [action of the symplectic group on the Siegel half space]**

$A^t + iB^t$  is invertible and

$$\Gamma^t := (C^t + iD^t)(A^t + iB^t)^{-1}$$

is symmetric complex, with positive definite imaginary part

# Wave packets for semiclassical Schrödinger operators

**Theorem (Hagedorn-Joye, Combescure-Robert)** *In the limit  $\hbar \rightarrow 0$ , and under general conditions on  $V$ ,*

$$e^{-i\frac{t}{\hbar}H(\hbar)}\psi_{z,\zeta}^{\hbar}(x)$$

*is well approximated by*

$$(\pi\hbar)^{-\frac{n}{4}}\gamma^t\mathcal{A}_t^{\hbar}(x)\exp\frac{i}{\hbar}\left(S^t + \zeta^t \cdot (x - z^t) + \frac{\Gamma^t}{2}(x - z^t) \cdot (x - z^t)\right)$$

# Wave packets for semiclassical Schrödinger operators

**Theorem (Hagedorn-Joye, Combescure-Robert)** *In the limit  $h \rightarrow 0$ , and under general conditions on  $V$ ,*

$$e^{-i\frac{t}{h}H(h)}\psi_{z,\zeta}^h(x)$$

*is well approximated by*

$$(\pi h)^{-\frac{n}{4}}\gamma^t \mathcal{A}_t^h(x) \exp \frac{i}{h} \left( S^t + \zeta^t \cdot (x - z^t) + \frac{\Gamma^t}{2} (x - z^t) \cdot (x - z^t) \right)$$

*for times  $|t| \leq C_0 |\ln h|$  ( $C_0$  dynamical constant). Here  $\gamma^t = \det(A_t + iB_t)^{-1/2}$ .*

# Wave packets for semiclassical Schrödinger operators

**Theorem (Hagedorn-Joye, Combescure-Robert)** *In the limit  $\hbar \rightarrow 0$ , and under general conditions on  $V$ ,*

$$e^{-i\frac{t}{\hbar}H(\hbar)}\psi_{z,\zeta}^{\hbar}(x)$$

*is well approximated by*

$$(\pi\hbar)^{-\frac{n}{4}}\gamma^t\mathcal{A}_t^{\hbar}(x)\exp\frac{i}{\hbar}\left(S^t+\zeta^t\cdot(x-z^t)+\frac{\Gamma^t}{2}(x-z^t)\cdot(x-z^t)\right)$$

*for times  $|t| \leq C_0|\ln \hbar|$  ( $C_0$  dynamical constant). Here  $\gamma^t = \det(A_t + iB_t)^{-1/2}$ . The amplitude is of the form*

$$\mathcal{A}_t^{\hbar}(x) \sim 1 + \sum_{j \geq 1} \hbar^{\frac{j}{2}} A_j\left(z, \zeta, t, \frac{x - z^t}{\hbar^{\frac{1}{2}}}\right)$$

*with  $A_j(z, \zeta, t, X)$  polynomial of degree  $\leq 3j$  in  $X$ , with coeff. depending on the classical trajectory  $t \mapsto (z^t, \zeta^t)$  and the Taylor expansion of  $V$  at  $z^t$*

# Wave packets for semiclassical Schrödinger operators

**Theorem (Hagedorn-Joye, Combescure-Robert)** *In the limit  $\hbar \rightarrow 0$ , and under general conditions on  $V$ ,*

$$e^{-i\frac{t}{\hbar}H(\hbar)}\psi_{z,\zeta}^{\hbar}(x)$$

*is well approximated by*

$$(\pi\hbar)^{-\frac{n}{4}}\gamma^t\mathcal{A}_t^{\hbar}(x)\exp\frac{i}{\hbar}\left(S^t+\zeta^t\cdot(x-z^t)+\frac{\Gamma^t}{2}(x-z^t)\cdot(x-z^t)\right)$$

*for times  $|t| \leq C_0|\ln \hbar|$  ( $C_0$  dynamical constant). Here  $\gamma^t = \det(A_t + iB_t)^{-1/2}$ . The amplitude is of the form*

$$\mathcal{A}_t^{\hbar}(x) \sim 1 + \sum_{j \geq 1} \hbar^{\frac{j}{2}} A_j\left(z, \zeta, t, \frac{x - z^t}{\hbar^{\frac{1}{2}}}\right)$$

*with  $A_j(z, \zeta, t, X)$  polynomial of degree  $\leq 3j$  in  $X$ , with coeff. depending on the classical trajectory  $t \mapsto (z^t, \zeta^t)$  and the Taylor expansion of  $V$  at  $z^t$*

**Rem.** The polynomial growth of the amplitude in  $(x - z^t)/\hbar^{\frac{1}{2}}$  is beaten by the exponential decay of the exponential since  $\text{Im}(\Gamma^t)$  is positive definite

# Wave packets for semiclassical Schrödinger operators

**Theorem (Hagedorn-Joye, Combescure-Robert)** *In the limit  $\hbar \rightarrow 0$ , and under general conditions on  $V$ ,*

$$e^{-i\frac{t}{\hbar}H(\hbar)}\psi_{z,\zeta}^{\hbar}(x)$$

*is well approximated by*

$$(\pi\hbar)^{-\frac{n}{4}}\gamma^t\mathcal{A}_t^{\hbar}(x)\exp\frac{i}{\hbar}\left(S^t+\zeta^t\cdot(x-z^t)+\frac{\Gamma^t}{2}(x-z^t)\cdot(x-z^t)\right)$$

*for times  $|t| \leq C_0|\ln \hbar|$  ( $C_0$  dynamical constant). Here  $\gamma^t = \det(A_t + iB_t)^{-1/2}$ . The amplitude is of the form*

$$\mathcal{A}_t^{\hbar}(x) \sim 1 + \sum_{j \geq 1} \hbar^{\frac{j}{2}} A_j\left(z, \zeta, t, \frac{x - z^t}{\hbar^{\frac{1}{2}}}\right)$$

*with  $A_j(z, \zeta, t, X)$  polynomial of degree  $\leq 3j$  in  $X$ , with coeff. depending on the classical trajectory  $t \mapsto (z^t, \zeta^t)$  and the Taylor expansion of  $V$  at  $z^t$*

**Rem.** The polynomial growth of the amplitude in  $(x - z^t)/\hbar^{\frac{1}{2}}$  is beaten by the exponential decay of the exponential since  $\text{Im}(\Gamma^t)$  is positive definite

$\Rightarrow$  Concentration near the classical trajectory,

# Wave packets for semiclassical Schrödinger operators

**Theorem (Hagedorn-Joye, Combes-Robert)** *In the limit  $h \rightarrow 0$ , and under general conditions on  $V$ ,*

$$e^{-i\frac{t}{h}H(h)}\psi_{z,\zeta}^h(x)$$

*is well approximated by*

$$(\pi h)^{-\frac{n}{4}}\gamma^t \mathcal{A}_t^h(x) \exp \frac{i}{h} \left( S^t + \zeta^t \cdot (x - z^t) + \frac{\Gamma^t}{2} (x - z^t) \cdot (x - z^t) \right)$$

*for times  $|t| \leq C_0 |\ln h|$  ( $C_0$  dynamical constant). Here  $\gamma^t = \det(A_t + iB_t)^{-1/2}$ . The amplitude is of the form*

$$\mathcal{A}_t^h(x) \sim 1 + \sum_{j \geq 1} h^{\frac{j}{2}} A_j \left( z, \zeta, t, \frac{x - z^t}{h^{\frac{1}{2}}} \right)$$

*with  $A_j(z, \zeta, t, X)$  polynomial of degree  $\leq 3j$  in  $X$ , with coeff. depending on the classical trajectory  $t \mapsto (z^t, \zeta^t)$  and the Taylor expansion of  $V$  at  $z^t$*

**Rem.** The polynomial growth of the amplitude in  $(x - z^t)/h^{\frac{1}{2}}$  is beaten by the exponential decay of the exponential since  $\text{Im}(\Gamma^t)$  is positive definite

$\Rightarrow$  Concentration near the classical trajectory, at least as long as  $\text{Im}(\Gamma^t) \gg h$



## Wave packets in semiclassical analysis

**Sketch of proof.**

**Lemma** *The matrix  $\Gamma^t$  satisfies the Riccati equation*

$$\dot{\Gamma}^t = -V^{(2)}(z^t) - (\Gamma^t)^2, \quad \Gamma^0 = iI_n,$$

## Wave packets in semiclassical analysis

**Sketch of proof.**

**Lemma** *The matrix  $\Gamma^t$  satisfies the Riccati equation*

$$\dot{\Gamma}^t = -V^{(2)}(z^t) - (\Gamma^t)^2, \quad \Gamma^0 = iI_n,$$

*and the function  $\gamma^t$  satisfies*

$$\dot{\gamma}^t = -\frac{\text{tr}(\Gamma^t)}{2}\gamma^t.$$

# Wave packets in semiclassical analysis

**Sketch of proof.**

**Lemma** *The matrix  $\Gamma^t$  satisfies the Riccati equation*

$$\dot{\Gamma}^t = -V^{(2)}(z^t) - (\Gamma^t)^2, \quad \Gamma^0 = iI_n,$$

*and the function  $\gamma^t$  satisfies*

$$\dot{\gamma}^t = -\frac{\text{tr}(\Gamma^t)}{2}\gamma^t.$$

Set

$$\varphi := S^t + \zeta^t \cdot (x - z^t) + \frac{\Gamma^t}{2}(x - z^t) \cdot (x - z^t).$$

# Wave packets in semiclassical analysis

**Sketch of proof.**

**Lemma** *The matrix  $\Gamma^t$  satisfies the Riccati equation*

$$\dot{\Gamma}^t = -V^{(2)}(z^t) - (\Gamma^t)^2, \quad \Gamma^0 = iI_n,$$

*and the function  $\gamma^t$  satisfies*

$$\dot{\gamma}^t = -\frac{\text{tr}(\Gamma^t)}{2}\gamma^t.$$

Set

$$\varphi := S^t + \zeta^t \cdot (x - z^t) + \frac{\Gamma^t}{2}(x - z^t) \cdot (x - z^t).$$

Then

$$H(h)\gamma^t e^{\frac{i}{h}\varphi}$$

# Wave packets in semiclassical analysis

**Sketch of proof.**

**Lemma** *The matrix  $\Gamma^t$  satisfies the Riccati equation*

$$\dot{\Gamma}^t = -V^{(2)}(z^t) - (\Gamma^t)^2, \quad \Gamma^0 = iI_n,$$

*and the function  $\gamma^t$  satisfies*

$$\dot{\gamma}^t = -\frac{\text{tr}(\Gamma^t)}{2}\gamma^t.$$

Set

$$\varphi := S^t + \zeta^t \cdot (x - z^t) + \frac{\Gamma^t}{2}(x - z^t) \cdot (x - z^t).$$

Then

$$H(h)\gamma^t e^{\frac{i}{h}\varphi} = \left[ \left( \dot{\varphi} + \frac{\nabla_x \varphi \cdot \nabla_x \varphi}{2} + V(x) \right) - ih \left( \frac{\dot{\gamma}^t}{\gamma^t} + \frac{\Delta \varphi}{2} \right) \right] \gamma^t e^{\frac{i}{h}\varphi}$$

## Wave packets in semiclassical analysis

**Sketch of proof.**

**Lemma** *The matrix  $\Gamma^t$  satisfies the Riccati equation*

$$\dot{\Gamma}^t = -V^{(2)}(z^t) - (\Gamma^t)^2, \quad \Gamma^0 = iI_n,$$

*and the function  $\gamma^t$  satisfies*

$$\dot{\gamma}^t = -\frac{\text{tr}(\Gamma^t)}{2}\gamma^t.$$

Set

$$\varphi := S^t + \zeta^t \cdot (x - z^t) + \frac{\Gamma^t}{2}(x - z^t) \cdot (x - z^t).$$

Then

$$\begin{aligned} H(h)\gamma^t e^{\frac{i}{h}\varphi} &= \left[ \left( \dot{\varphi} + \frac{\nabla_x \varphi \cdot \nabla_x \varphi}{2} + V(x) \right) - ih \left( \frac{\dot{\gamma}^t}{\gamma^t} + \frac{\Delta \varphi}{2} \right) \right] \gamma^t e^{\frac{i}{h}\varphi} \\ &= \left[ V(x) - V(z^t) - V^{(1)}(z^t) \cdot (x - z^t) - \frac{V^{(2)}(z^t)}{2}(x - z^t) \cdot (x - z^t) \right] \gamma^t e^{\frac{i}{h}\varphi} \end{aligned}$$

# Wave packets in semiclassical analysis

**Sketch of proof.**

**Lemma** *The matrix  $\Gamma^t$  satisfies the Riccati equation*

$$\dot{\Gamma}^t = -V^{(2)}(z^t) - (\Gamma^t)^2, \quad \Gamma^0 = iI_n,$$

*and the function  $\gamma^t$  satisfies*

$$\dot{\gamma}^t = -\frac{\text{tr}(\Gamma^t)}{2}\gamma^t.$$

Set

$$\varphi := S^t + \zeta^t \cdot (x - z^t) + \frac{\Gamma^t}{2}(x - z^t) \cdot (x - z^t).$$

Then

$$\begin{aligned} H(h)\gamma^t e^{\frac{i}{h}\varphi} &= \left[ \left( \dot{\varphi} + \frac{\nabla_x \varphi \cdot \nabla_x \varphi}{2} + V(x) \right) - ih \left( \frac{\dot{\gamma}^t}{\gamma^t} + \frac{\Delta \varphi}{2} \right) \right] \gamma^t e^{\frac{i}{h}\varphi} \\ &= \left[ V(x) - V(z^t) - V^{(1)}(z^t) \cdot (x - z^t) - \frac{V^{(2)}(z^t)}{2}(x - z^t) \cdot (x - z^t) \right] \gamma^t e^{\frac{i}{h}\varphi} \\ &= O(|x - z^t|^3) \gamma^t e^{\frac{i}{h}\varphi} \end{aligned}$$

# Wave packets in semiclassical analysis

Sketch of proof.

**Lemma** The matrix  $\Gamma^t$  satisfies the Riccati equation

$$\dot{\Gamma}^t = -V^{(2)}(z^t) - (\Gamma^t)^2, \quad \Gamma^0 = iI_n,$$

and the function  $\gamma^t$  satisfies

$$\dot{\gamma}^t = -\frac{\text{tr}(\Gamma^t)}{2}\gamma^t.$$

Set

$$\varphi := S^t + \zeta^t \cdot (x - z^t) + \frac{\Gamma^t}{2}(x - z^t) \cdot (x - z^t).$$

Then

$$\begin{aligned} H(h)\gamma^t e^{\frac{i}{h}\varphi} &= \left[ \left( \dot{\varphi} + \frac{\nabla_x \varphi \cdot \nabla_x \varphi}{2} + V(x) \right) - ih \left( \frac{\dot{\gamma}^t}{\gamma^t} + \frac{\Delta \varphi}{2} \right) \right] \gamma^t e^{\frac{i}{h}\varphi} \\ &= \left[ V(x) - V(z^t) - V^{(1)}(z^t) \cdot (x - z^t) - \frac{V^{(2)}(z^t)}{2}(x - z^t) \cdot (x - z^t) \right] \gamma^t e^{\frac{i}{h}\varphi} \\ &= O(|x - z^t|^3) \gamma^t e^{\frac{i}{h}\varphi} \\ &= h^{3/2} O\left(\frac{|x - z^t|^3}{h^{3/2}}\right) \gamma^t e^{\frac{i}{h}\varphi} \end{aligned}$$



# Wave packets on Riemannian manifolds

# Wave packets on Riemannian manifolds

**Goal:** to emulate the construction on  $\mathbb{R}^n$

# Wave packets on Riemannian manifolds

**Goal:** to emulate the construction on  $\mathbb{R}^n$

Previous related works:

- ▶ **Construction of quasimodes:** by propagating a *single* wave packet along a *closed* geodesic (Babich-Lazutkin, Ralston, Paul-Uribe, Nonnenmacher-Eswarathan...).

# Wave packets on Riemannian manifolds

**Goal:** to emulate the construction on  $\mathbb{R}^n$

Previous related works:

- ▶ **Construction of quasimodes:** by propagating a *single* wave packet along a *closed* geodesic (Babich-Lazutkin, Ralston, Paul-Uribe, Nonnenmacher-Eswarathan...). Allows to use Fermi coordinates.

# Wave packets on Riemannian manifolds

**Goal:** to emulate the construction on  $\mathbb{R}^n$

Previous related works:

- ▶ **Construction of quasimodes:** by propagating a *single* wave packet along a *closed* geodesic (Babich-Lazutkin, Ralston, Paul-Uribe, Nonnenmacher-Eswarathan...). Allows to use Fermi coordinates.
- ▶ **More general propagation results:** Paul-Uribe, Guillemin-Uribe-Wang: qualitative description of wave packets and their evolutions (for Hamiltonians with non homogeneous symbols).

# Wave packets on Riemannian manifolds

**Goal:** to emulate the construction on  $\mathbb{R}^n$

Previous related works:

- ▶ **Construction of quasimodes:** by propagating a *single* wave packet along a *closed* geodesic (Babich-Lazutkin, Ralston, Paul-Uribe, Nonnenmacher-Eswarathan...). Allows to use Fermi coordinates.
- ▶ **More general propagation results:** Paul-Uribe, Guillemin-Urbe-Wang: qualitative description of wave packets and their evolutions (for Hamiltonians with non homogeneous symbols). General but not so explicit,

# Wave packets on Riemannian manifolds

**Goal:** to emulate the construction on  $\mathbb{R}^n$

Previous related works:

- ▶ **Construction of quasimodes:** by propagating a *single* wave packet along a *closed* geodesic (Babich-Lazutkin, Ralston, Paul-Urbe, Nonnenmacher-Eswarathan...). Allows to use Fermi coordinates.
- ▶ **More general propagation results:** Paul-Urbe, Guillemin-Urbe-Wang: qualitative description of wave packets and their evolutions (for Hamiltonians with non homogeneous symbols). General but not so explicit, using local coordinates

# Wave packets on Riemannian manifolds

**Goal:** to emulate the construction on  $\mathbb{R}^n$

Previous related works:

- ▶ **Construction of quasimodes:** by propagating a *single* wave packet along a *closed* geodesic (Babich-Lazutkin, Ralston, Paul-Urbe, Nonnenmacher-Eswarathan...). Allows to use Fermi coordinates.
- ▶ **More general propagation results:** Paul-Urbe, Guillemin-Urbe-Wang: qualitative description of wave packets and their evolutions (for Hamiltonians with non homogeneous symbols). General but not so explicit, using local coordinates and given for finite times



# Wave packets on Riemannian manifolds

**Goal:** to emulate the construction on  $\mathbb{R}^n$

Previous related works:

- ▶ **Construction of quasimodes:** by propagating a *single* wave packet along a *closed* geodesic (Babich-Lazutkin, Ralston, Paul-Uribe, Nonnenmacher-Eswarathan...). Allows to use Fermi coordinates.
- ▶ **More general propagation results:** Paul-Uribe, Guillemin-Urbe-Wang: qualitative description of wave packets and their evolutions (for Hamiltonians with non homogeneous symbols). General but not so explicit, using local coordinates and given for finite times

**Motivations and interests:**

1. Consider more than the propagation along a *single* trajectory  $\Rightarrow$  vary  $(z, \zeta)$

# Wave packets on Riemannian manifolds

**Goal:** to emulate the construction on  $\mathbb{R}^n$

Previous related works:

- ▶ **Construction of quasimodes:** by propagating a *single* wave packet along a *closed* geodesic (Babich-Lazutkin, Ralston, Paul-Uribe, Nonnenmacher-Eswarathan...). Allows to use Fermi coordinates.
- ▶ **More general propagation results:** Paul-Uribe, Guillemin-Uribe-Wang: qualitative description of wave packets and their evolutions (for Hamiltonians with non homogeneous symbols). General but not so explicit, using local coordinates and given for finite times

**Motivations and interests:**

1. Consider more than the propagation along a *single* trajectory  $\Rightarrow$  vary  $(z, \zeta)$
2. Get an (at most as possible) intrinsic description of wave packets propagation

# Wave packets on Riemannian manifolds

**Goal:** to emulate the construction on  $\mathbb{R}^n$

Previous related works:

- ▶ **Construction of quasimodes:** by propagating a *single* wave packet along a *closed* geodesic (Babich-Lazutkin, Ralston, Paul-Uribe, Nonnenmacher-Eswarathanan...). Allows to use Fermi coordinates.
- ▶ **More general propagation results:** Paul-Uribe, Guillemin-Uribe-Wang: qualitative description of wave packets and their evolutions (for Hamiltonians with non homogeneous symbols). General but not so explicit, using local coordinates and given for finite times

**Motivations and interests:**

1. Consider more than the propagation along a *single* trajectory  $\Rightarrow$  vary  $(z, \zeta)$
2. Get an (at most as possible) intrinsic description of wave packets propagation
3. Get (relatively) explicit approximation of  $e^{i\hbar H(h)/\hbar}$  as a single integral

# Wave packets on Riemannian manifolds

**Goal:** to emulate the construction on  $\mathbb{R}^n$

Previous related works:

- ▶ **Construction of quasimodes:** by propagating a *single* wave packet along a *closed* geodesic (Babich-Lazutkin, Ralston, Paul-Uribe, Nonnenmacher-Eswarathanan...). Allows to use Fermi coordinates.
- ▶ **More general propagation results:** Paul-Uribe, Guillemin-Uribe-Wang: qualitative description of wave packets and their evolutions (for Hamiltonians with non homogeneous symbols). General but not so explicit, using local coordinates and given for finite times

**Motivations and interests:**

1. Consider more than the propagation along a *single* trajectory  $\Rightarrow$  vary  $(z, \zeta)$
2. Get an (at most as possible) intrinsic description of wave packets propagation
3. Get (relatively) explicit approximation of  $e^{itH(h)/h}$  as a single integral, without need to go to the universal cover

# Wave packets on Riemannian manifolds

**Goal:** to emulate the construction on  $\mathbb{R}^n$

Previous related works:

- ▶ **Construction of quasimodes:** by propagating a *single* wave packet along a *closed* geodesic (Babich-Lazutkin, Ralston, Paul-Uribe, Nonnenmacher-Eswarathanan...). Allows to use Fermi coordinates.
- ▶ **More general propagation results:** Paul-Uribe, Guillemin-Uribe-Wang: qualitative description of wave packets and their evolutions (for Hamiltonians with non homogeneous symbols). General but not so explicit, using local coordinates and given for finite times

**Motivations and interests:**

1. Consider more than the propagation along a *single* trajectory  $\Rightarrow$  vary  $(z, \zeta)$
2. Get an (at most as possible) intrinsic description of wave packets propagation
3. Get (relatively) explicit approximation of  $e^{itH(h)/h}$  as a single integral, without need to go to the universal cover, up to  $|t| \leq C_0 |\log h|$

# Wave packets on Riemannian manifolds

**Goal:** to emulate the construction on  $\mathbb{R}^n$

Previous related works:

- ▶ **Construction of quasimodes:** by propagating a *single* wave packet along a *closed* geodesic (Babich-Lazutkin, Ralston, Paul-Uribe, Nonnenmacher-Eswarathanan...). Allows to use Fermi coordinates.
- ▶ **More general propagation results:** Paul-Uribe, Guillemin-Urbe-Wang: qualitative description of wave packets and their evolutions (for Hamiltonians with non homogeneous symbols). General but not so explicit, using local coordinates and given for finite times

**Motivations and interests:**

1. Consider more than the propagation along a *single* trajectory  $\Rightarrow$  vary  $(z, \zeta)$
2. Get an (at most as possible) intrinsic description of wave packets propagation
3. Get (relatively) explicit approximation of  $e^{itH(h)/h}$  as a single integral, without need to go to the universal cover, up to  $|t| \leq C_0 |\log h|$
4. See e.g. quite explicitly the effect of (negative) curvature

# Wave packets on Riemannian manifolds

**Goal:** to emulate the construction on  $\mathbb{R}^n$

Previous related works:

- ▶ **Construction of quasimodes:** by propagating a *single* wave packet along a *closed* geodesic (Babich-Lazutkin, Ralston, Paul-Uribe, Nonnenmacher-Eswarathanan...). Allows to use Fermi coordinates.
- ▶ **More general propagation results:** Paul-Uribe, Guillemin-Uribe-Wang: qualitative description of wave packets and their evolutions (for Hamiltonians with non homogeneous symbols). General but not so explicit, using local coordinates and given for finite times

**Motivations and interests:**

1. Consider more than the propagation along a *single* trajectory  $\Rightarrow$  vary  $(z, \zeta)$
2. Get an (at most as possible) intrinsic description of wave packets propagation
3. Get (relatively) explicit approximation of  $e^{itH(h)/h}$  as a single integral, without need to go to the universal cover, up to  $|t| \leq C_0 |\log h|$
4. See e.g. quite explicitly the effect of (negative) curvature
5. ...

# Wave packets on Riemannian manifolds

Let  $(M^n, g)$  be a Riemannian manifold with **bounded geometry**



## Wave packets on Riemannian manifolds

Let  $(M^n, g)$  be a Riemannian manifold with **bounded geometry** i.e.

1. injectivity radius bounded from below by  $r_0 > 0$

## Wave packets on Riemannian manifolds

Let  $(M^n, g)$  be a Riemannian manifold with **bounded geometry** i.e.

1. injectivity radius bounded from below by  $r_0 > 0$
2. all covariant derivatives of the Riemann curvature tensor bounded on  $M$

# Wave packets on Riemannian manifolds

Let  $(M^n, g)$  be a Riemannian manifold with **bounded geometry** i.e.

1. injectivity radius bounded from below by  $r_0 > 0$
2. all covariant derivatives of the Riemann curvature tensor bounded on  $M$
3. complete (for simplicity)

## Wave packets on Riemannian manifolds

Let  $(M^n, g)$  be a Riemannian manifold with **bounded geometry** i.e.

1. injectivity radius bounded from below by  $r_0 > 0$
2. all covariant derivatives of the Riemann curvature tensor bounded on  $M$
3. complete (for simplicity)

**Example.** Any closed Riemannian manifold

# Wave packets on Riemannian manifolds

Let  $(M^n, g)$  be a Riemannian manifold with **bounded geometry** i.e.

1. injectivity radius bounded from below by  $r_0 > 0$
2. all covariant derivatives of the Riemann curvature tensor bounded on  $M$
3. complete (for simplicity)

**Example.** Any closed Riemannian manifold

**Lemma [Inverse exponential map close to the diagonal of  $M \times M$ ]** *If  $d_g(z, m) < r_0$ , there is a unique  $W_z^m \in T_z M$  such that*

$$m = \exp_z (W_z^m).$$

# Wave packets on Riemannian manifolds

Let  $(M^n, g)$  be a Riemannian manifold with **bounded geometry** i.e.

1. injectivity radius bounded from below by  $r_0 > 0$
2. all covariant derivatives of the Riemann curvature tensor bounded on  $M$
3. complete (for simplicity)

**Example.** Any closed Riemannian manifold

**Lemma [Inverse exponential map close to the diagonal of  $M \times M$ ]** *If  $d_g(z, m) < r_0$ , there is a unique  $W_z^m \in T_z M$  such that*

$$m = \exp_z (W_z^m).$$

*For fixed  $m$ ,  $z \mapsto W_z^m$  is a vector field*

## Wave packets on Riemannian manifolds

Let  $(M^n, g)$  be a Riemannian manifold with **bounded geometry** i.e.

1. injectivity radius bounded from below by  $r_0 > 0$
2. all covariant derivatives of the Riemann curvature tensor bounded on  $M$
3. complete (for simplicity)

**Example.** Any closed Riemannian manifold

**Lemma [Inverse exponential map close to the diagonal of  $M \times M$ ]** *If  $d_g(z, m) < r_0$ , there is a unique  $W_z^m \in T_z M$  such that*

$$m = \exp_z (W_z^m).$$

*For fixed  $m$ ,  $z \mapsto W_z^m$  is a vector field and one can expand its covariant derivative*

$$\nabla W_z^m \sim -I + \frac{1}{3}R_z(\cdot, W_z^m)W_z^m + \frac{1}{12}(\nabla R)_z(W_z^m; \cdot, W_z^m)W_z^m + \dots$$

## Wave packets on Riemannian manifolds

Let  $(M^n, g)$  be a Riemannian manifold with **bounded geometry** i.e.

1. injectivity radius bounded from below by  $r_0 > 0$
2. all covariant derivatives of the Riemann curvature tensor bounded on  $M$
3. complete (for simplicity)

**Example.** Any closed Riemannian manifold

**Lemma [Inverse exponential map close to the diagonal of  $M \times M$ ]** *If  $d_g(z, m) < r_0$ , there is a unique  $W_z^m \in T_z M$  such that*

$$m = \exp_z (W_z^m).$$

*For fixed  $m$ ,  $z \mapsto W_z^m$  is a vector field and one can expand its covariant derivative*

$$\nabla W_z^m \sim -I + \frac{1}{3}R_z(\cdot, W_z^m)W_z^m + \frac{1}{12}(\nabla R)_z(W_z^m; \cdot, W_z^m)W_z^m + \dots$$

*All tensors in this expansion are bounded*



## Wave packets on Riemannian manifolds

Let  $(M^n, g)$  be a Riemannian manifold with **bounded geometry** i.e.

1. injectivity radius bounded from below by  $r_0 > 0$
2. all covariant derivatives of the Riemann curvature tensor bounded on  $M$
3. complete (for simplicity)

**Example.** Any closed Riemannian manifold

**Lemma [Inverse exponential map close to the diagonal of  $M \times M$ ]** *If  $d_g(z, m) < r_0$ , there is a unique  $W_z^m \in T_z M$  such that*

$$m = \exp_z (W_z^m).$$

*For fixed  $m$ ,  $z \mapsto W_z^m$  is a vector field and one can expand its covariant derivative*

$$\nabla W_z^m \sim -I + \frac{1}{3}R_z(\cdot, W_z^m)W_z^m + \frac{1}{12}(\nabla R)_z(W_z^m; \cdot, W_z^m)W_z^m + \dots$$

*All tensors in this expansion are bounded (similar result for higher covariant derivatives)*

## Wave packets on Riemannian manifolds

Let  $(M^n, g)$  be a Riemannian manifold with **bounded geometry** i.e.

1. injectivity radius bounded from below by  $r_0 > 0$
2. all covariant derivatives of the Riemann curvature tensor bounded on  $M$
3. complete (for simplicity)

**Example.** Any closed Riemannian manifold

**Lemma [Inverse exponential map close to the diagonal of  $M \times M$ ]** *If  $d_g(z, m) < r_0$ , there is a unique  $W_z^m \in T_z M$  such that*

$$m = \exp_z (W_z^m).$$

*For fixed  $m$ ,  $z \mapsto W_z^m$  is a vector field and one can expand its covariant derivative*

$$\nabla W_z^m \sim -I + \frac{1}{3}R_z(\cdot, W_z^m)W_z^m + \frac{1}{12}(\nabla R)_z(W_z^m; \cdot, W_z^m)W_z^m + \dots$$

*All tensors in this expansion are bounded (similar result for higher covariant derivatives)*

**Rem:** on  $\mathbb{R}^n$ ,  $W_z^m = m - z$ .

## Wave packets on Riemannian manifolds

Consider  $V \in C^\infty(M, \mathbb{R})$  and

$$H(h) := -h^2 \frac{\Delta_g}{2} + V$$

## Wave packets on Riemannian manifolds

Consider  $V \in C^\infty(M, \mathbb{R})$  and

$$H(h) := -h^2 \frac{\Delta_g}{2} + V$$

$$(z^t, \zeta^t) = \Phi^t(z, \zeta), \quad \text{Hamiltonian flow of } \frac{|\xi|_m^2}{2} + V(m)$$

## Wave packets on Riemannian manifolds

Consider  $V \in C^\infty(M, \mathbb{R})$  and

$$H(h) := -h^2 \frac{\Delta_g}{2} + V$$

$$(z^t, \zeta^t) = \Phi^t(z, \zeta), \quad \text{Hamiltonian flow of } \frac{|\xi|_m^2}{2} + V(m)$$

**Proposition.** *Let  $U$  be a coordinate patch, with coordinates  $y_1, \dots, y_n$ .*

## Wave packets on Riemannian manifolds

Consider  $V \in C^\infty(M, \mathbb{R})$  and

$$H(h) := -h^2 \frac{\Delta_g}{2} + V$$

$$(z^t, \zeta^t) = \Phi^t(z, \zeta), \quad \text{Hamiltonian flow of } \frac{|\xi|_m^2}{2} + V(m)$$

**Proposition.** *Let  $U$  be a coordinate patch, with coordinates  $y_1, \dots, y_n$ . Along each trajectory starting at  $(z, \zeta) \in T^*U$ , one can define intrinsically*

$$\Gamma^t : T_{z^t} M^{\mathbb{C}} \rightarrow T_{z^t} M^{\mathbb{C}}, \quad \text{where } T_{z^t} M^{\mathbb{C}} = T_{z^t} M \otimes \mathbb{C}$$

## Wave packets on Riemannian manifolds

Consider  $V \in C^\infty(M, \mathbb{R})$  and

$$H(h) := -h^2 \frac{\Delta_g}{2} + V$$

$$(z^t, \zeta^t) = \Phi^t(z, \zeta), \quad \text{Hamiltonian flow of } \frac{|\xi|_m^2}{2} + V(m)$$

**Proposition.** *Let  $U$  be a coordinate patch, with coordinates  $y_1, \dots, y_n$ . Along each trajectory starting at  $(z, \zeta) \in T^*U$ , one can define intrinsically*

$$\Gamma^t : T_{z^t} M^{\mathbb{C}} \rightarrow T_{z^t} M^{\mathbb{C}}, \quad \text{where } T_{z^t} M^{\mathbb{C}} = T_{z^t} M \otimes \mathbb{C}$$

(i.e.  $\Gamma^t$  is a complex tensor along the curve  $t \mapsto z^t$ )

## Wave packets on Riemannian manifolds

Consider  $V \in C^\infty(M, \mathbb{R})$  and

$$H(h) := -h^2 \frac{\Delta_g}{2} + V$$

$$(z^t, \zeta^t) = \Phi^t(z, \zeta), \quad \text{Hamiltonian flow of } \frac{|\xi|_m^2}{2} + V(m)$$

**Proposition.** *Let  $U$  be a coordinate patch, with coordinates  $y_1, \dots, y_n$ . Along each trajectory starting at  $(z, \zeta) \in T^*U$ , one can define intrinsically*

$$\Gamma^t : T_{z^t} M^{\mathbb{C}} \rightarrow T_{z^t} M^{\mathbb{C}}, \quad \text{where } T_{z^t} M^{\mathbb{C}} = T_{z^t} M \otimes \mathbb{C}$$

(i.e.  $\Gamma^t$  is a complex tensor along the curve  $t \mapsto z^t$ ) which is



## Wave packets on Riemannian manifolds

Consider  $V \in C^\infty(M, \mathbb{R})$  and

$$H(h) := -h^2 \frac{\Delta_g}{2} + V$$

$$(z^t, \zeta^t) = \Phi^t(z, \zeta), \quad \text{Hamiltonian flow of } \frac{|\xi|_m^2}{2} + V(m)$$

**Proposition.** *Let  $U$  be a coordinate patch, with coordinates  $y_1, \dots, y_n$ . Along each trajectory starting at  $(z, \zeta) \in T^*U$ , one can define intrinsically*

$$\Gamma^t : T_{z^t} M^{\mathbb{C}} \rightarrow T_{z^t} M^{\mathbb{C}}, \quad \text{where } T_{z^t} M^{\mathbb{C}} = T_{z^t} M \otimes \mathbb{C}$$

(i.e.  $\Gamma^t$  is a complex tensor along the curve  $t \mapsto z^t$ ) which is **symmetric**

$$\langle \Gamma^t X, Y \rangle_{z^t} = \langle X, \Gamma^t Y \rangle_{z^t}, \quad X, Y \in T_{z^t} M$$

## Wave packets on Riemannian manifolds

Consider  $V \in C^\infty(M, \mathbb{R})$  and

$$H(h) := -h^2 \frac{\Delta_g}{2} + V$$

$$(z^t, \zeta^t) = \Phi^t(z, \zeta), \quad \text{Hamiltonian flow of } \frac{|\xi|_m^2}{2} + V(m)$$

**Proposition.** *Let  $U$  be a coordinate patch, with coordinates  $y_1, \dots, y_n$ . Along each trajectory starting at  $(z, \zeta) \in T^*U$ , one can define intrinsically*

$$\Gamma^t : T_{z^t} M^{\mathbb{C}} \rightarrow T_{z^t} M^{\mathbb{C}}, \quad \text{where } T_{z^t} M^{\mathbb{C}} = T_{z^t} M \otimes \mathbb{C}$$

(i.e.  $\Gamma^t$  is a complex tensor along the curve  $t \mapsto z^t$ ) which is **symmetric**

$$\langle \Gamma^t X, Y \rangle_{z^t} = \langle X, \Gamma^t Y \rangle_{z^t}, \quad X, Y \in T_{z^t} M$$

has **positive definite imaginary part**

$$\text{Im} \langle \Gamma^t X, X \rangle_{z^t} > 0, \quad X \neq 0, X \in T_{z^t} M$$

## Wave packets on Riemannian manifolds

Consider  $V \in C^\infty(M, \mathbb{R})$  and

$$H(h) := -h^2 \frac{\Delta_g}{2} + V$$

$$(z^t, \zeta^t) = \Phi^t(z, \zeta), \quad \text{Hamiltonian flow of } \frac{|\xi|_m^2}{2} + V(m)$$

**Proposition.** *Let  $U$  be a coordinate patch, with coordinates  $y_1, \dots, y_n$ . Along each trajectory starting at  $(z, \zeta) \in T^*U$ , one can define intrinsically*

$$\Gamma^t : T_{z^t} M^{\mathbb{C}} \rightarrow T_{z^t} M^{\mathbb{C}}, \quad \text{where } T_{z^t} M^{\mathbb{C}} = T_{z^t} M \otimes \mathbb{C}$$

(i.e.  $\Gamma^t$  is a complex tensor along the curve  $t \mapsto z^t$ ) which is **symmetric**

$$\langle \Gamma^t X, Y \rangle_{z^t} = \langle X, \Gamma^t Y \rangle_{z^t}, \quad X, Y \in T_{z^t} M$$

has **positive definite imaginary part**

$$\text{Im} \langle \Gamma^t X, X \rangle_{z^t} > 0, \quad X \neq 0, X \in T_{z^t} M$$

and satisfies the **Riccati equation**

$$\nabla_{z^t} \Gamma^t = -\text{Hess}(V)_{z^t} - R_{z^t}(\cdot, \dot{z}^t) \dot{z}^t - (\Gamma^t)^2$$

where  $R_{z^t}$  is the Riemann tensor at  $z^t$

## Wave packets on Riemannian manifolds

**Proof.**

To construct  $\Gamma^t$  on  $\mathbb{R}^n$ , we have used the natural identifications

$$T_{(z,\zeta)}(T^*\mathbb{R}^n) = \mathbb{R}^n \oplus \mathbb{R}^n, \quad T_{(z^t,\zeta^t)}(T^*\mathbb{R}^n) = \mathbb{R}^n \oplus \mathbb{R}^n$$

## Wave packets on Riemannian manifolds

### Proof.

To construct  $\Gamma^t$  on  $\mathbb{R}^n$ , we have used the natural identifications

$$T_{(z,\zeta)}(T^*\mathbb{R}^n) = \mathbb{R}^n \oplus \mathbb{R}^n, \quad T_{(z^t,\zeta^t)}(T^*\mathbb{R}^n) = \mathbb{R}^n \oplus \mathbb{R}^n$$

**How to proceed on a manifold ?**

# Wave packets on Riemannian manifolds

## Proof.

To construct  $\Gamma^t$  on  $\mathbb{R}^n$ , we have used the natural identifications

$$T_{(z,\zeta)}(T^*\mathbb{R}^n) = \mathbb{R}^n \oplus \mathbb{R}^n, \quad T_{(z^t,\zeta^t)}(T^*\mathbb{R}^n) = \mathbb{R}^n \oplus \mathbb{R}^n$$

## How to proceed on a manifold ?

1. At starting points  $(z, \zeta)$  with  $z \in U$ , we split

$$T_{(z,\zeta)}(T^*M) \approx \mathbb{R}_y^n \oplus \mathbb{R}_\eta^n$$

using the (symplectic) coordinates  $(y_1, \dots, y_n, \eta_1, \dots, \eta_n)$  on  $T^*U$

# Wave packets on Riemannian manifolds

## Proof.

To construct  $\Gamma^t$  on  $\mathbb{R}^n$ , we have used the natural identifications

$$T_{(z,\zeta)}(T^*\mathbb{R}^n) = \mathbb{R}^n \oplus \mathbb{R}^n, \quad T_{(z^t,\zeta^t)}(T^*\mathbb{R}^n) = \mathbb{R}^n \oplus \mathbb{R}^n$$

## How to proceed on a manifold ?

1. At starting points  $(z, \zeta)$  with  $z \in U$ , we split

$$T_{(z,\zeta)}(T^*M) \approx \mathbb{R}_y^n \oplus \mathbb{R}_\eta^n$$

using the (symplectic) coordinates  $(y_1, \dots, y_n, \eta_1, \dots, \eta_n)$  on  $T^*U$

2. At points  $(z^t, \zeta^t)$ , we use the (global) identification  $\mathcal{I}_g : T^*M \rightarrow TM$

# Wave packets on Riemannian manifolds

## Proof.

To construct  $\Gamma^t$  on  $\mathbb{R}^n$ , we have used the natural identifications

$$T_{(z,\zeta)}(T^*\mathbb{R}^n) = \mathbb{R}^n \oplus \mathbb{R}^n, \quad T_{(z^t,\zeta^t)}(T^*\mathbb{R}^n) = \mathbb{R}^n \oplus \mathbb{R}^n$$

## How to proceed on a manifold ?

1. At starting points  $(z, \zeta)$  with  $z \in U$ , we split

$$T_{(z,\zeta)}(T^*M) \approx \mathbb{R}_y^n \oplus \mathbb{R}_\eta^n$$

using the (symplectic) coordinates  $(y_1, \dots, y_n, \eta_1, \dots, \eta_n)$  on  $T^*U$

2. At points  $(z^t, \zeta^t)$ , we use the (global) identification  $\mathcal{I}_g : T^*M \rightarrow TM$

$$\mathcal{I}_g(z^t, \zeta^t) = (z^t, \dot{z}^t)$$



# Wave packets on Riemannian manifolds

## Proof.

To construct  $\Gamma^t$  on  $\mathbb{R}^n$ , we have used the natural identifications

$$T_{(z,\zeta)}(T^*\mathbb{R}^n) = \mathbb{R}^n \oplus \mathbb{R}^n, \quad T_{(z^t,\zeta^t)}(T^*\mathbb{R}^n) = \mathbb{R}^n \oplus \mathbb{R}^n$$

## How to proceed on a manifold ?

1. At starting points  $(z, \zeta)$  with  $z \in U$ , we split

$$T_{(z,\zeta)}(T^*M) \approx \mathbb{R}_y^n \oplus \mathbb{R}_\eta^n$$

using the (symplectic) coordinates  $(y_1, \dots, y_n, \eta_1, \dots, \eta_n)$  on  $T^*U$

2. At points  $(z^t, \zeta^t)$ , we use the (global) identification  $\mathcal{I}_g : T^*M \rightarrow TM$

$$\mathcal{I}_g(z^t, \zeta^t) = (z^t, \dot{z}^t)$$

and split along **horizontal** and **vertical** spaces

$$T_{(z^t, \dot{z}^t)}(\mathcal{I}_g T^*M) = \mathcal{H}_{(z^t, \dot{z}^t)} \oplus \mathcal{V}_{(z^t, \dot{z}^t)}$$

# Wave packets on Riemannian manifolds

## Proof.

To construct  $\Gamma^t$  on  $\mathbb{R}^n$ , we have used the natural identifications

$$T_{(z,\zeta)}(T^*\mathbb{R}^n) = \mathbb{R}^n \oplus \mathbb{R}^n, \quad T_{(z^t,\zeta^t)}(T^*\mathbb{R}^n) = \mathbb{R}^n \oplus \mathbb{R}^n$$

## How to proceed on a manifold ?

1. At starting points  $(z, \zeta)$  with  $z \in U$ , we split

$$T_{(z,\zeta)}(T^*M) \approx \mathbb{R}_y^n \oplus \mathbb{R}_\eta^n$$

using the (symplectic) coordinates  $(y_1, \dots, y_n, \eta_1, \dots, \eta_n)$  on  $T^*U$

2. At points  $(z^t, \zeta^t)$ , we use the (global) identification  $\mathcal{I}_g : T^*M \rightarrow TM$

$$\mathcal{I}_g(z^t, \zeta^t) = (z^t, \dot{z}^t)$$

and split along **horizontal** and **vertical** spaces

$$T_{(z^t, \dot{z}^t)}(\mathcal{I}_g T^*M) = \mathcal{H}_{(z^t, \dot{z}^t)} \oplus \mathcal{V}_{(z^t, \dot{z}^t)}$$

This gives a natural block decomposition

$$d(\mathcal{I}_g \circ \Phi^t) = \begin{pmatrix} \mathcal{L}_A & \mathcal{L}_B \\ \mathcal{L}_C & \mathcal{L}_D \end{pmatrix} : \mathbb{R}_y^n \oplus \mathbb{R}_\eta^n \rightarrow \mathcal{H}_{(z^t, \dot{z}^t)} \oplus \mathcal{V}_{(z^t, \dot{z}^t)}$$

## Wave packets on Riemannian manifolds

**Proof (continued).** One can then define

$$(\mathcal{L}_C + i\mathcal{L}_D)(\mathcal{L}_A + i\mathcal{L}_B)^{-1} : \mathcal{H}_{(z^t, \dot{z}^t)}^{\mathbb{C}} \rightarrow \mathcal{V}_{(z^t, \dot{z}^t)}^{\mathbb{C}}$$

## Wave packets on Riemannian manifolds

**Proof (continued).** One can then define

$$(\mathcal{L}_C + i\mathcal{L}_D)(\mathcal{L}_A + i\mathcal{L}_B)^{-1} : \mathcal{H}_{(z^t, \dot{z}^t)}^{\mathbb{C}} \rightarrow \mathcal{V}_{(z^t, \dot{z}^t)}^{\mathbb{C}}$$

and then define  $\Gamma^t$  by composition with the natural isomorphisms

$$T_{z^t} M^{\mathbb{C}} \rightarrow \mathcal{H}_{(z^t, \dot{z}^t)}^{\mathbb{C}}, \quad \mathcal{V}_{(z^t, \dot{z}^t)}^{\mathbb{C}} \rightarrow T_{z^t} M^{\mathbb{C}}$$

## Wave packets on Riemannian manifolds

**Proof (continued).** One can then define

$$(\mathcal{L}_C + i\mathcal{L}_D)(\mathcal{L}_A + i\mathcal{L}_B)^{-1} : \mathcal{H}_{(z^t, \dot{z}^t)}^{\mathbb{C}} \rightarrow \mathcal{V}_{(z^t, \dot{z}^t)}^{\mathbb{C}}$$

and then define  $\Gamma^t$  by composition with the natural isomorphisms

$$T_{z^t}M^{\mathbb{C}} \rightarrow \mathcal{H}_{(z^t, \dot{z}^t)}^{\mathbb{C}}, \quad \mathcal{V}_{(z^t, \dot{z}^t)}^{\mathbb{C}} \rightarrow T_{z^t}M^{\mathbb{C}}$$

More concretely, using local coordinates  $(x_1, \dots, x_n)$  near  $z^t$ , the matrix of  $\Gamma^t$  reads

## Wave packets on Riemannian manifolds

**Proof (continued).** One can then define

$$(\mathcal{L}_C + i\mathcal{L}_D)(\mathcal{L}_A + i\mathcal{L}_B)^{-1} : \mathcal{H}_{(z^t, \dot{z}^t)}^{\mathbb{C}} \rightarrow \mathcal{V}_{(z^t, \dot{z}^t)}^{\mathbb{C}}$$

and then define  $\Gamma^t$  by composition with the natural isomorphisms

$$T_{z^t}M^{\mathbb{C}} \rightarrow \mathcal{H}_{(z^t, \dot{z}^t)}^{\mathbb{C}}, \quad \mathcal{V}_{(z^t, \dot{z}^t)}^{\mathbb{C}} \rightarrow T_{z^t}M^{\mathbb{C}}$$

More concretely, using local coordinates  $(x_1, \dots, x_n)$  near  $z^t$ , the matrix of  $\Gamma^t$  reads

$$G^{-1}(C^t + iD^t)(A^t + iB^t)^{-1} - G^{-1}\Sigma^t$$

## Wave packets on Riemannian manifolds

**Proof (continued).** One can then define

$$(\mathcal{L}_C + i\mathcal{L}_D)(\mathcal{L}_A + i\mathcal{L}_B)^{-1} : \mathcal{H}_{(z^t, \dot{z}^t)}^{\mathbb{C}} \rightarrow \mathcal{V}_{(z^t, \dot{z}^t)}^{\mathbb{C}}$$

and then define  $\Gamma^t$  by composition with the natural isomorphisms

$$T_{z^t}M^{\mathbb{C}} \rightarrow \mathcal{H}_{(z^t, \dot{z}^t)}^{\mathbb{C}}, \quad \mathcal{V}_{(z^t, \dot{z}^t)}^{\mathbb{C}} \rightarrow T_{z^t}M^{\mathbb{C}}$$

More concretely, using local coordinates  $(x_1, \dots, x_n)$  near  $z^t$ , the matrix of  $\Gamma^t$  reads

$$G^{-1}(C^t + iD^t)(A^t + iB^t)^{-1} - G^{-1}\Sigma^t$$

with

$$G^{-1} = (g^{ij}(x^t)),$$

## Wave packets on Riemannian manifolds

**Proof (continued).** One can then define

$$(\mathcal{L}_C + i\mathcal{L}_D)(\mathcal{L}_A + i\mathcal{L}_B)^{-1} : \mathcal{H}_{(z^t, \dot{z}^t)}^{\mathbb{C}} \rightarrow \mathcal{V}_{(z^t, \dot{z}^t)}^{\mathbb{C}}$$

and then define  $\Gamma^t$  by composition with the natural isomorphisms

$$T_{z^t}M^{\mathbb{C}} \rightarrow \mathcal{H}_{(z^t, \dot{z}^t)}^{\mathbb{C}}, \quad \mathcal{V}_{(z^t, \dot{z}^t)}^{\mathbb{C}} \rightarrow T_{z^t}M^{\mathbb{C}}$$

More concretely, using local coordinates  $(x_1, \dots, x_n)$  near  $z^t$ , the matrix of  $\Gamma^t$  reads

$$G^{-1}(C^t + iD^t)(A^t + iB^t)^{-1} - G^{-1}\Sigma^t$$

with

$$G^{-1} = (g^{ij}(x^t)), \quad \Sigma_{ij}^t = \sum_{k,l} g_{kl}(x^t) \Gamma_{ij}^l(x^t) \dot{x}_k^t, \quad x^t = x(z^t)$$



## Wave packets on Riemannian manifolds

**Proof (continued).** One can then define

$$(\mathcal{L}_C + i\mathcal{L}_D)(\mathcal{L}_A + i\mathcal{L}_B)^{-1} : \mathcal{H}_{(z^t, \dot{z}^t)}^{\mathbb{C}} \rightarrow \mathcal{V}_{(z^t, \dot{z}^t)}^{\mathbb{C}}$$

and then define  $\Gamma^t$  by composition with the natural isomorphisms

$$T_{z^t}M^{\mathbb{C}} \rightarrow \mathcal{H}_{(z^t, \dot{z}^t)}^{\mathbb{C}}, \quad \mathcal{V}_{(z^t, \dot{z}^t)}^{\mathbb{C}} \rightarrow T_{z^t}M^{\mathbb{C}}$$

More concretely, using local coordinates  $(x_1, \dots, x_n)$  near  $z^t$ , the matrix of  $\Gamma^t$  reads

$$G^{-1}(C^t + iD^t)(A^t + iB^t)^{-1} - G^{-1}\Sigma^t$$

with

$$G^{-1} = (g^{ij}(x^t)), \quad \Sigma_{ij}^t = \sum_{k,l} g_{kl}(x^t) \Gamma_{ij}^l(x^t) \dot{x}_k^t, \quad x^t = x(z^t)$$

and

$$\begin{pmatrix} A^t & B^t \\ C^t & D^t \end{pmatrix} = \begin{pmatrix} \partial x^t / \partial y & \partial x^t / \partial \eta \\ \partial \xi^t / \partial y & \partial \xi^t / \partial \eta \end{pmatrix}$$

## Wave packets on Riemannian manifolds

**Proof (continued).** One can then define

$$(\mathcal{L}_C + i\mathcal{L}_D)(\mathcal{L}_A + i\mathcal{L}_B)^{-1} : \mathcal{H}_{(z^t, \dot{z}^t)}^{\mathbb{C}} \rightarrow \mathcal{V}_{(z^t, \dot{z}^t)}^{\mathbb{C}}$$

and then define  $\Gamma^t$  by composition with the natural isomorphisms

$$T_{z^t}M^{\mathbb{C}} \rightarrow \mathcal{H}_{(z^t, \dot{z}^t)}^{\mathbb{C}}, \quad \mathcal{V}_{(z^t, \dot{z}^t)}^{\mathbb{C}} \rightarrow T_{z^t}M^{\mathbb{C}}$$

More concretely, using local coordinates  $(x_1, \dots, x_n)$  near  $z^t$ , the matrix of  $\Gamma^t$  reads

$$G^{-1}(C^t + iD^t)(A^t + iB^t)^{-1} - G^{-1}\Sigma^t$$

with

$$G^{-1} = (g^{ij}(x^t)), \quad \Sigma_{ij}^t = \sum_{k,l} g_{kl}(x^t) \Gamma_{ij}^l(x^t) \dot{x}_k^t, \quad x^t = x(z^t)$$

and

$$\begin{pmatrix} A^t & B^t \\ C^t & D^t \end{pmatrix} = \begin{pmatrix} \partial x^t / \partial y & \partial x^t / \partial \eta \\ \partial \xi^t / \partial y & \partial \xi^t / \partial \eta \end{pmatrix}$$

$\implies$  Symmetry of  $\Gamma^t$ ,

## Wave packets on Riemannian manifolds

**Proof (continued).** One can then define

$$(\mathcal{L}_C + i\mathcal{L}_D)(\mathcal{L}_A + i\mathcal{L}_B)^{-1} : \mathcal{H}_{(z^t, \dot{z}^t)}^{\mathbb{C}} \rightarrow \mathcal{V}_{(z^t, \dot{z}^t)}^{\mathbb{C}}$$

and then define  $\Gamma^t$  by composition with the natural isomorphisms

$$T_{z^t}M^{\mathbb{C}} \rightarrow \mathcal{H}_{(z^t, \dot{z}^t)}^{\mathbb{C}}, \quad \mathcal{V}_{(z^t, \dot{z}^t)}^{\mathbb{C}} \rightarrow T_{z^t}M^{\mathbb{C}}$$

More concretely, using local coordinates  $(x_1, \dots, x_n)$  near  $z^t$ , the matrix of  $\Gamma^t$  reads

$$G^{-1}(C^t + iD^t)(A^t + iB^t)^{-1} - G^{-1}\Sigma^t$$

with

$$G^{-1} = (g^{ij}(x^t)), \quad \Sigma_{ij}^t = \sum_{k,l} g_{kl}(x^t) \Gamma_{ij}^l(x^t) \dot{x}_k^t, \quad x^t = x(z^t)$$

and

$$\begin{pmatrix} A^t & B^t \\ C^t & D^t \end{pmatrix} = \begin{pmatrix} \partial x^t / \partial y & \partial x^t / \partial \eta \\ \partial \xi^t / \partial y & \partial \xi^t / \partial \eta \end{pmatrix}$$

$\implies$  Symmetry of  $\Gamma^t$ , positivity of  $\text{Im}(\Gamma^t)$

## Wave packets on Riemannian manifolds

**Proof (continued).** One can then define

$$(\mathcal{L}_C + i\mathcal{L}_D)(\mathcal{L}_A + i\mathcal{L}_B)^{-1} : \mathcal{H}_{(z^t, \dot{z}^t)}^{\mathbb{C}} \rightarrow \mathcal{V}_{(z^t, \dot{z}^t)}^{\mathbb{C}}$$

and then define  $\Gamma^t$  by composition with the natural isomorphisms

$$T_{z^t}M^{\mathbb{C}} \rightarrow \mathcal{H}_{(z^t, \dot{z}^t)}^{\mathbb{C}}, \quad \mathcal{V}_{(z^t, \dot{z}^t)}^{\mathbb{C}} \rightarrow T_{z^t}M^{\mathbb{C}}$$

More concretely, using local coordinates  $(x_1, \dots, x_n)$  near  $z^t$ , the matrix of  $\Gamma^t$  reads

$$G^{-1}(C^t + iD^t)(A^t + iB^t)^{-1} - G^{-1}\Sigma^t$$

with

$$G^{-1} = (g^{ij}(x^t)), \quad \Sigma_{ij}^t = \sum_{k,l} g_{kl}(x^t) \Gamma_{ij}^l(x^t) \dot{x}_k^t, \quad x^t = x(z^t)$$

and

$$\begin{pmatrix} A^t & B^t \\ C^t & D^t \end{pmatrix} = \begin{pmatrix} \partial x^t / \partial y & \partial x^t / \partial \eta \\ \partial \xi^t / \partial y & \partial \xi^t / \partial \eta \end{pmatrix}$$

$\implies$  Symmetry of  $\Gamma^t$ , positivity of  $\text{Im}(\Gamma^t)$  + Riccati equation by direct computation #

## Wave packets on Riemannian manifolds

**Proof (continued).** One can then define

$$(\mathcal{L}_C + i\mathcal{L}_D)(\mathcal{L}_A + i\mathcal{L}_B)^{-1} : \mathcal{H}_{(z^t, \dot{z}^t)}^{\mathbb{C}} \rightarrow \mathcal{V}_{(z^t, \dot{z}^t)}^{\mathbb{C}}$$

and then define  $\Gamma^t$  by composition with the natural isomorphisms

$$T_{z^t} M^{\mathbb{C}} \rightarrow \mathcal{H}_{(z^t, \dot{z}^t)}^{\mathbb{C}}, \quad \mathcal{V}_{(z^t, \dot{z}^t)}^{\mathbb{C}} \rightarrow T_{z^t} M^{\mathbb{C}}$$

More concretely, using local coordinates  $(x_1, \dots, x_n)$  near  $z^t$ , the matrix of  $\Gamma^t$  reads

$$G^{-1}(C^t + iD^t)(A^t + iB^t)^{-1} - G^{-1}\Sigma^t$$

with

$$G^{-1} = (g^{ij}(x^t)), \quad \Sigma_{ij}^t = \sum_{k,l} g_{kl}(x^t) \Gamma_{ij}^l(x^t) \dot{x}_k^t, \quad x^t = x(z^t)$$

and

$$\begin{pmatrix} A^t & B^t \\ C^t & D^t \end{pmatrix} = \begin{pmatrix} \partial x^t / \partial y & \partial x^t / \partial \eta \\ \partial \xi^t / \partial y & \partial \xi^t / \partial \eta \end{pmatrix}$$

$\implies$  Symmetry of  $\Gamma^t$ , positivity of  $\text{Im}(\Gamma^t)$  + Riccati equation by direct computation #

**Rem.** If  $(\tilde{y}_1, \dots, \tilde{y}_n)$  are other coordinates on  $U$ , the matrix of  $\Gamma^t$  is changed into

$$G^{-1}(\tilde{C}^t + \tilde{D}^t Z)(\tilde{A}^t + \tilde{B}^t Z)^{-1} - G^{-1}\Sigma^t,$$

## Wave packets on Riemannian manifolds

**Proof (continued).** One can then define

$$(\mathcal{L}_C + i\mathcal{L}_D)(\mathcal{L}_A + i\mathcal{L}_B)^{-1} : \mathcal{H}_{(z^t, \dot{z}^t)}^{\mathbb{C}} \rightarrow \mathcal{V}_{(z^t, \dot{z}^t)}^{\mathbb{C}}$$

and then define  $\Gamma^t$  by composition with the natural isomorphisms

$$T_{z^t} M^{\mathbb{C}} \rightarrow \mathcal{H}_{(z^t, \dot{z}^t)}^{\mathbb{C}}, \quad \mathcal{V}_{(z^t, \dot{z}^t)}^{\mathbb{C}} \rightarrow T_{z^t} M^{\mathbb{C}}$$

More concretely, using local coordinates  $(x_1, \dots, x_n)$  near  $z^t$ , the matrix of  $\Gamma^t$  reads

$$G^{-1}(C^t + iD^t)(A^t + iB^t)^{-1} - G^{-1}\Sigma^t$$

with

$$G^{-1} = (g^{ij}(x^t)), \quad \Sigma_{ij}^t = \sum_{k,l} g_{kl}(x^t) \Gamma_{ij}^l(x^t) \dot{x}_k^t, \quad x^t = x(z^t)$$

and

$$\begin{pmatrix} A^t & B^t \\ C^t & D^t \end{pmatrix} = \begin{pmatrix} \partial x^t / \partial y & \partial x^t / \partial \eta \\ \partial \xi^t / \partial y & \partial \xi^t / \partial \eta \end{pmatrix}$$

$\implies$  Symmetry of  $\Gamma^t$ , positivity of  $\text{Im}(\Gamma^t)$  + Riccati equation by direct computation #

**Rem.** If  $(\tilde{y}_1, \dots, \tilde{y}_n)$  are other coordinates on  $U$ , the matrix of  $\Gamma^t$  is changed into

$$G^{-1}(\tilde{C}^t + \tilde{D}^t Z)(\tilde{A}^t + \tilde{B}^t Z)^{-1} - G^{-1}\Sigma^t, \quad Z = \left( \frac{\partial \tilde{\eta}}{\partial y} + i \frac{\partial \tilde{\eta}}{\partial \eta} \right) \left( \frac{\partial \tilde{y}}{\partial y} + i \frac{\partial \tilde{y}}{\partial \eta} \right)^{-1}$$

# Wave packets on Riemannian manifolds

## Definition of gaussian wave packets

## Wave packets on Riemannian manifolds

**Definition of gaussian wave packets** Let  $\rho \in C_0^\infty(-r_0, r_0)$ , equal to 1 near 0.



## Wave packets on Riemannian manifolds

**Definition of gaussian wave packets** Let  $\rho \in C_0^\infty(-r_0, r_0)$ , equal to 1 near 0.

$$\Psi_{z,\zeta}^h(m) := (\pi h)^{-\frac{n}{4}} \gamma^0 \exp \frac{i}{h} \left( \zeta \cdot W_z^m + \frac{1}{2} \langle \Gamma^0 W_z^m, W_z^m \rangle_z \right) \rho(d_g(z, m)),$$

for  $m \in M$  and  $(z, \zeta) \in T^*U$  (i.e.  $\zeta \in T_z^*U$ )

$$\gamma^0 = \det(g_{jk}(y(z)))^{-\frac{1}{4}}$$

## Wave packets on Riemannian manifolds

**Definition of gaussian wave packets** Let  $\rho \in C_0^\infty(-r_0, r_0)$ , equal to 1 near 0.

$$\Psi_{z,\zeta}^h(m) := (\pi h)^{-\frac{n}{4}} \gamma^0 \exp \frac{i}{h} \left( \zeta \cdot W_z^m + \frac{1}{2} \langle \Gamma^0 W_z^m, W_z^m \rangle_z \right) \rho(d_g(z, m)),$$

for  $m \in M$  and  $(z, \zeta) \in T^*U$  (i.e.  $\zeta \in T_z^*U$ )

$$\gamma^0 = \det(g_{jk}(y(z)))^{-\frac{1}{4}}$$

**Rem.**  $\Psi_{z,\zeta}^h(m) = 0$  if  $d_g(z, m) \geq r_0$ .

## Wave packets on Riemannian manifolds

**Definition of gaussian wave packets** Let  $\rho \in C_0^\infty(-r_0, r_0)$ , equal to 1 near 0.

$$\Psi_{z,\zeta}^h(m) := (\pi h)^{-\frac{n}{4}} \gamma^0 \exp \frac{i}{h} \left( \zeta \cdot W_z^m + \frac{1}{2} \langle \Gamma^0 W_z^m, W_z^m \rangle_z \right) \rho(d_g(z, m)),$$

for  $m \in M$  and  $(z, \zeta) \in T^*U$  (i.e.  $\zeta \in T_z^*U$ )

$$\gamma^0 = \det(g_{jk}(y(z)))^{-\frac{1}{4}}$$

**Rem.**  $\Psi_{z,\zeta}^h(m) = 0$  if  $d_g(z, m) \geq r_0$ .

**Proposition [Wave packet decomposition - Approximate Bargmann transform]**

## Wave packets on Riemannian manifolds

**Definition of gaussian wave packets** Let  $\rho \in C_0^\infty(-r_0, r_0)$ , equal to 1 near 0.

$$\Psi_{z,\zeta}^h(m) := (\pi h)^{-\frac{n}{4}} \gamma^0 \exp \frac{i}{h} \left( \zeta \cdot W_z^m + \frac{1}{2} \langle \Gamma^0 W_z^m, W_z^m \rangle_z \right) \rho(d_g(z, m)),$$

for  $m \in M$  and  $(z, \zeta) \in T^*U$  (i.e.  $\zeta \in T_z^*U$ )

$$\gamma^0 = \det(g_{jk}(y(z)))^{-\frac{1}{4}}$$

**Rem.**  $\Psi_{z,\zeta}^h(m) = 0$  if  $d_g(z, m) \geq r_0$ .

**Proposition [Wave packet decomposition - Approximate Bargmann transform]** Set

$$B_h u(z, \zeta) := \left\langle \Psi_{z,\zeta}^h, u \right\rangle_{L^2(M)}, \quad u \in C_0^\infty(U)$$

## Wave packets on Riemannian manifolds

**Definition of gaussian wave packets** Let  $\rho \in C_0^\infty(-r_0, r_0)$ , equal to 1 near 0.

$$\Psi_{z,\zeta}^h(m) := (\pi h)^{-\frac{n}{4}} \gamma^0 \exp \frac{i}{h} \left( \zeta \cdot W_z^m + \frac{1}{2} \langle \Gamma^0 W_z^m, W_z^m \rangle_z \right) \rho(d_g(z, m)),$$

for  $m \in M$  and  $(z, \zeta) \in T^*U$  (i.e.  $\zeta \in T_z^*U$ )

$$\gamma^0 = \det(g_{jk}(y(z)))^{-\frac{1}{4}}$$

**Rem.**  $\Psi_{z,\zeta}^h(m) = 0$  if  $d_g(z, m) \geq r_0$ .

**Proposition [Wave packet decomposition - Approximate Bargmann transform]** Set

$$B_h u(z, \zeta) := \left\langle \Psi_{z,\zeta}^h, u \right\rangle_{L^2(M)}, \quad u \in C_0^\infty(U)$$

Then

$$(2\pi h)^{-n} B_h^* B_h u = a(h)u$$

## Wave packets on Riemannian manifolds

**Definition of gaussian wave packets** Let  $\rho \in C_0^\infty(-r_0, r_0)$ , equal to 1 near 0.

$$\Psi_{z,\zeta}^h(m) := (\pi h)^{-\frac{n}{4}} \gamma^0 \exp \frac{i}{h} \left( \zeta \cdot W_z^m + \frac{1}{2} \langle \Gamma^0 W_z^m, W_z^m \rangle_z \right) \rho(d_g(z, m)),$$

for  $m \in M$  and  $(z, \zeta) \in T^*U$  (i.e.  $\zeta \in T_z^*U$ )

$$\gamma^0 = \det(g_{jk}(y(z)))^{-\frac{1}{4}}$$

**Rem.**  $\Psi_{z,\zeta}^h(m) = 0$  if  $d_g(z, m) \geq r_0$ .

**Proposition [Wave packet decomposition - Approximate Bargmann transform]** Set

$$B_h u(z, \zeta) := \left\langle \Psi_{z,\zeta}^h, u \right\rangle_{L^2(M)}, \quad u \in C_0^\infty(U)$$

Then

$$(2\pi h)^{-n} B_h^* B_h u = a(h)u = \left( 1 + h^{\frac{1}{2}} a_1 + h^1 a_2 + \dots \right) u$$

with  $a(h), a_1, a_2, \dots \in C^\infty$

## Wave packets on Riemannian manifolds

**Definition of gaussian wave packets** Let  $\rho \in C_0^\infty(-r_0, r_0)$ , equal to 1 near 0.

$$\Psi_{z,\zeta}^h(m) := (\pi h)^{-\frac{n}{4}} \gamma^0 \exp \frac{i}{h} \left( \zeta \cdot W_z^m + \frac{1}{2} \langle \Gamma^0 W_z^m, W_z^m \rangle_z \right) \rho(d_g(z, m)),$$

for  $m \in M$  and  $(z, \zeta) \in T^*U$  (i.e.  $\zeta \in T_z^*U$ )

$$\gamma^0 = \det(g_{jk}(y(z)))^{-\frac{1}{4}}$$

**Rem.**  $\Psi_{z,\zeta}^h(m) = 0$  if  $d_g(z, m) \geq r_0$ .

**Proposition [Wave packet decomposition - Approximate Bargmann transform]** Set

$$B_h u(z, \zeta) := \left\langle \Psi_{z,\zeta}^h, u \right\rangle_{L^2(M)}, \quad u \in C_0^\infty(U)$$

Then

$$(2\pi h)^{-n} B_h^* B_h u = a(h)u = \left( 1 + h^{\frac{1}{2}} a_1 + h^1 a_2 + \dots \right) u$$

with  $a(h), a_1, a_2, \dots \in C^\infty$ , i.e.

$$(2\pi h)^{-n} \int \int_{T^*U} B_h u(z, \zeta) \Psi_{z,\zeta}^h dz d\zeta = a(h)u$$

# Wave packets on Riemannian manifolds

**Theorem [Propagation of gaussian wave packets]**



## Wave packets on Riemannian manifolds

**Theorem [Propagation of gaussian wave packets]** *In the limit  $\hbar \rightarrow 0$ , and under general conditions on  $V$  (e.g. all covariant derivatives bounded),*

## Wave packets on Riemannian manifolds

**Theorem [Propagation of gaussian wave packets]** *In the limit  $\hbar \rightarrow 0$ , and under general conditions on  $V$  (e.g. all covariant derivatives bounded),*

$$e^{-i\frac{t}{\hbar}H(\hbar)}\psi_{z,\zeta}^{\hbar}(m)$$

## Wave packets on Riemannian manifolds

**Theorem [Propagation of gaussian wave packets]** *In the limit  $h \rightarrow 0$ , and under general conditions on  $V$  (e.g. all covariant derivatives bounded),*

$$e^{-i\frac{t}{h}H(h)}\psi_{z,\zeta}^h(m)$$

*is well approximated by*

$$(\pi h)^{-\frac{n}{4}}\gamma^t\mathcal{A}_t^h(x)\exp\frac{i}{h}\left(S^t + \zeta^t \cdot W_{z^t}^m + \frac{1}{2}\langle \Gamma^t W_{z^t}^m, W_{z^t}^m \rangle_{z^t}\right)\rho(d_g(z_t, m))$$

## Wave packets on Riemannian manifolds

**Theorem [Propagation of gaussian wave packets]** *In the limit  $h \rightarrow 0$ , and under general conditions on  $V$  (e.g. all covariant derivatives bounded),*

$$e^{-i\frac{t}{h}H(h)}\psi_{z,\zeta}^h(m)$$

*is well approximated by*

$$(\pi h)^{-\frac{n}{4}}\gamma^t\mathcal{A}_t^h(x)\exp\frac{i}{h}\left(S^t + \zeta^t \cdot W_{z^t}^m + \frac{1}{2}\langle \Gamma^t W_{z^t}^m, W_{z^t}^m \rangle_{z^t}\right)\rho(d_g(z_t, m))$$

*with*

$$\gamma^t = \det(g_{jk}(x^t))^{-1/4}\det(A^t + iB^t)^{-1/2}$$

## Wave packets on Riemannian manifolds

**Theorem [Propagation of gaussian wave packets]** *In the limit  $h \rightarrow 0$ , and under general conditions on  $V$  (e.g. all covariant derivatives bounded),*

$$e^{-i\frac{t}{h}H(h)}\psi_{z,\zeta}^h(m)$$

*is well approximated by*

$$(\pi h)^{-\frac{n}{4}}\gamma^t\mathcal{A}_t^h(x)\exp\frac{i}{h}\left(S^t + \zeta^t \cdot W_{z^t}^m + \frac{1}{2}\langle \Gamma^t W_{z^t}^m, W_{z^t}^m \rangle_{z^t}\right)\rho(d_g(z_t, m))$$

*with*

$$\gamma^t = \det(g_{jk}(x^t))^{-1/4}\det(A^t + iB^t)^{-1/2}$$

*and an amplitude of the form*

$$\mathcal{A}_t^h(x) \sim 1 + \sum_{j \geq 1} h^{\frac{j}{2}} T_j \left( t, z^t, \zeta^t, \frac{W_{z^t}^m}{h^{\frac{1}{2}}} \right)$$

## Wave packets on Riemannian manifolds

**Theorem [Propagation of gaussian wave packets]** *In the limit  $h \rightarrow 0$ , and under general conditions on  $V$  (e.g. all covariant derivatives bounded),*

$$e^{-i\frac{t}{h}H(h)}\psi_{z,\zeta}^h(m)$$

*is well approximated by*

$$(\pi h)^{-\frac{n}{4}}\gamma^t\mathcal{A}_t^h(x)\exp\frac{i}{h}\left(S^t + \zeta^t \cdot W_{z^t}^m + \frac{1}{2}\langle \Gamma^t W_{z^t}^m, W_{z^t}^m \rangle_{z^t}\right)\rho(d_g(z_t, m))$$

*with*

$$\gamma^t = \det(g_{jk}(x^t))^{-1/4}\det(A^t + iB^t)^{-1/2}$$

*and an amplitude of the form*

$$\mathcal{A}_t^h(x) \sim 1 + \sum_{j \geq 1} h^{\frac{j}{2}} T_j\left(t, z^t, \zeta^t, \frac{W_{z^t}^m}{h^{\frac{1}{2}}}\right)$$

*for times  $|t| \leq C_0|\ln h|$*

## Wave packets on Riemannian manifolds

**Theorem [Propagation of gaussian wave packets]** *In the limit  $h \rightarrow 0$ , and under general conditions on  $V$  (e.g. all covariant derivatives bounded),*

$$e^{-i\frac{t}{h}H(h)}\psi_{z,\zeta}^h(m)$$

*is well approximated by*

$$(\pi h)^{-\frac{n}{4}}\gamma^t\mathcal{A}_t^h(x)\exp\frac{i}{h}\left(S^t + \zeta^t \cdot W_{z^t}^m + \frac{1}{2}\langle \Gamma^t W_{z^t}^m, W_{z^t}^m \rangle_{z^t}\right)\rho(d_g(z_t, m))$$

*with*

$$\gamma^t = \det(g_{jk}(x^t))^{-1/4}\det(A^t + iB^t)^{-1/2}$$

*and an amplitude of the form*

$$\mathcal{A}_t^h(x) \sim 1 + \sum_{j \geq 1} h^{\frac{j}{2}} T_j\left(t, z^t, \zeta^t, \frac{W_{z^t}^m}{h^{\frac{1}{2}}}\right)$$

*for times  $|t| \leq C_0|\ln h|$  with  $T_j(t, z^t, \zeta^t, \cdot)$  polynomial (i.e. sum of tensors) of degree at most  $3j$ ,*

## Wave packets on Riemannian manifolds

**Theorem [Propagation of gaussian wave packets]** *In the limit  $h \rightarrow 0$ , and under general conditions on  $V$  (e.g. all covariant derivatives bounded),*

$$e^{-i\frac{t}{h}H(h)}\psi_{z,\zeta}^h(m)$$

*is well approximated by*

$$(\pi h)^{-\frac{n}{4}}\gamma^t\mathcal{A}_t^h(x)\exp\frac{i}{h}\left(S^t + \zeta^t \cdot W_{z^t}^m + \frac{1}{2}\langle \Gamma^t W_{z^t}^m, W_{z^t}^m \rangle_{z^t}\right)\rho(d_g(z_t, m))$$

*with*

$$\gamma^t = \det(g_{jk}(x^t))^{-1/4}\det(A^t + iB^t)^{-1/2}$$

*and an amplitude of the form*

$$\mathcal{A}_t^h(x) \sim 1 + \sum_{j \geq 1} h^{\frac{j}{2}} T_j\left(t, z^t, \zeta^t, \frac{W_{z^t}^m}{h^{\frac{1}{2}}}\right)$$

*for times  $|t| \leq C_0|\ln h|$  with  $T_j(t, z^t, \zeta^t, \cdot)$  polynomial (i.e. sum of tensors) of degree at most  $3j$ , depending on the classical trajectory and the Taylor expansions of  $V$  and  $W^m$  at  $z^t$ .*



## Wave packets on Riemannian manifolds

**Theorem [Propagation of gaussian wave packets]** *In the limit  $h \rightarrow 0$ , and under general conditions on  $V$  (e.g. all covariant derivatives bounded),*

$$e^{-i\frac{t}{h}H(h)}\psi_{z,\zeta}^h(m)$$

*is well approximated by*

$$(\pi h)^{-\frac{n}{4}}\gamma^t\mathcal{A}_t^h(x)\exp\frac{i}{h}\left(S^t + \zeta^t \cdot W_{z^t}^m + \frac{1}{2}\langle \Gamma^t W_{z^t}^m, W_{z^t}^m \rangle_{z^t}\right)\rho(d_g(z_t, m))$$

*with*

$$\gamma^t = \det(g_{jk}(x^t))^{-1/4}\det(A^t + iB^t)^{-1/2}$$

*and an amplitude of the form*

$$\mathcal{A}_t^h(x) \sim 1 + \sum_{j \geq 1} h^{\frac{j}{2}} T_j\left(t, z^t, \zeta^t, \frac{W_{z^t}^m}{h^{\frac{1}{2}}}\right)$$

*for times  $|t| \leq C_0|\ln h|$  with  $T_j(t, z^t, \zeta^t, \cdot)$  polynomial (i.e. sum of tensors) of degree at most  $3j$ , depending on the classical trajectory and the Taylor expansions of  $V$  and  $W^m$  at  $z^t$ .*

# Wave packets on Riemannian manifolds

**Remark on the proof:** The transport equations

## Wave packets on Riemannian manifolds

**Remark on the proof:** The transport equations are of the form

$$(\nabla_{\dot{z}^t} T)(\underbrace{., \dots, .}_{k \text{ factors}}) + \underbrace{T[\Gamma^t \cdot, \dots] + \dots + T[\dots, \Gamma^t \cdot]}_{k \text{ terms}} = F[., \dots, .]$$

## Wave packets on Riemannian manifolds

**Remark on the proof:** The transport equations are of the form

$$(\nabla_{\dot{z}^t} T)(\underbrace{., \dots, .}_{k \text{ factors}}) + \underbrace{T[\Gamma^t, \dots] + \dots + T[\dots, \Gamma^t]}_{k \text{ terms}} = F[., \dots, .]$$

which turns out to be equivalent to

$$\frac{d}{dt} (T[E_t, \dots, E_t]) = F[E_t, \dots, E_t]$$

with  $E_t := d\pi(\mathcal{L}_A + i\mathcal{L}_B) : \mathbb{C}^n \rightarrow T_{z^t}M \otimes \mathbb{C}$  ( $d\pi =$  projection from the horizontal space at  $(z^t, \dot{z}^t)$  to the tangent space at  $z^t$ )

$\implies$  Control on the exponential growth in time of  $T_j(t, z^t, \zeta^t, .)$ .

# Wave packets on Riemannian manifolds

## Theorem [Propagator approximation]

## Wave packets on Riemannian manifolds

**Theorem [Propagator approximation]** *If  $A_h$  is a pseudodifferential operator supported in  $U$ ,*

## Wave packets on Riemannian manifolds

**Theorem [Propagator approximation]** *If  $A_h$  is a pseudodifferential operator supported in  $U$ , with principal symbol  $\chi$ ,*

## Wave packets on Riemannian manifolds

**Theorem [Propagator approximation]** *If  $A_h$  is a pseudodifferential operator supported in  $U$ , with principal symbol  $\chi$ , then (the kernel of)  $e^{-i\frac{\epsilon}{h}H(h)}A_h$*



## Wave packets on Riemannian manifolds

**Theorem [Propagator approximation]** *If  $A_h$  is a pseudodifferential operator supported in  $U$ , with principal symbol  $\chi$ , then (the kernel of)  $e^{-i\frac{t}{h}H(h)}A_h$  is well approximated by*

$$K_t^h(m, m') = h^{-\frac{3n}{2}} \int \int_{T^*U} b_h(t, z, \zeta, m, m') \exp \frac{i}{h} F(t, z, \zeta, m, m') dz d\zeta$$

for times  $|t| \leq C_0 |\log h|$ .

## Wave packets on Riemannian manifolds

**Theorem [Propagator approximation]** If  $A_h$  is a pseudodifferential operator supported in  $U$ , with principal symbol  $\chi$ , then (the kernel of)  $e^{-i\frac{t}{h}H(h)}A_h$  is well approximated by

$$K_t^h(m, m') = h^{-\frac{3n}{2}} \int \int_{T^*U} b_h(t, z, \zeta, m, m') \exp \frac{i}{h} F(t, z, \zeta, m, m') dz d\zeta$$

for times  $|t| \leq C_0 |\log h|$ . The **phase** reads

$$F = S_{(z, \zeta)}^t + \zeta^t \cdot W_{z^t}^m + \frac{1}{2} \left\langle \Gamma_{(z, \zeta)}^t W_{z^t}^m, W_{z^t}^m \right\rangle_{z^t} - \zeta \cdot W_z^{m'} + \frac{1}{2} \left\langle \widetilde{\Gamma_{(z, \zeta)}^0} W_z^{m'}, W_z^{m'} \right\rangle_z$$

## Wave packets on Riemannian manifolds

**Theorem [Propagator approximation]** If  $A_h$  is a pseudodifferential operator supported in  $U$ , with principal symbol  $\chi$ , then (the kernel of)  $e^{-i\frac{t}{h}H(h)}A_h$  is well approximated by

$$K_t^h(m, m') = h^{-\frac{3n}{2}} \int \int_{T^*U} b_h(t, z, \zeta, m, m') \exp \frac{i}{h} F(t, z, \zeta, m, m') dz d\zeta$$

for times  $|t| \leq C_0 |\log h|$ . The **phase** reads

$$F = S_{(z, \zeta)}^t + \zeta^t \cdot W_{z^t}^m + \frac{1}{2} \left\langle \Gamma_{(z, \zeta)}^t W_{z^t}^m, W_{z^t}^m \right\rangle_{z^t} - \zeta \cdot W_z^{m'} + \frac{1}{2} \left\langle \widetilde{\Gamma_{(z, \zeta)}^0} W_z^{m'}, W_z^{m'} \right\rangle_z$$

where

$$\left\langle \widetilde{\Gamma_{(z, \zeta)}^0} W_z^{m'}, W_z^{m'} \right\rangle_z = -\operatorname{Re} \left\langle \Gamma_{(z, \zeta)}^0 W_z^{m'}, W_z^{m'} \right\rangle_z + i \operatorname{Im} \left\langle \Gamma_{(z, \zeta)}^0 W_z^{m'}, W_z^{m'} \right\rangle_z$$

## Wave packets on Riemannian manifolds

**Theorem [Propagator approximation]** If  $A_h$  is a pseudodifferential operator supported in  $U$ , with principal symbol  $\chi$ , then (the kernel of)  $e^{-i\frac{t}{h}H(h)}A_h$  is well approximated by

$$K_t^h(m, m') = h^{-\frac{3n}{2}} \int \int_{T^*U} b_h(t, z, \zeta, m, m') \exp \frac{i}{h} F(t, z, \zeta, m, m') dz d\zeta$$

for times  $|t| \leq C_0 |\log h|$ . The **phase** reads

$$F = S_{(z, \zeta)}^t + \zeta^t \cdot W_{z^t}^m + \frac{1}{2} \left\langle \Gamma_{(z, \zeta)}^t W_{z^t}^m, W_{z^t}^m \right\rangle_{z^t} - \zeta \cdot W_z^{m'} + \frac{1}{2} \left\langle \widetilde{\Gamma_{(z, \zeta)}^0} W_z^{m'}, W_z^{m'} \right\rangle_z$$

where

$$\left\langle \widetilde{\Gamma_{(z, \zeta)}^0} W_z^{m'}, W_z^{m'} \right\rangle_z = -\operatorname{Re} \left\langle \Gamma_{(z, \zeta)}^0 W_z^{m'}, W_z^{m'} \right\rangle_z + i \operatorname{Im} \left\langle \Gamma_{(z, \zeta)}^0 W_z^{m'}, W_z^{m'} \right\rangle_z$$

The **amplitude**  $b_h(t, z, \zeta, m, m')$  reads  $b_0(t, z, \zeta, m, m') + O_t(h^{1/2})$ ,

$$b_0 = \det((g_{jk}(x^t))^{1/2} (A^t + iB^t))^{-\frac{1}{2}} \det(g_{jk}(y))^{-\frac{1}{4}} \chi(z, \zeta) \rho(d_g(z, m')) \rho(d_g(z^t, m))$$

## Wave packets on Riemannian manifolds

**Theorem [Propagator approximation]** If  $A_h$  is a pseudodifferential operator supported in  $U$ , with principal symbol  $\chi$ , then (the kernel of)  $e^{-i\frac{t}{h}H(h)}A_h$  is well approximated by

$$K_t^h(m, m') = h^{-\frac{3n}{2}} \int \int_{T^*U} b_h(t, z, \zeta, m, m') \exp \frac{i}{h} F(t, z, \zeta, m, m') dz d\zeta$$

for times  $|t| \leq C_0 |\log h|$ . The **phase** reads

$$F = S_{(z, \zeta)}^t + \zeta^t \cdot W_{z^t}^m + \frac{1}{2} \left\langle \Gamma_{(z, \zeta)}^t W_{z^t}^m, W_{z^t}^m \right\rangle_{z^t} - \zeta \cdot W_z^{m'} + \frac{1}{2} \left\langle \widetilde{\Gamma_{(z, \zeta)}^0} W_z^{m'}, W_z^{m'} \right\rangle_z$$

where

$$\left\langle \widetilde{\Gamma_{(z, \zeta)}^0} W_z^{m'}, W_z^{m'} \right\rangle_z = -\operatorname{Re} \left\langle \Gamma_{(z, \zeta)}^0 W_z^{m'}, W_z^{m'} \right\rangle_z + i \operatorname{Im} \left\langle \Gamma_{(z, \zeta)}^0 W_z^{m'}, W_z^{m'} \right\rangle_z$$

The **amplitude**  $b_h(t, z, \zeta, m, m')$  reads  $b_0(t, z, \zeta, m, m') + O_t(h^{1/2})$ ,

$$b_0 = \det((g_{jk}(x^t))^{1/2} (A^t + iB^t))^{-\frac{1}{2}} \det(g_{jk}(y))^{-\frac{1}{4}} \chi(z, \zeta) \rho(d_g(z, m')) \rho(d_g(z^t, m))$$

**Proof:**

$$e^{-i\frac{t}{h}H(h)}A_h u = (2\pi h)^{-n} \int \int_{T^*U} e^{-i\frac{t}{h}H(h)} \psi_{z, \zeta}^h \left\langle A_h^* a_h^{-1} \psi_{z, \zeta}^h, u \right\rangle_{L^2(M)} dz d\zeta$$

**Thank you for your attention**